

Finite temperature critical point in heavy-quark QCD

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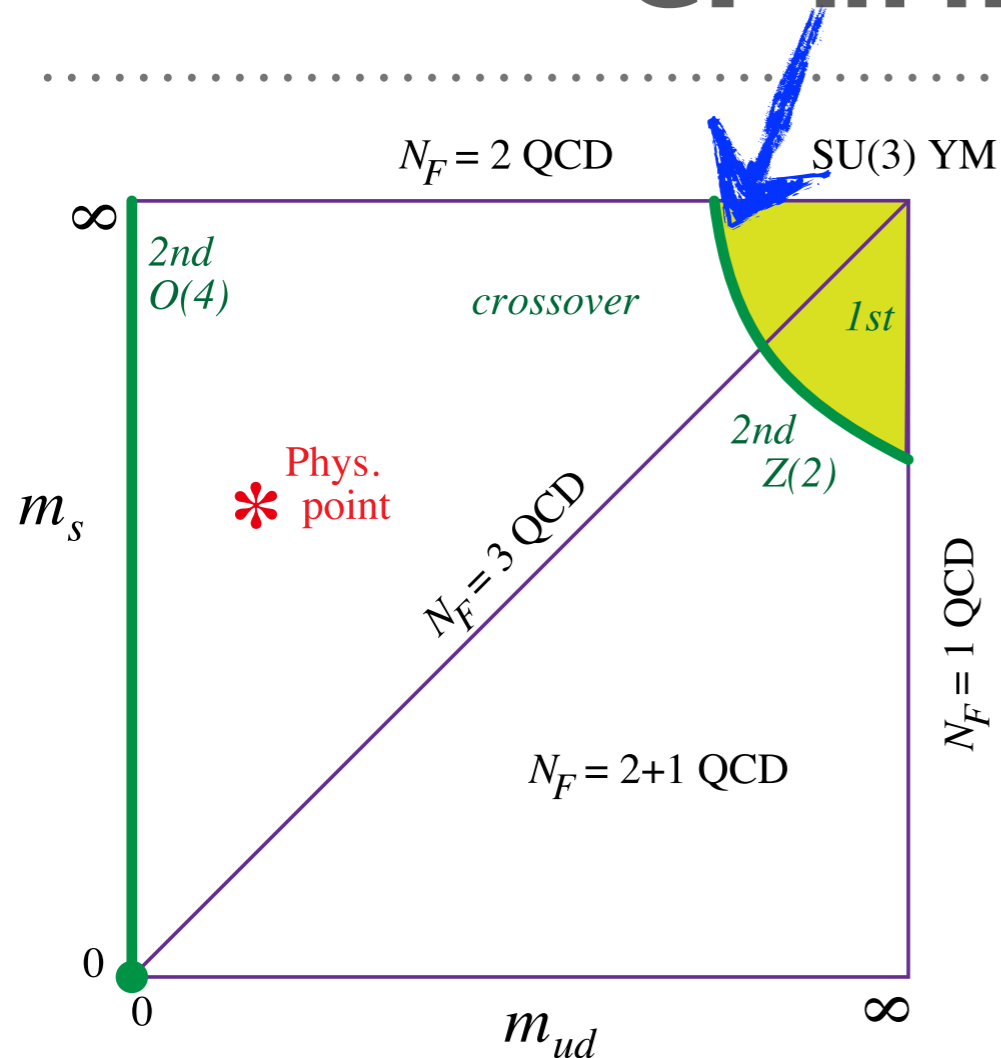
with R. Ashikawa²⁾, S. Ejiri³⁾, M. Kitazawa^{4,5)}, H. Sugawara³⁾

(WHOT-QCD Collaboration)

1) Univ. Tsukuba, 2) Osaka Univ., 3) Niigata Univ., 4) Kyoto Univ., 5) KEK J-PARC branch



CP in heavy-quark QCD



Finite- T QCD transition is important in understanding various phenomena: the early evolution of the Universe, neutron stars, quark matter in relativistic heavy-ion collisions, etc.

\Leftarrow Columbia plot summarizing the nature of the QCD trans. as function of m_q 's.

We focus on the CP near the upper-right corner (hq-QCD).

A powerful way to determine CP:

Binder cumulant analysis based on the expected $Z(2)$ FSS,

assuming (approx.) dominance of the leading sing. in $L \rightarrow \infty$ lim.

- ▶ **Large spatial lattices & high statistics**
to identify the leading FSS clearly.
- ▶ **Reweighting to vary coupling parameters continuously**
as required by Binder cumulant analyses.

Lattice setup

▶ Lattice action: plaquette gauge + standard Wilson quarks

▶ Wilson quark kernel: $M_{xy}(\kappa) = \delta_{xy} - \kappa \sum_{\mu} \left[(1 - \gamma_{\mu}) U_{x,\mu} \delta_{y,x+\hat{\mu}} + (1 + \gamma_{\mu}) U_{y,\mu}^{\dagger} \delta_{y,x-\hat{\mu}} \right]$

$$= \delta_{xy} - \kappa B_{xy} \quad \text{where } \kappa = 1/(2am_q + 8)$$

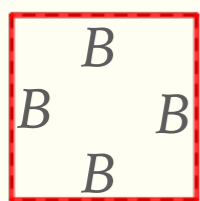
▶ Quark contribution to the effective action: $\ln \det M(\kappa) = -\frac{1}{N_{\text{site}} n} \sum_{n=1}^{\infty} \text{Tr}[B^n] \kappa^n$

● each term given by closed loops of B with κ [loop length]

▶ Hopping Parameter Expansion

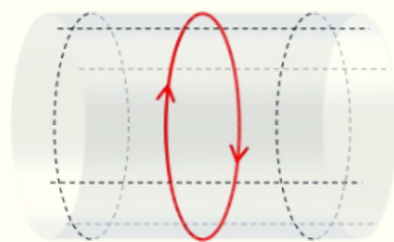
LO:

plaquette



$$\kappa^4$$

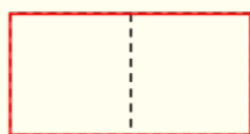
Polyakov loop Ω



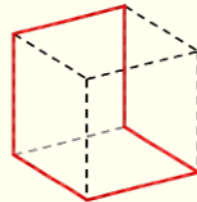
$$\kappa^{N_t}$$

NLO:

