

Finite temperature QCD phase transition with 3 flavors of Möbius domain wall fermions

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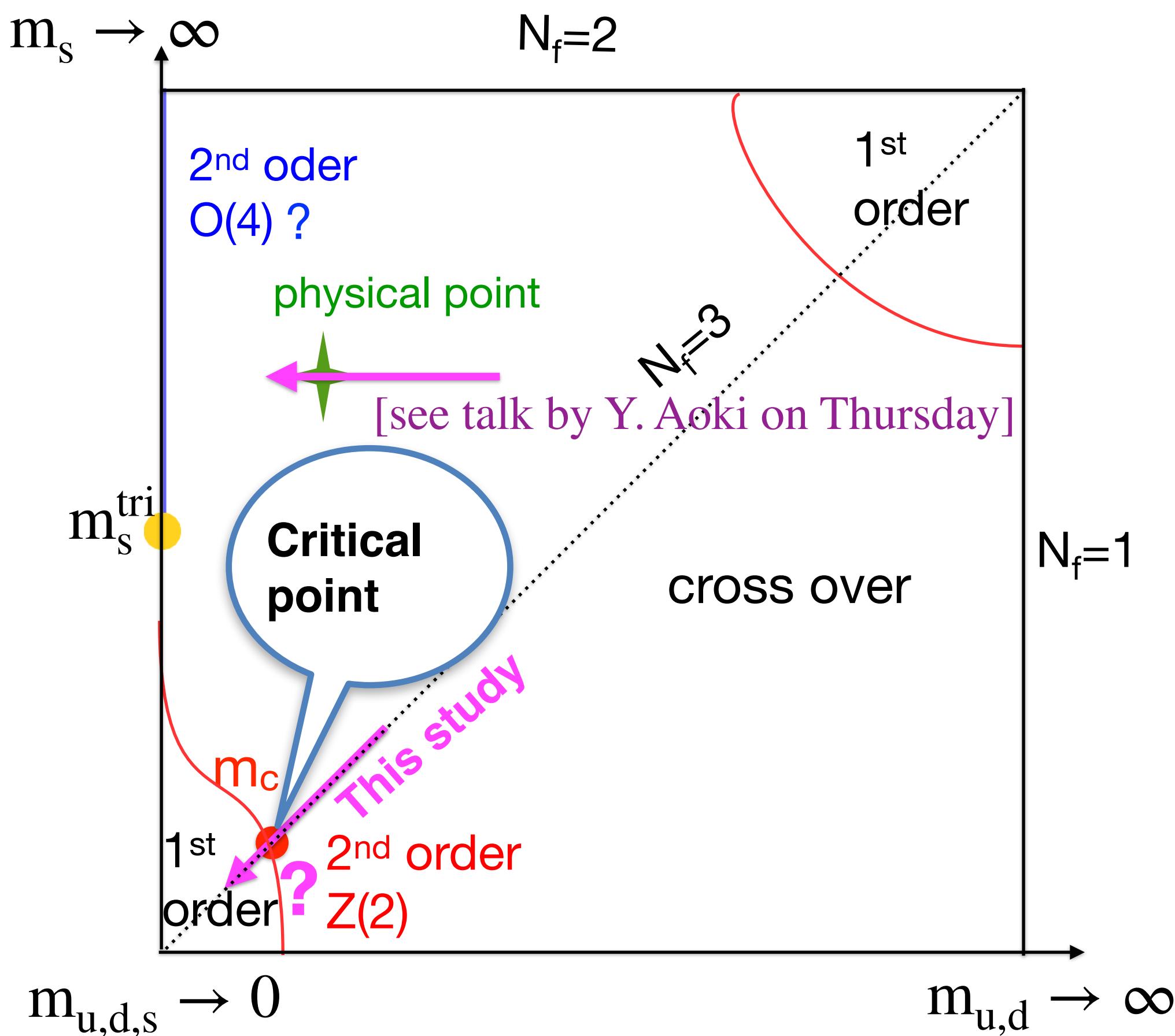
In collaboration with

Y. Aoki, S. Hashimoto, I. Kanamori, T. Kaneko, Y. Nakamura

German Japanese Workshop
Mainz, Sep 25-27, 2024

The nature of QCD phase transition at $\mu_B = 0$

Columbia plot



- ϵ expansion: 1st order phase transition in the chiral limit for $N_f = 3$
[R. D. Pisarski, F. Wilczek PRD 84]
- Possible 2nd order phase transition in the $N_f = 3$ chiral limit:

[G. Fejos, PRD 22]
[S. R. Kousvos, A. Stergiou SciPost 23]
[J. Bernhardt, C. S. Fischer PRD 23]
[R. D. Pisarski, F. Rennecke PRL 24]
[G. Fejos, T. Hatsuda PRD 24]

Need to be checked by lattice QCD

Critical point on the $N_f = 3$ chiral region:

- Location? (existence?)
- Universality class?

Previous Nf=3 lattice QCD studies

Action	N_t	$m_\pi^{Z_2}$ [MeV]	Ref.
Staggered,standard	4	290	Karsch et al. (2001)
Staggered,standard	6	150	de Forcrand et al. (2007)
Staggered, HISQ	6	$\lesssim 50$	Bazavov et al. (2017)
Staggered, stout	4-6	0?	Varnhost (2014)
Wilson, standard	4	$\lesssim 670$	Iwasaki et al. (1996)
Wilson-Clover	6-10	$\lesssim 170$	Jin et al. (2017)
Wilson-Clover	6-12	$\lesssim 110$	Kuramashi et al. (2020)

→ Strong cutoff and discretization effects

Evidence for chiral limit to feature 2nd order PT with staggered and HISQ fermion [F. Cuteri et al. JHEP 2021; S. Sharma et al. PRD 2022]
[see talk by O. Philipsen on Thursday]