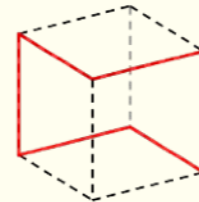
(a) rectangle



(b) chair



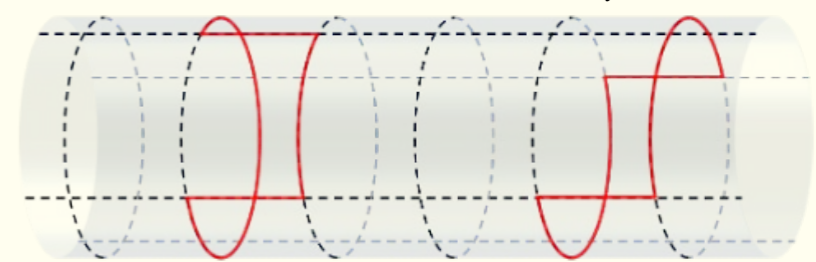
(c) crown



$$\kappa^6$$

Wilson loops

bent Polyakov loops Ω_i



$$\kappa^{N_t+2}$$

Polyakov-loop type loops

● HPE $\approx 1/(am_q)$ expansion

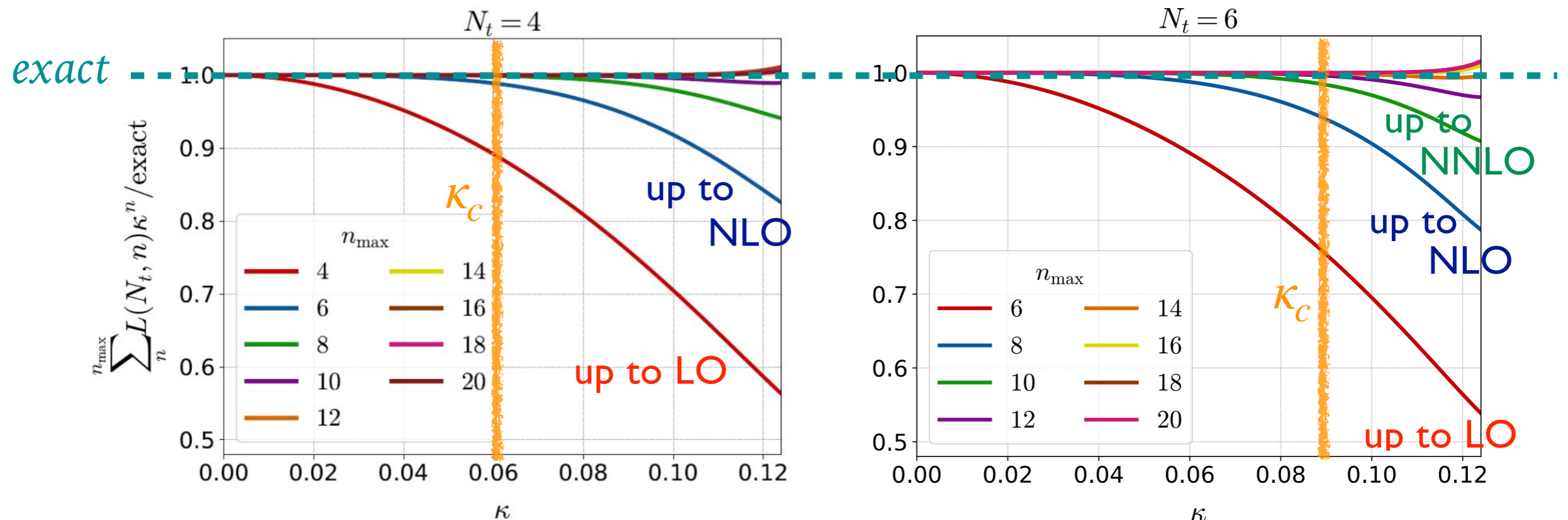
● HPE worsens with $a \rightarrow 0$ ($N_t \rightarrow \infty$) \Rightarrow higher orders required.

Convergence of HPE

Wakabayashi, Ejiri, KK, Kitazawa, PTEP 2022, 033B05 (2022)

Deviation from true value due to truncation of HPE in the worst convergent case:

Polyakov-loop type contribution to the eff. action: (Wilson-loop type contribution similar.)

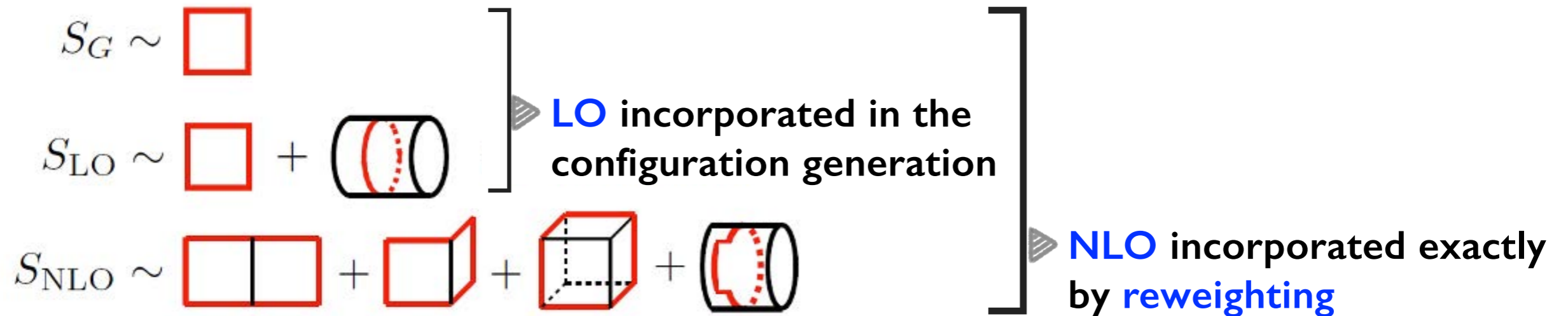


For studies around CP:

Ashikawa+, arXiv:2407.09156 (2024)

- **Nt=4:** $\kappa_c = 0.0603(4)$ [Kiyohara+ ('21)]
 \Rightarrow **LO may have at worst 10% error in the eff. action, NLO is $\geq 98\%$ accurate.**
- **Nt=6:** $\kappa_c = 0.08769(7)_{-0}^{+11}$ [Ashikawa+ ('24)], $0.0877(9)$ [Cuteri+ ('22)]
 \Rightarrow **NLO is $\geq 93\%$ accurate, NNLO is $\geq 97\%$ accurate.**
- **Nt=8:** $\kappa_c = 0.09024(46)$ [Sugawara+ ('24)], $0.1135(8)$ [Cuteri+ ('22)] \Rightarrow **NNLO for $> 95\%$**



Simulation up to NLO exactly



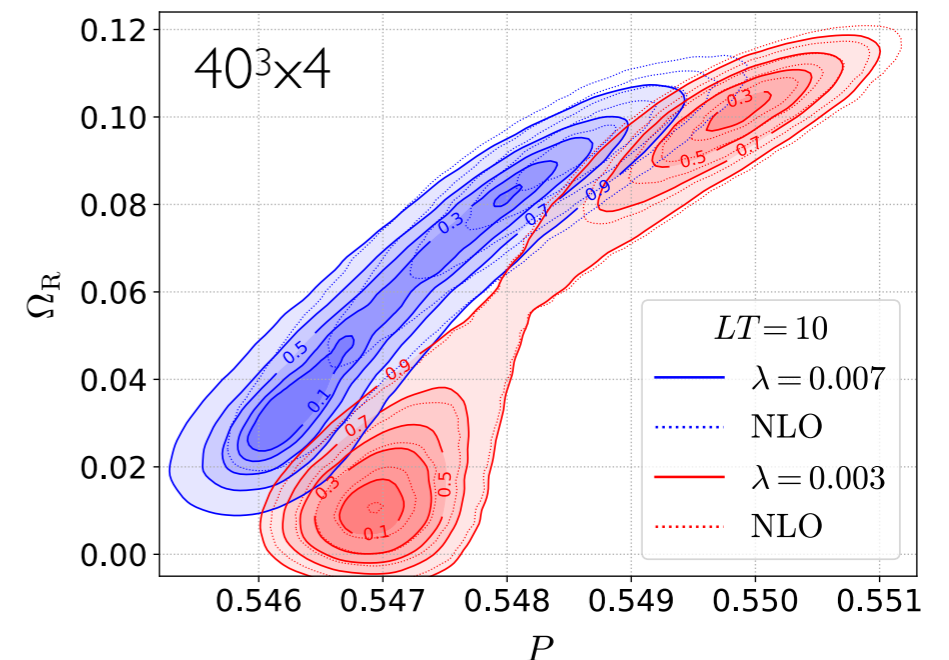
$$\beta \rightarrow \beta^* = \beta + 48N_f\kappa^4$$

$\lambda \sum_{\mathbf{x}} \Omega(\mathbf{x})$ term can be incorporated in PHB+OR algorithm *à la* pure YM simulations

Kiyohara, Kitazawa, Ejiri, KK PRD 104, 1144509 (2021)

-  Simulation cost \ll f-QCD simulations
 -  Overlap problem of reweighting resolved by the LO in configuration generation
- \Rightarrow 1st order transition on large lattices

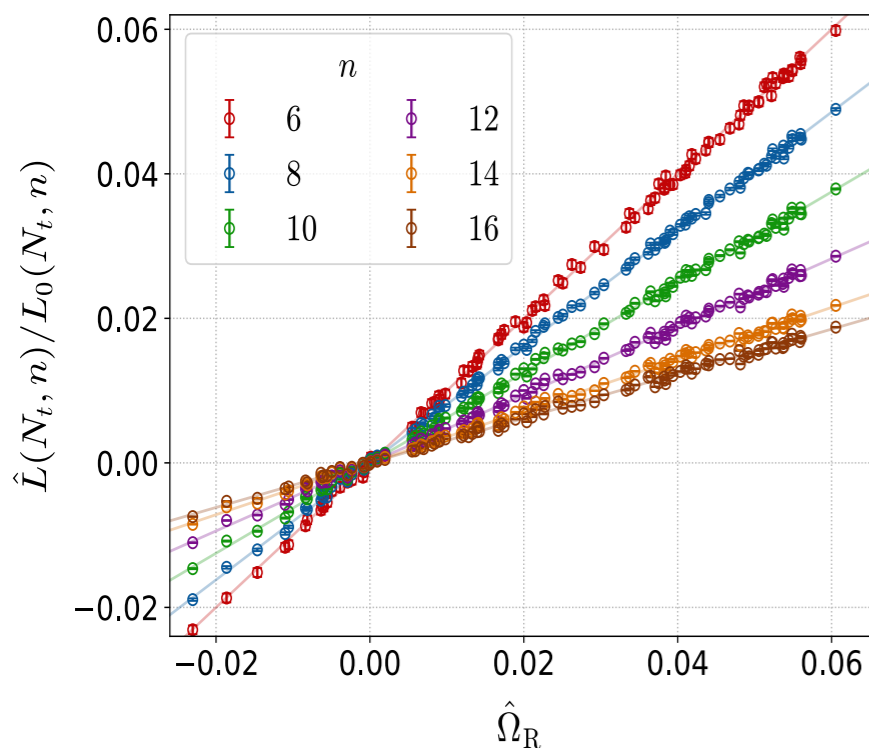
Kiyohara+ (2021)



Effective incorporation of higher orders

Basic observation:

strong linear correlation among different order terms of the HPE.



Scatter plot of n-th order Polyakov-loop type terms of HPE vs LO Polyakov-loop, observed on an $N_t=6$ lattice near the CP.