We propose to use chiral fermion (Möbius domain wall fermion)

- Exact chiral symmetry at finite a for infinite Ls
- Reduced χ_{SB} parameterized by residual mass when Ls is finite

Lattice Setup

- $N_f=3$ Möbius Domain Wall Fermion
- Tree-level Symanzik improved gauge action and stout smearing

★ $T=0$:

$$\beta = 4.0, 24^3 \times 48 \times 16 : \quad 0.02 \leq am_q \leq 0.045, am_{\text{res}}(\text{estimated}) \approx 0.006$$

$$\beta = 4.1, 24^3 \times 48 \times 16 : \quad 0.015 \leq am_q \leq 0.040$$

$$\beta = 4.17, 32^3 \times 64 \times 16 : \quad 0.012 \leq am_q \leq 0.026:$$

★ $T=121(2)$ MeV ($\beta = 4.0$, ($a = 0.1361(20)$ fm, determined from Wilson flow t_0))

$$N_s^3 \times 12 \times 16 : \quad N_s = 48, -0.004 \leq am_q \leq -0.003$$

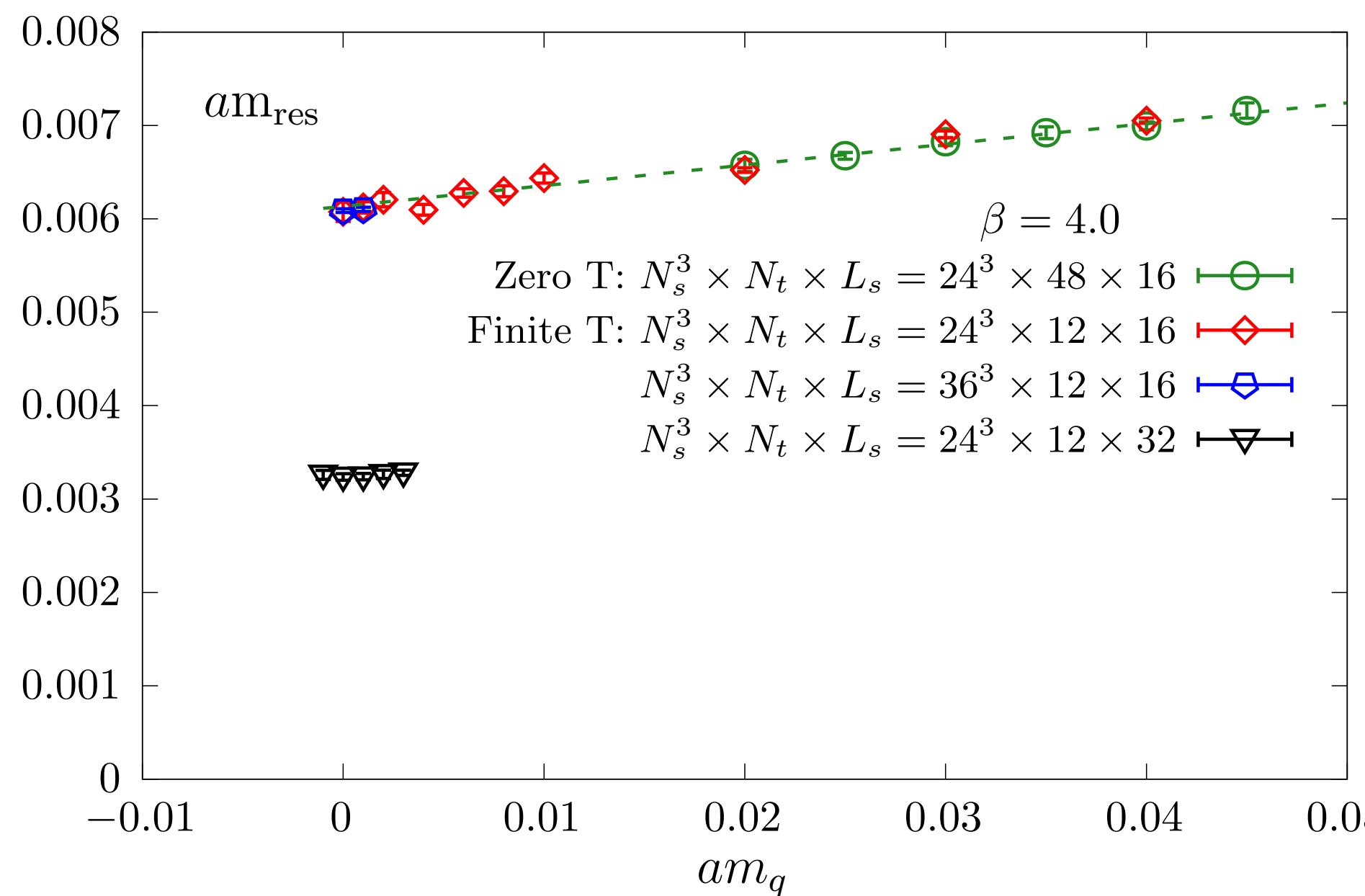
$$N_s = 36, -0.005 \leq am_q \leq 0.001$$