Ashikawa,+ (2024)

We may approximate

measured from the slopes

n-th order term $\approx C_n \times$ low-order term

Wilson-loop type terms show weaker but similar correlation.

► **eff.[LO] method**

Exact up to LO.

Effectively incorporate NLO and higher orders by shifting the couplings in S_{G+LO} .

Wakabayashi, Ejiri, KK, Kitazawa, PTEP 2022, 033B05 (2022)

► **eff.[NLO] method**

Exact up to NLO.

Effectively incorporate NNLO and higher by shifting the NLO couplings in S_{NLO} .

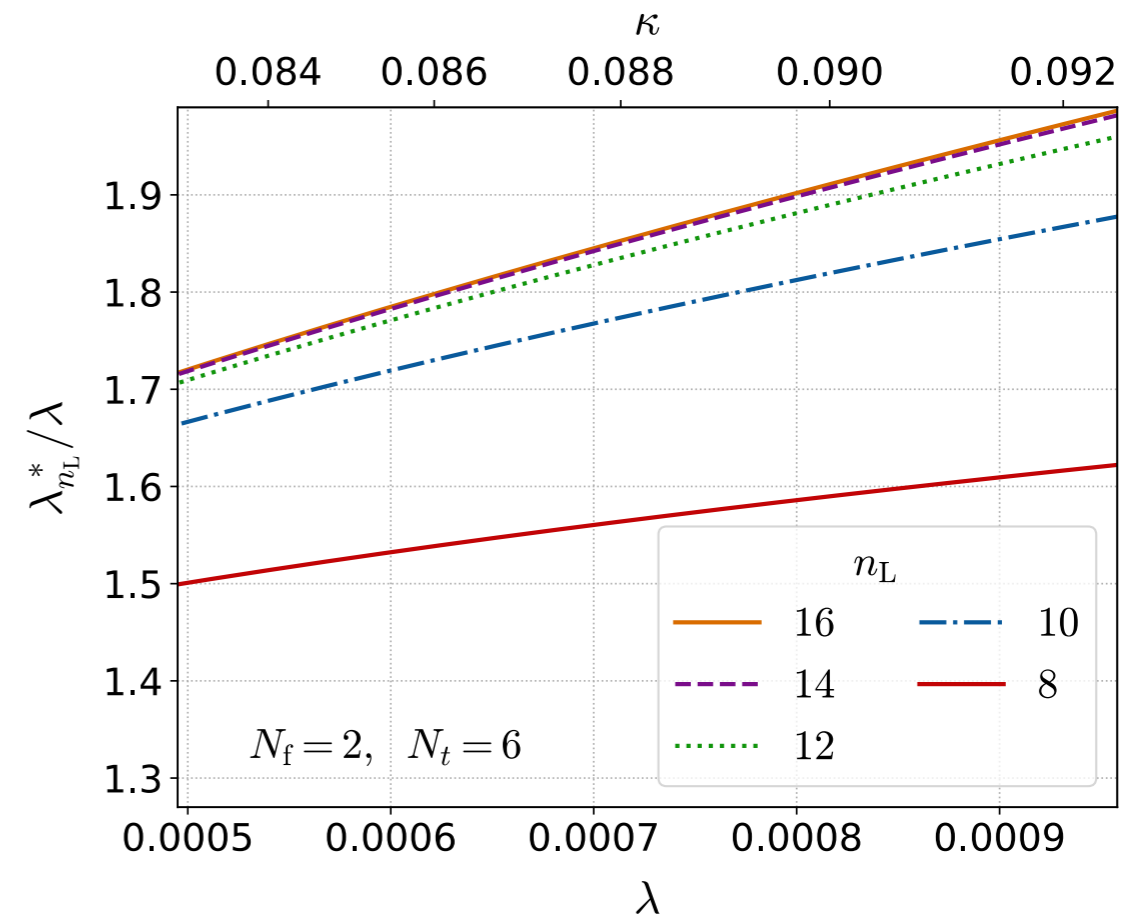
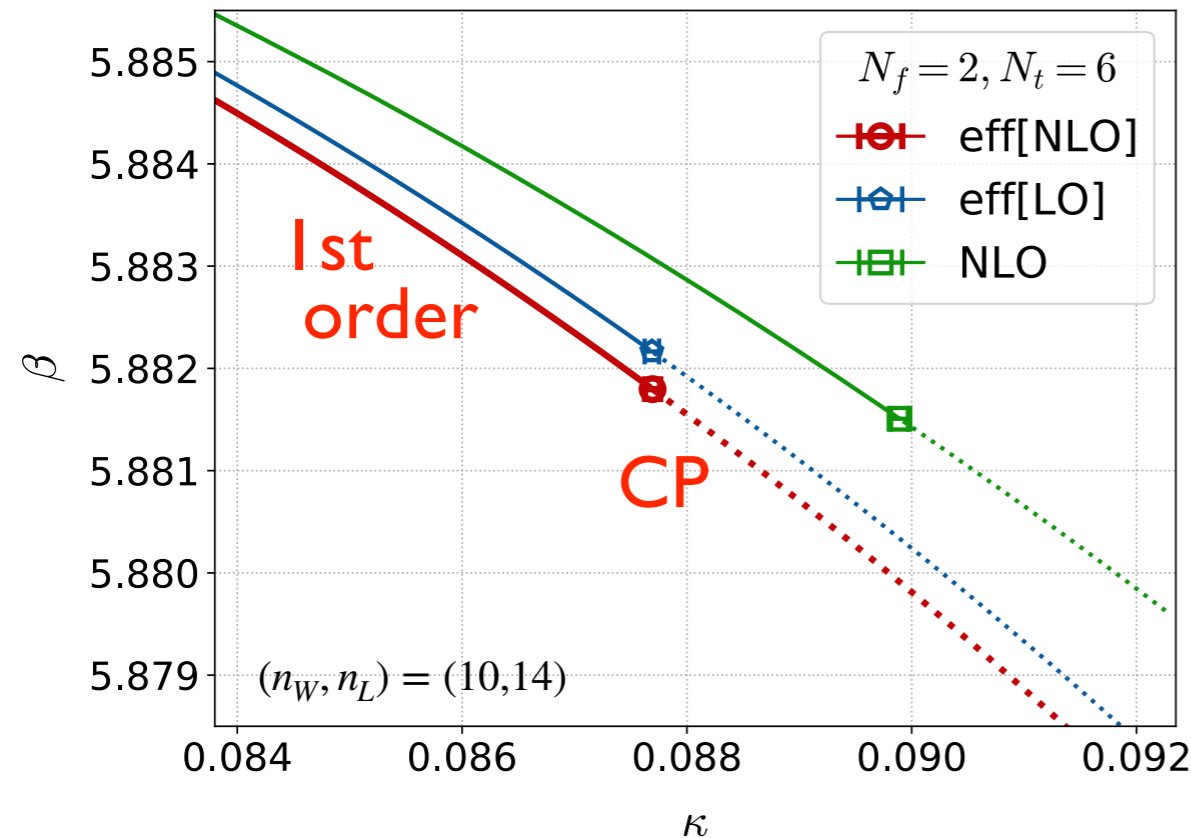
Better because NLO is exact and correlation is stronger with smaller order-differences.