$$N_s = 24, -0.006 \leq am_q \leq 0.1$$

$$N_s^3 \times 12 \times 32 : \quad N_s = 24, -0.001 \leq am_q \leq 0.003$$

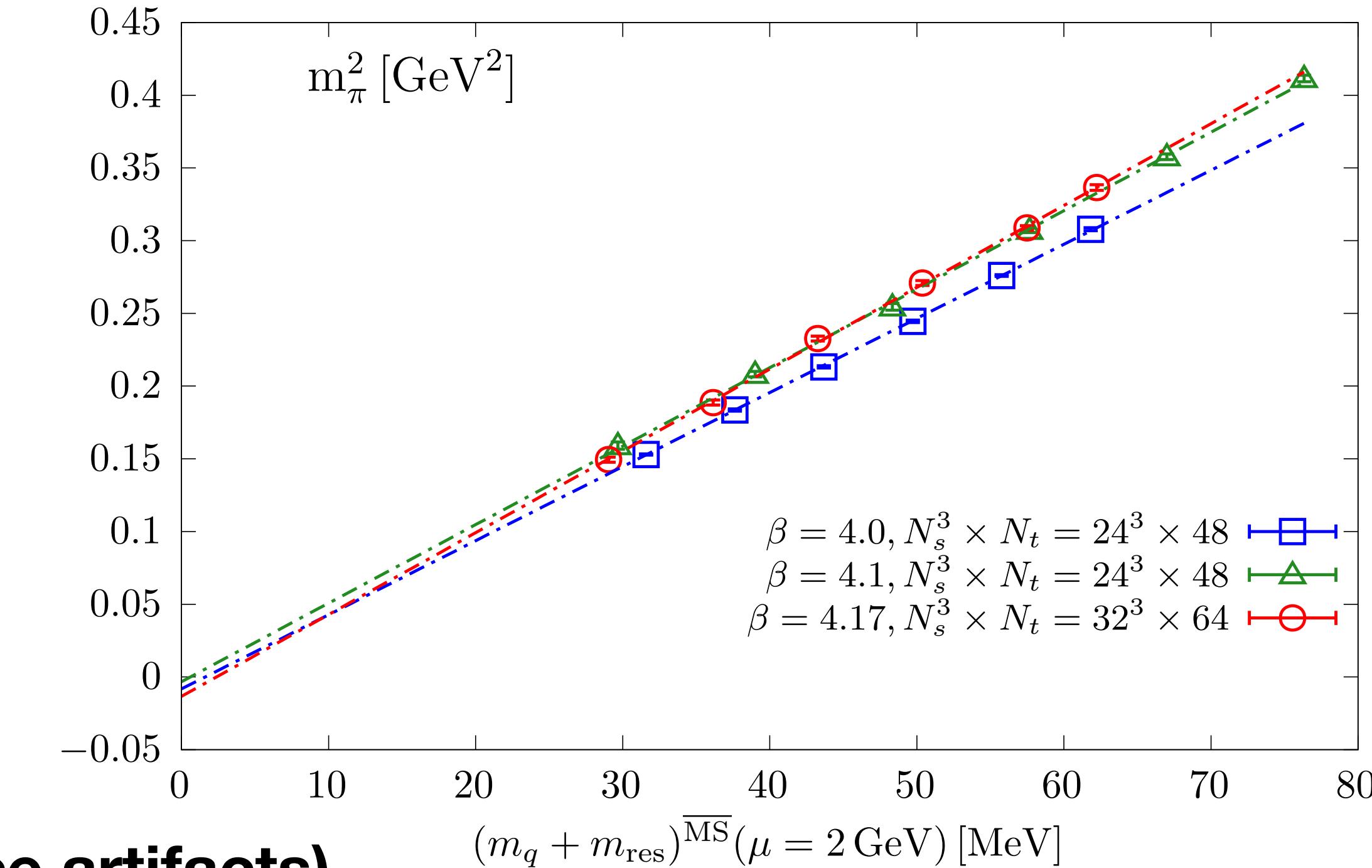
Residual chiral symmetry breaking

$$m_{\text{res}} = \left. \frac{\left\langle \sum_{\vec{x}} J_{5q}^a(\vec{x}, t) P^a(\vec{0}, 0) \right\rangle}{\left\langle \sum_{\vec{x}} P^a(\vec{x}, t) P^a(\vec{0}, 0) \right\rangle} \right|_{t \geq t_{\min}}$$



- m_{res} has a linear dependence on m_q (lattice artifacts)
 - At $am_q = 0$: $am_{\text{res}} = 0.00613(9)$ for $L_s = 16$, $am_{\text{res}} = 0.00324(3)$ for $L_s = 32$
- At strong coupling, m_{res} dominated by gauge field dislocations, suppressed by $1/L_s$
- m_π almost vanishes at chiral limit

- For finite L_s chiral symmetry is broken, leading to an additive renormalization of the mass: $m_q \rightarrow m_q + m_{\text{res}}$



Chiral condensate

$$\langle \bar{\psi} \psi \rangle|_{DWF} \sim \langle \bar{\psi} \psi \rangle|_{\text{cont.}} + C^D \frac{m_q + xm_{res}}{a^2} + \dots$$

- $x = \mathcal{O}(1)$ but $x \neq 1$ [S. Sharpe, arXiv: 0706.0218]

- Additive divergence remains by $m = m_q + m_{res} \rightarrow 0$: [see talk by Y. Aoki on Thursday]