Ashikawa, Kitazawa, Ejiri, KK, arXiv:2407.09156 (2024)

Test of eff. methods using final observables

Results of the phase diagram at $Nt = 6$ (discussed later)

Ashikawa, Kitazawa, Ejiri, KK, arXiv:2407.09156 (2024)



- Transition line and CP with $\text{eff.}[LO/NLO]$ shift from NLO
- ⇒ **NNLO and higher orders important at $Nt \geq 6$.**
- $\text{eff.}[LO] \approx \text{eff.}[NLO]$
- ⇒ **effective incorporation of NLO works**

- Dependence on the truncation order in $\text{eff.}[LO/NLO]$
- ⇒ **convergent if $(n_W, n_L) \geq (10, 14)$ for $Nt=6$.**

We adopt $(n_W, n_L) = (10, 14)$ at $Nt=6$.

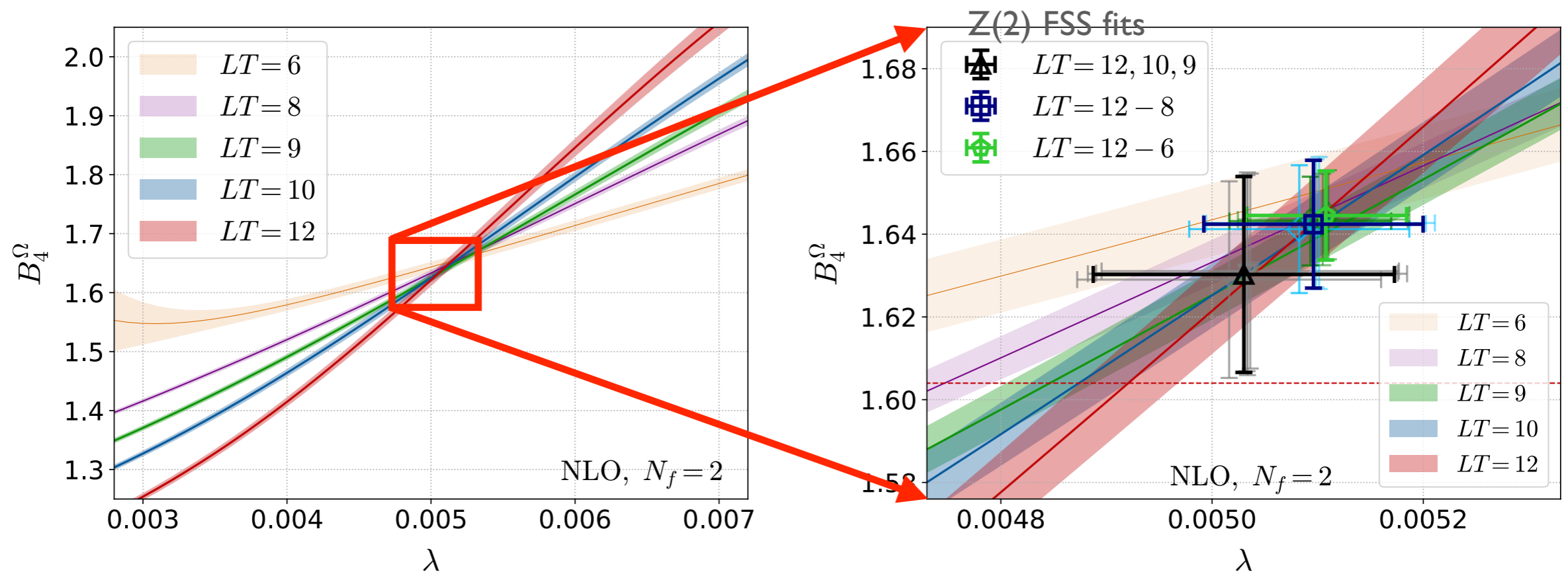
Study at $N_t = 4$

Kiyohara, Kitazawa, Ejiri, KK, PRD 104 (2021)

Binder cumulant of the Polyakov loop $B_4^\Omega = \frac{\langle \Omega_R^4 \rangle_c}{\langle \Omega_R^2 \rangle_c^2} + 3$ along the transition line