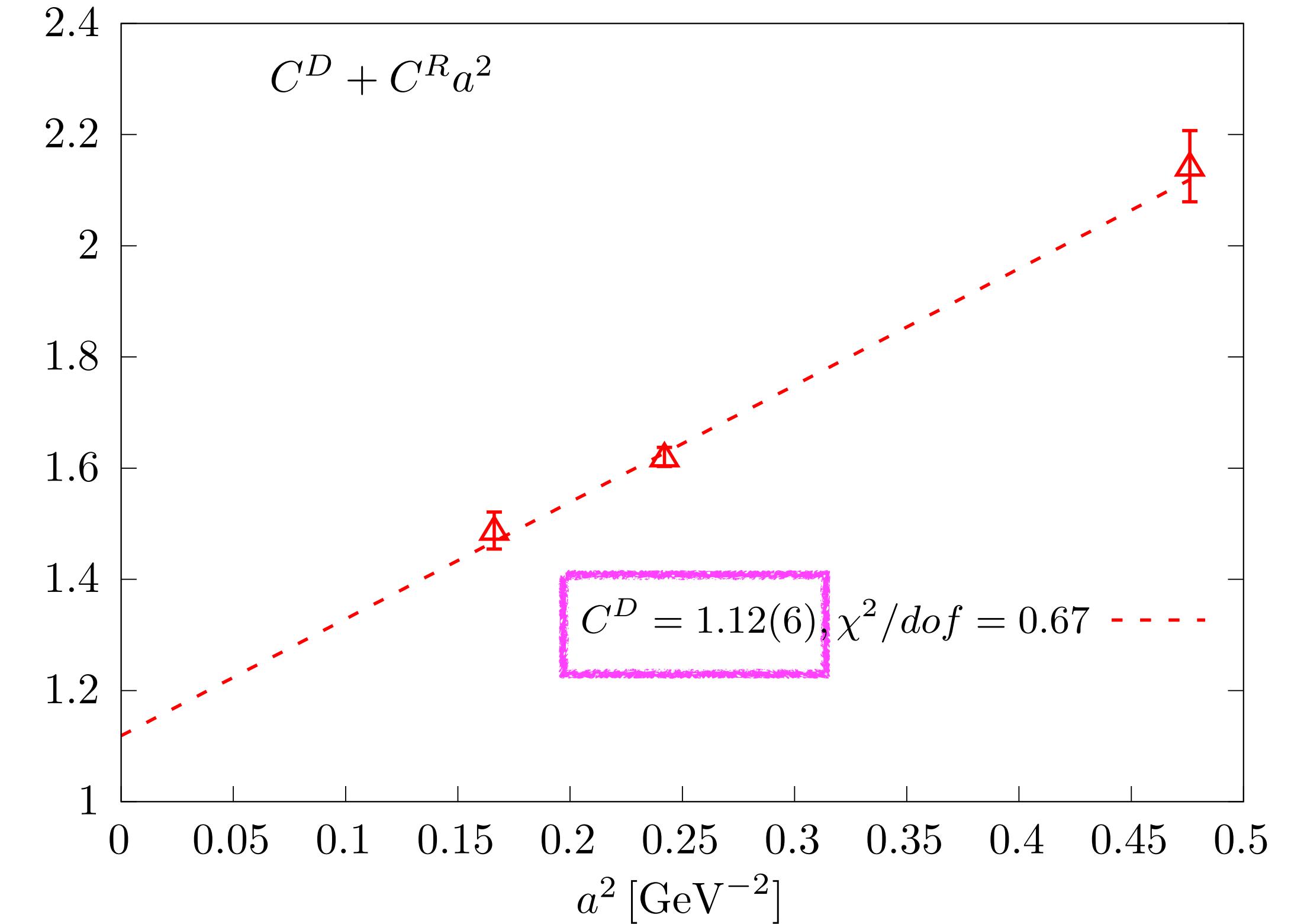
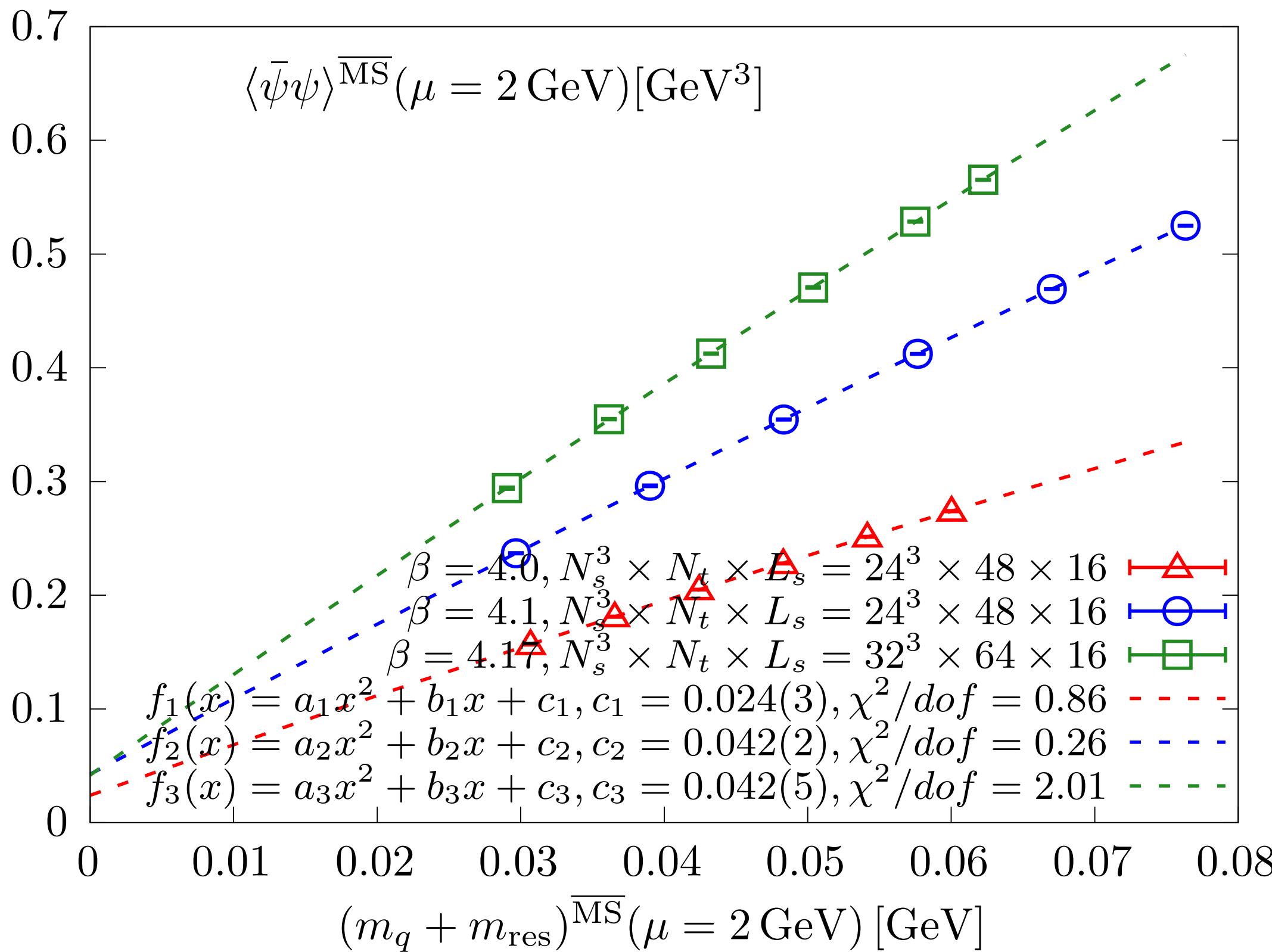
$$\lim_{m \rightarrow 0} \lim_{L \rightarrow 0} \langle \bar{\psi} \psi \rangle|_{DWF} \sim \langle \bar{\psi} \psi \rangle|_{\text{cont.}} + C^D \frac{(x - 1)m_{res}}{a^2}$$

Two ways of subtracting divergences:

- $\langle \bar{\psi} \psi \rangle^{\text{ren}} = Z_m^{-1} [\langle \bar{\psi} \psi \rangle^{\text{T}>0} - \langle \bar{\psi} \psi \rangle^{\text{T}=0}]$
- $\langle \bar{\psi} \psi \rangle^{\text{ren}} = Z_m^{-1} \left[\langle \bar{\psi} \psi \rangle - C^D \frac{m_q + xm_{res}}{a^2} \right]$, If we know C^D and x

Zero T results

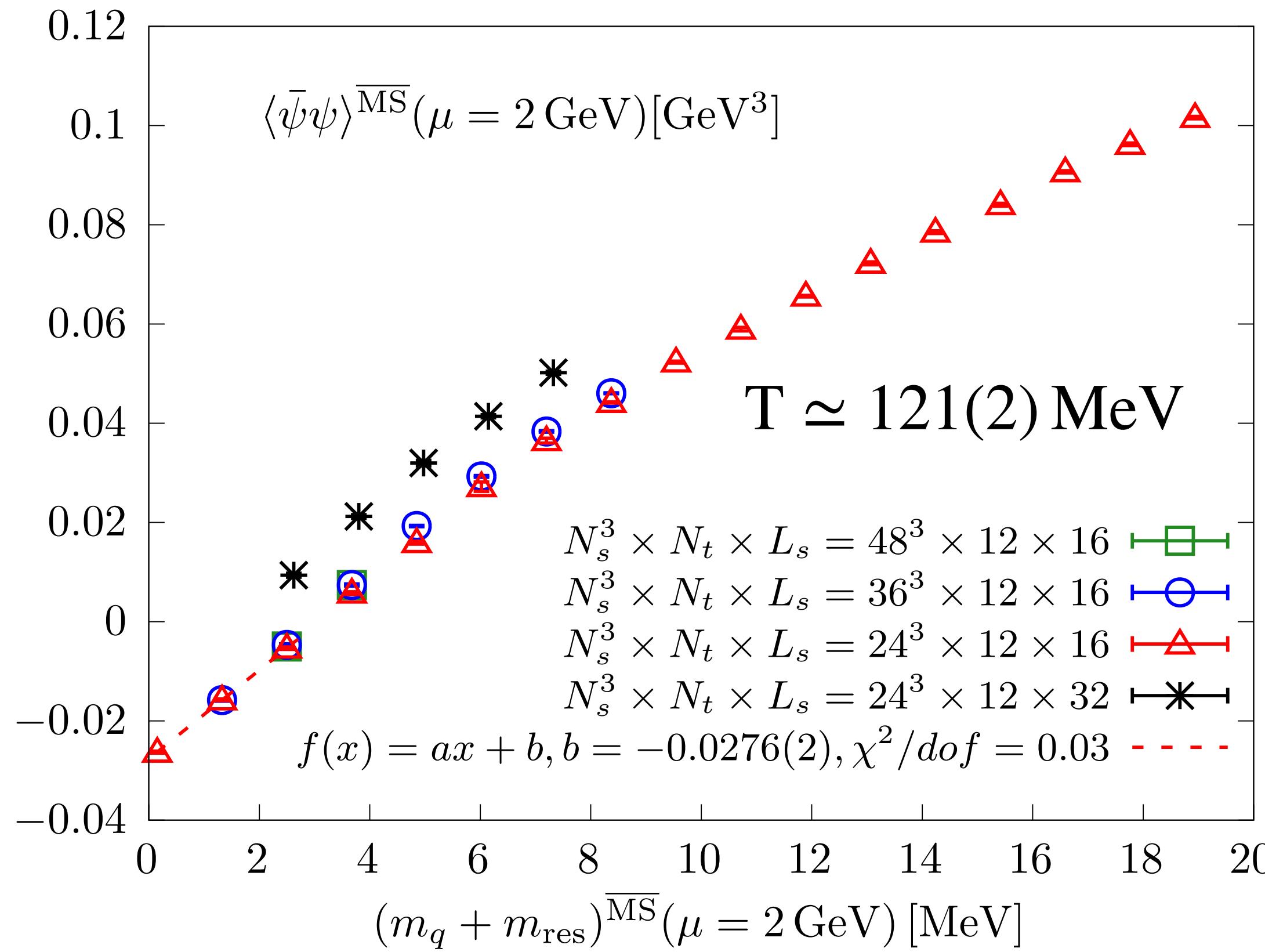
Calculate C^D for subtracting UV divergence term $C^D \frac{m_q + xm_{res}}{a^2}$



$$\begin{aligned} \langle\bar{\psi}\psi\rangle(m_q + m_{\text{res}}) &= \langle\bar{\psi}\psi\rangle(0) + C^D \frac{m_q + xm_{\text{res}}}{a^2} + C^R(m_q + m_{\text{res}}) + A(m_q + m_{\text{res}})^2 \\ &= \langle\bar{\psi}\psi\rangle(0) + (C^D + C^R a^2) \frac{m_q + m_{\text{res}}}{a^2} + C^D \frac{(x - 1)m_{\text{res}}}{a^2} + A(m_q + m_{\text{res}})^2 \end{aligned}$$