► Results at $N_t = 4$ with HPE up to NLO

$\lambda = 48N_f N_t \kappa^4$ for $N_t=4$



★ large lattices with $N_s/N_t = LT \geq 9$ required for Z(2) FSS

★ $\nu = 0.614(48)(3)$, $b_4 = 1.630(24)(2)$ using $N_s/N_t = 9-12$

consistent with the Z(2) values 0.630, 1.604 within 1σ

★ $\lambda_c = 0.00503(14)(2)$ [$\kappa_c = 0.0603(4)$] for $N_t=4, N_f=2$

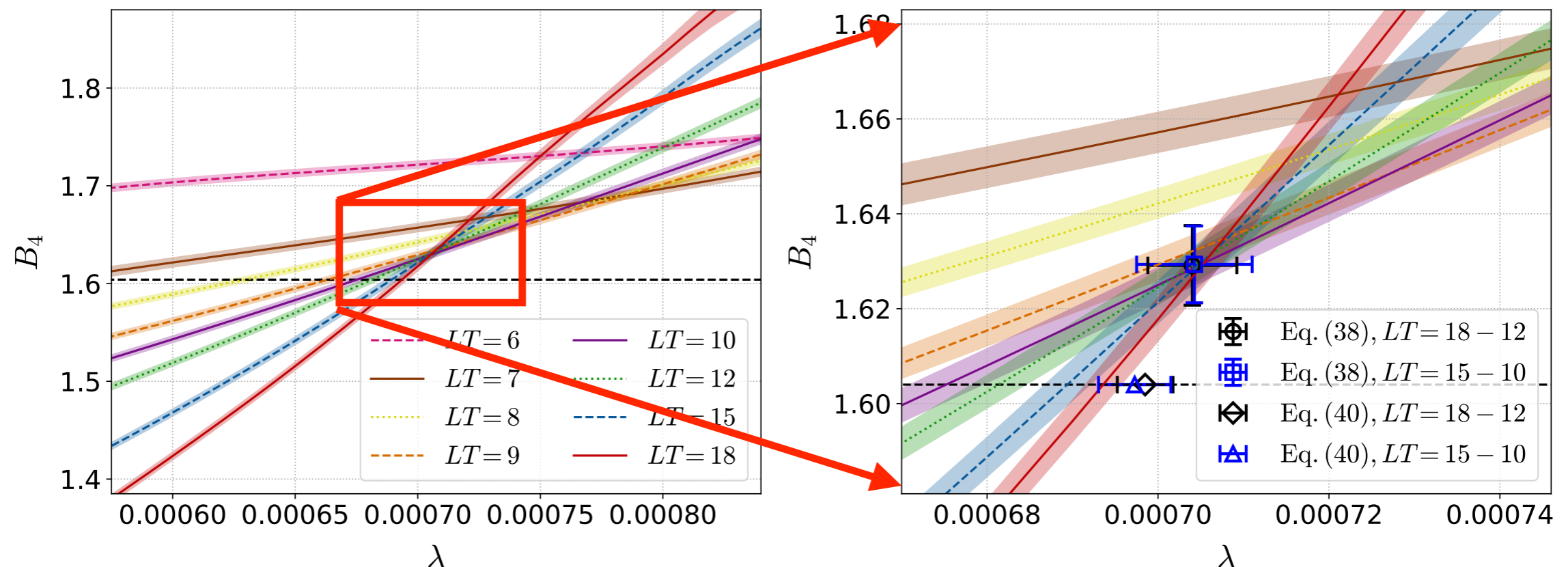
Study at $N_t = 6$

Ashikawa, Kitazawa, Ejiri, KK, arXiv:2407.09156 (2024)

▶ $N_t=6$, $N_s/N_t = LT = 6, 7, 8, 9, 10, 12, 15, 18$ ($N_s = 36-108$)

▶ B_4^Ω

$\lambda = 128N_f N_t \kappa^6$ for $N_t=6$



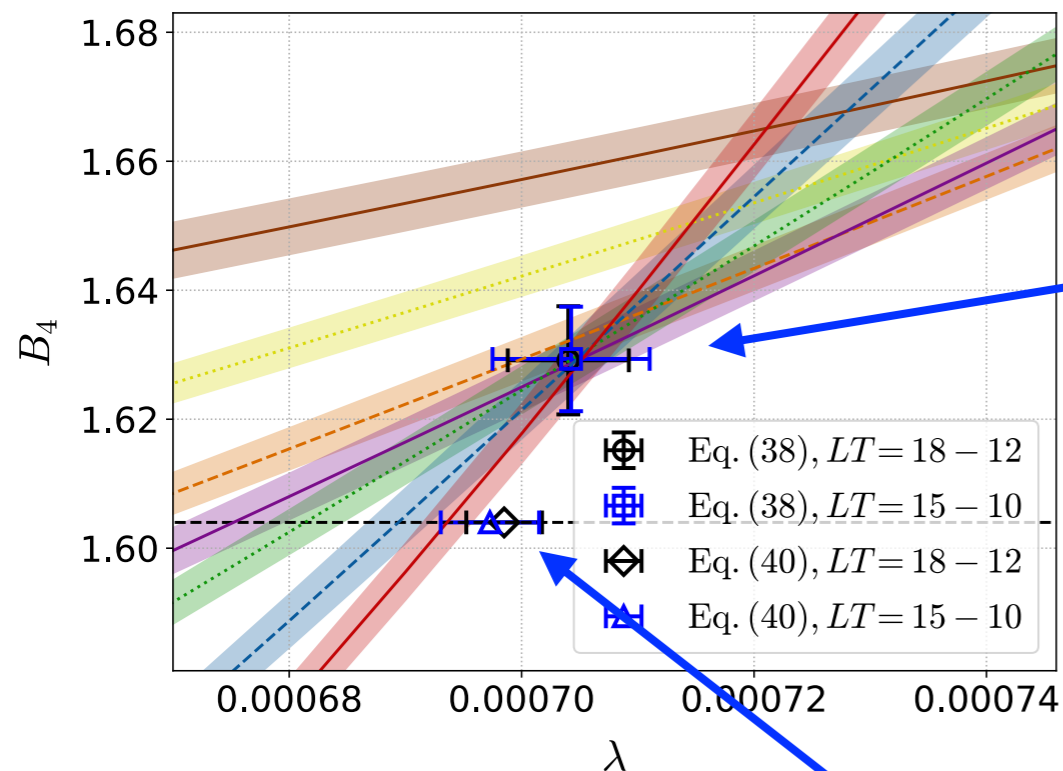
★ Violation of FSS larger on finer lattice \Rightarrow larger LT (≥ 10) required