Finite T results

Calculate x for subtracting UV divergence term $C^D \frac{m_q + xm_{res}}{a^2}$



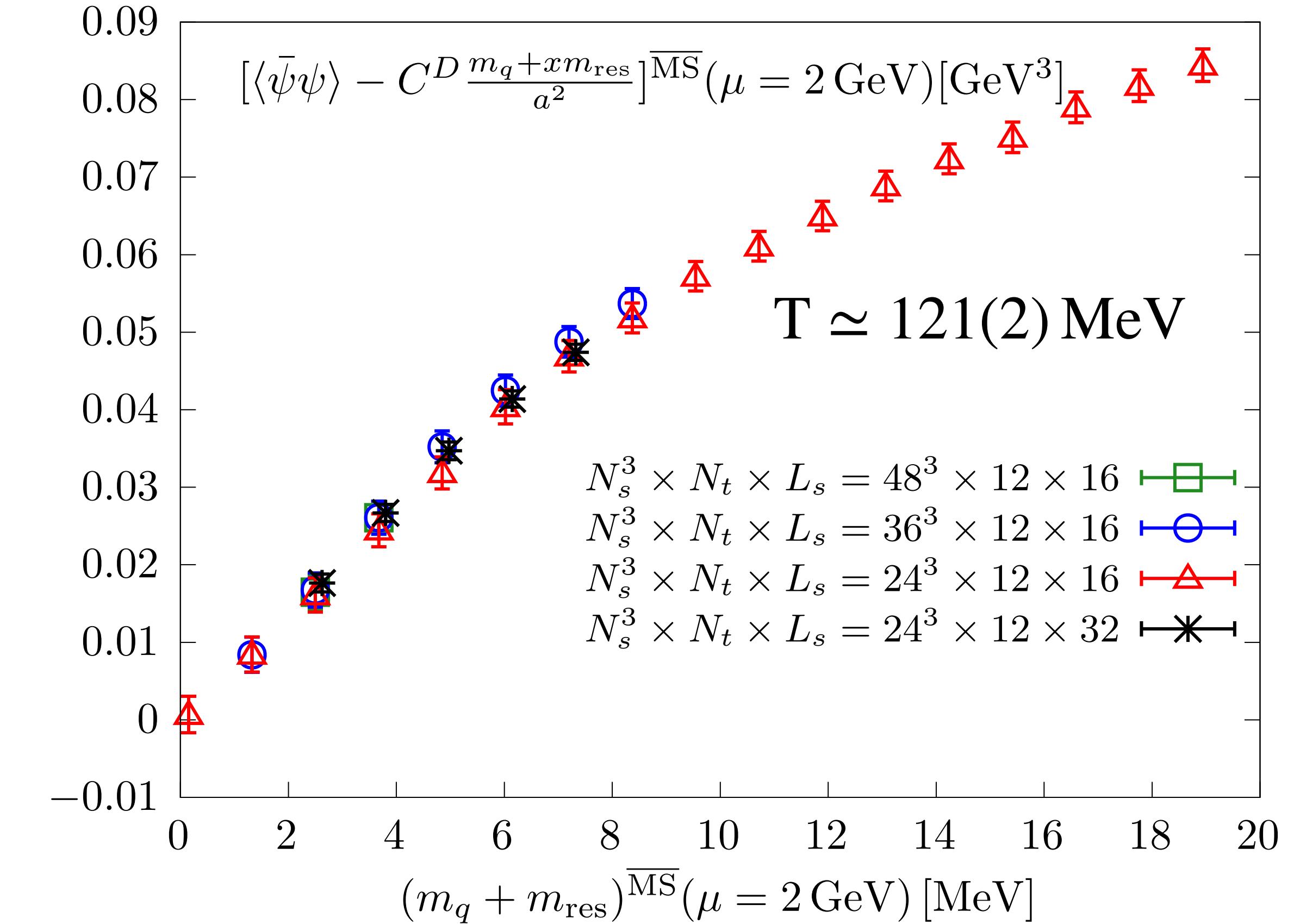
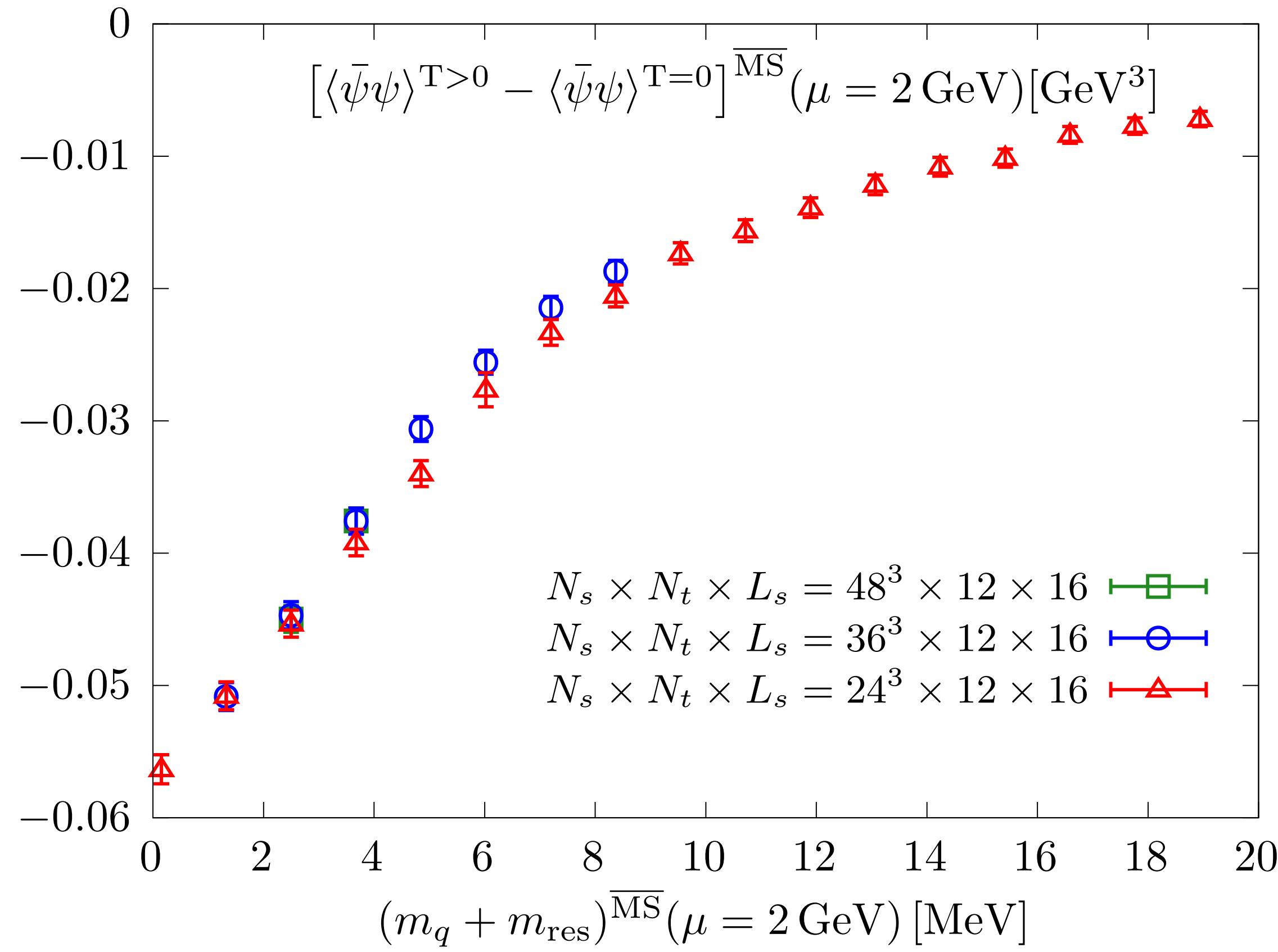
$$\lim_{(m_q + m_{\text{res}}) \rightarrow 0} \langle \bar{\psi} \psi \rangle|_{DWF} \sim \langle \bar{\psi} \psi \rangle|_{\text{cont.}} + C^D \frac{(x - 1)m_{\text{res}}}{a^2}$$

$x = -0.6(1)$ for $T > T_c$

$$T_c = 98^{+3}_{-6} \text{ MeV}$$

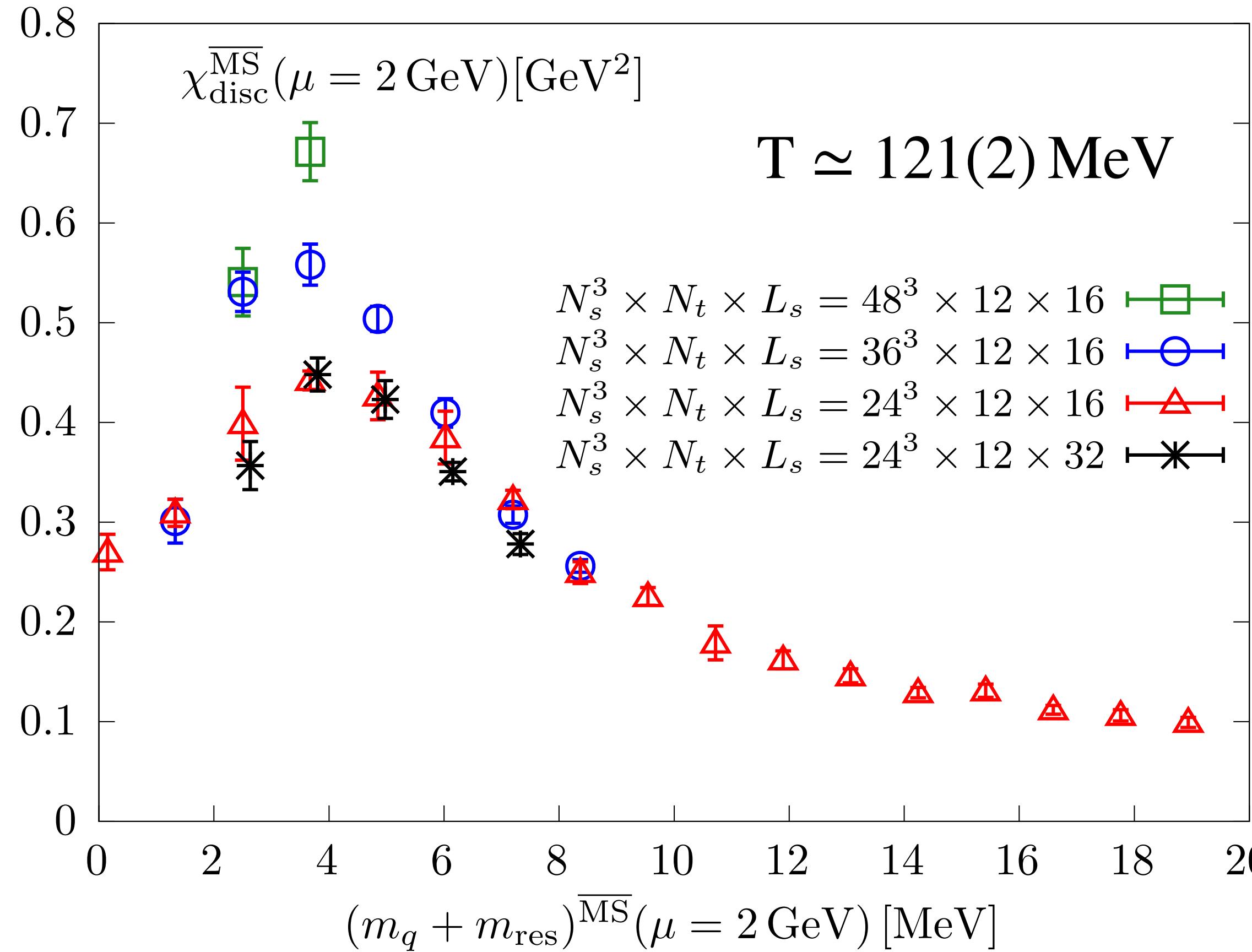
[S. Sharma et al. PRD 2022]