★ $\nu = 0.627(19)(5)$, $b_4 = 1.6297(84)(6)$ with $LT = 12-18$

(cf.) $Z(2)$ values: 0.630, 1.604

Study at $N_t = 6$

Ashikawa, Kitazawa, Ejiri, KK, arXiv:2407.09156 (2024)



FSS fit using standard ansatz for mag. op.:

$$B_4 = (b_4 + c(\lambda - \lambda_c)(LT)^{-1/\nu})$$

★ $\nu = 0.627(19)(5)$, $b_4 = 1.6297(84)(6)$
[cf. Z(2) values: 0.630, 1.604]

FSS fit taking mixing with energy-like op.:

$$B_4 = (b_4 + c(\lambda - \lambda_c)(LT)^{-1/\nu})(1 + d(LT)^Y)$$

$Y = -0.896$ for Z(2)

Full 6 parameter fits were unstable

=> fit fixing b_4, ν, Y to Z(2) values

★ fits acceptable when $LT \geq 10$.

note: w/ mixing, the crossing pt. moves with LT .

★ $\kappa_c = 0.08769(7) \left(\begin{smallmatrix} +11 \\ -0 \end{smallmatrix} \right)$

$LT=12-18$, eff.[NLO], $N_F=2$

[cf. 0.0877(9) Cuteri+ (21) fQCD, $LT=4-7$]

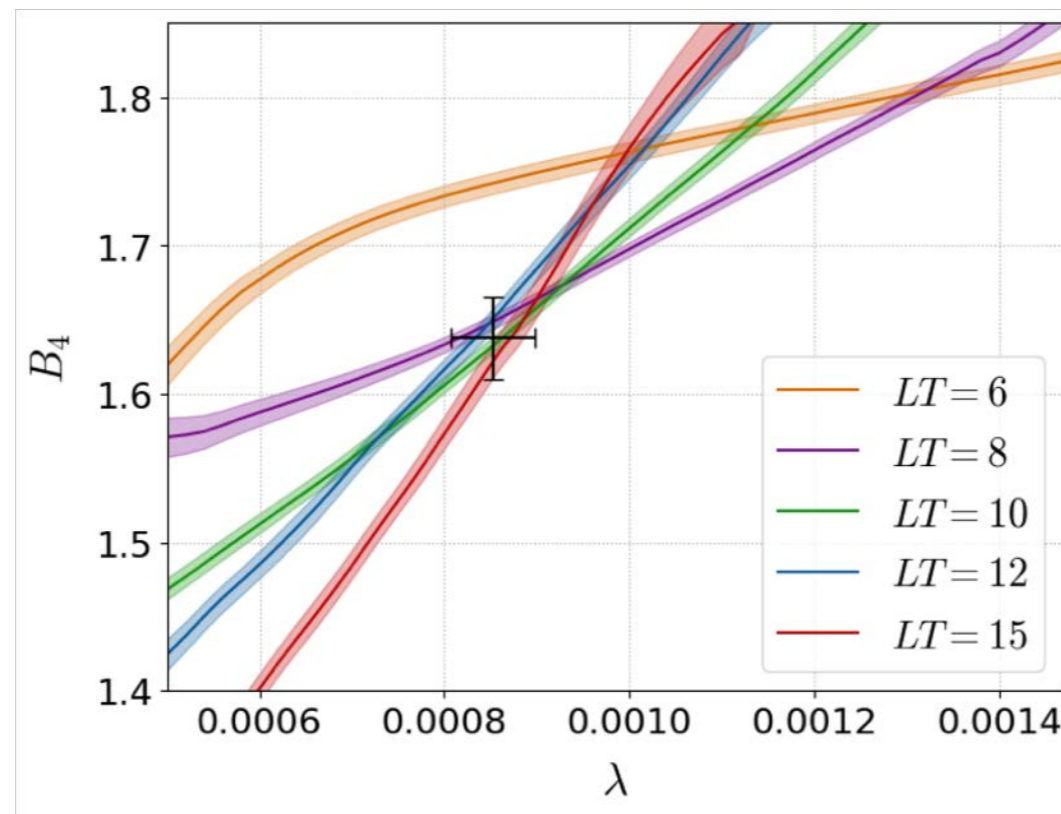
LT	$\lambda_c \times 10^4$	χ^2/dof
18, 15, 12	6.986(32)(9)	0.51
15, 12, 10	6.973(42)(5)	0.83
18 – 10	6.984(25)(8)	1.07
18 – 9	6.998(21)(5)	4.23
18 – 8	7.035(19)(3)	36.4

Study at $N_t = 8$ (in preparation)

Sugawara+, in preparation

▶ $N_t=8$, $N_s/N_t = LT = 6, 8, 10, 12, 15$ ($N_s = 48 - 120$)

▶ B_4^Ω



$$\lambda = 384 N_f N_t \kappa^8 \text{ for } N_t=8$$

★ $\nu = 0.72(27)$, $b_4 = 1.638(28)$ [$\chi^2/\text{dof} = 1.474$] with $LT = 10-15$

★ $b_4 = 1.637(24)$ fixing $\nu = 0.630$ [$\chi^2/\text{dof} = 1.343$] with $LT = 10-15$

consistent with $Z(2)$ within 1σ

★ $\kappa_c = 0.09024(46)$ for $N_F=2$ with eff.[LO] up to 20th order

[cf. 1.1135(8) Cuteri+ ('22) fQCD, $LT=4-6(7,10)$]

Conclusions

.....
 We determined CP in heavy-quark QCD by the Binder cumulant:

★ **Large spatial lattices with $LT = N_s/N_t \gtrsim 10$ required.**

★ **HPE + eff. method to incorporate high orders.**

★ $N_t = 4$: $\kappa_c = 0.0603(4)$ for $N_F=2$ using $LT=9-12$

★ $N_t = 6$: $\kappa_c = 0.08769(7)_{-0}^{+11}$ for $N_F=2$ using $LT=10-18$

★ $N_t = 8$: $\kappa_c = 0.09024(46)$ for $N_F=2$ using $LT=10-15$

★ N_F -dep. known analytically in HPE \Rightarrow easy to translate them to 2+1 flavors etc.

★ **CP in physical units using m_{PS} at $T=0$.** [Cuteri+ ('21), Itagaki+('19) unpublished]

★ $N_t = 4$: $m_{PS}^{(CP)}/T_c = 16.30(3)$ for $N_F=2$ using $LT=9-12$

★ $N_t = 6$: $m_{PS}^{(CP)}/T_c = 18.07(2)_{-2}^{+0}$ for $N_F=2$ using $LT=10-18$

★ $N_t = 8$: $m_{PS}^{(CP)}/T_c = 17.2(2)$ for $N_F=2$ using $LT=10-15$

★ N_t -dep. (a -dep.) looks small in this combination.

 The method should work at least up to $N_t \sim 10$.

 Application to finite-density QCD.

We miss

Yusuke Taniguchi

07.04.1968 — 22.07.2022

who was joining the German-Japanese Seminars since 2010, and was the main coordinator of the Seminars in the Japanese side since 2018 until he went into a coma in Nov. 2019.

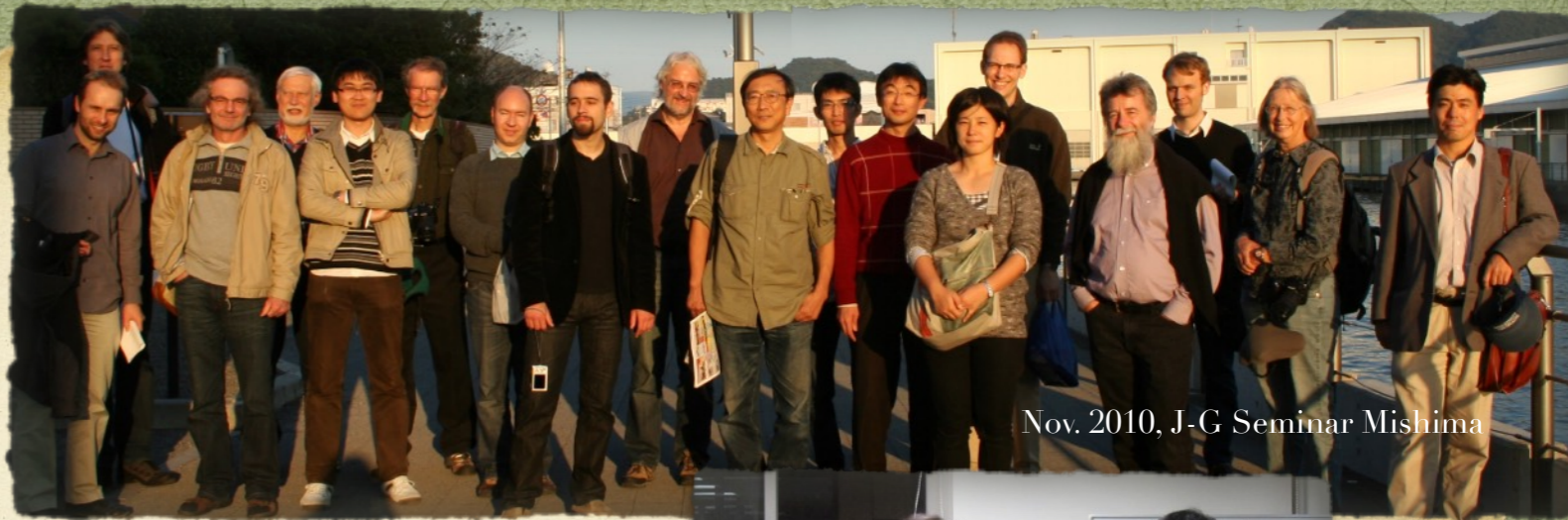
His application of the Seminar for FY 2020 at Mainz was successful, but the Seminar was postponed and finally withdrawn due to the COVID-19 pandemic.



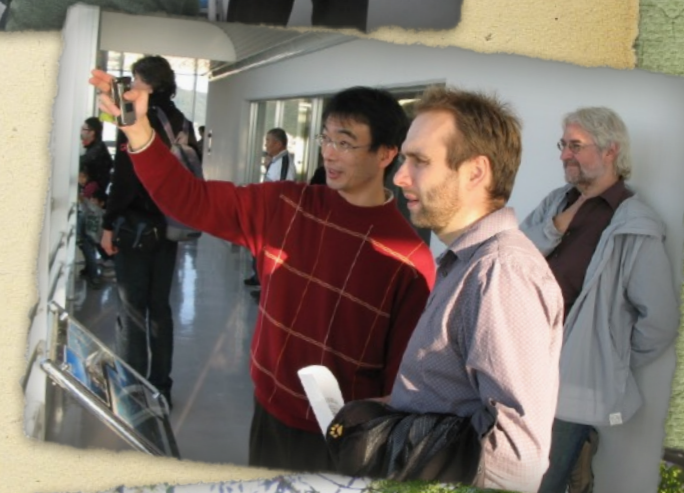
15.03.2019, Fukuoka, Japan

brief history of YT

- 1997 PhD, Kyoto University.
- 1997-1999 JSPS fellow (Univ. Tsukuba)
- 1999-2008 Assistant, Univ. Tsukuba
- 2003-2004 JSPS Visiting Scientist (CERN)
- 2008-2018 Assistant Professor (Lecturer), Univ. Tsukuba
- 2010 Japanese-German Seminar, Mishima, Japan
- 2013 German-Japanese Seminar, Regensburg, Germany
- 2016 Japanese-German Seminar, Niigata, Japan
- 2018- Associate Professor, Univ. Tsukuba



Nov. 2010, J-G Seminar Mishima



Sept. 2016, J-G Seminar Niigata



31.07.2016, Plymouth, UK



18.09.2019, Yamagata, Japan



09.09.2019, Wuhan, China



15.07.2011, Squaw Valley, USA



28.07.2016, Stone Henge, UK



20.06.2017, Granada, Spain