Renormalized chiral condensate



- Subtracted chiral condensate vanishes in the chiral limit
- Finite volume effect is visible at low T
- Subtracted chiral condensate remains the same with same total quark mass for varying Ls

Disconnected chiral susceptibility



- **Large finite volume effect near the transition point, but the change in peak height is not as large as anticipated from a real phase transition**
- **Consistent with the crossover transition**

- **The transition mass point is around 3.6 MeV**

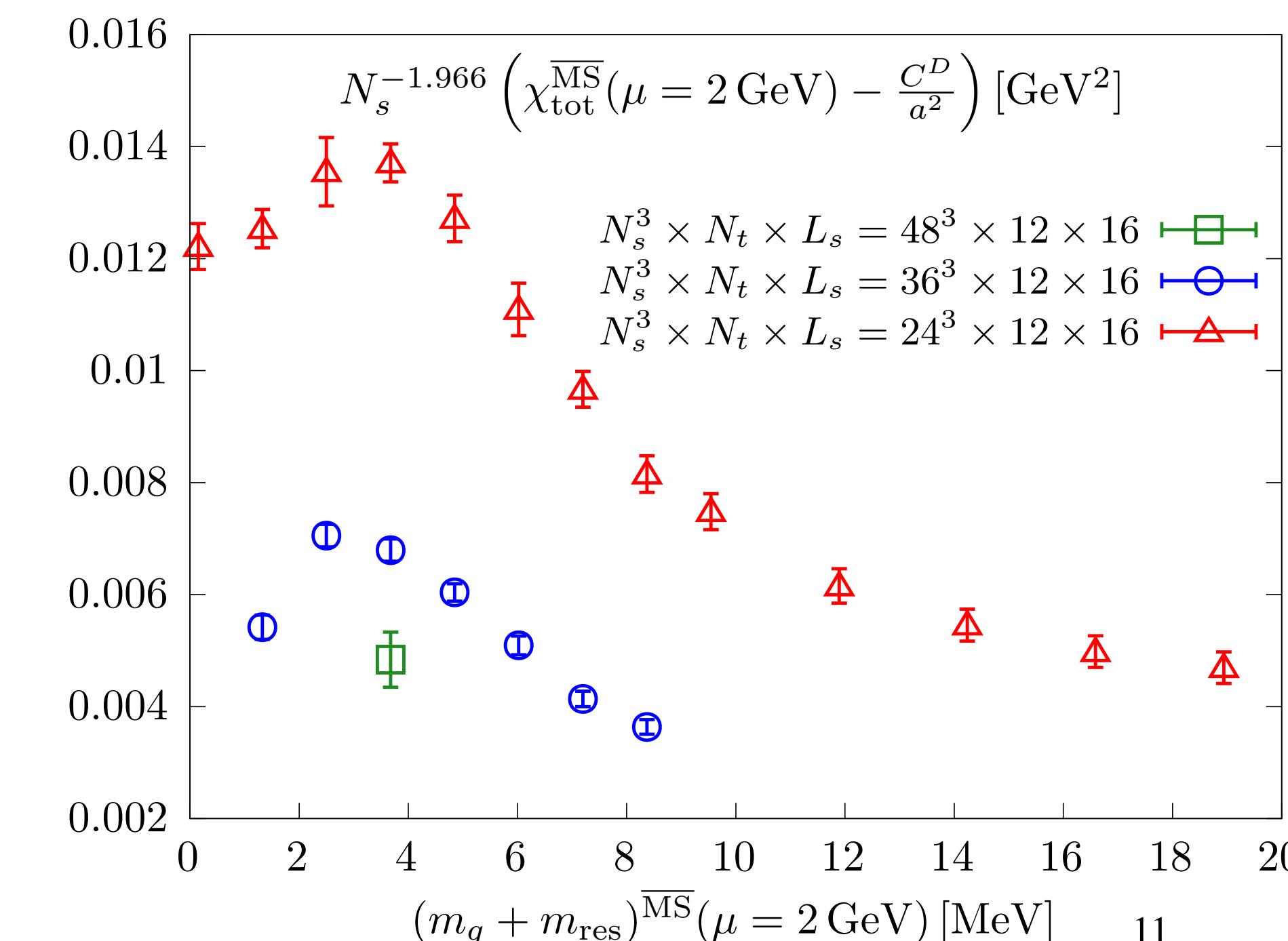
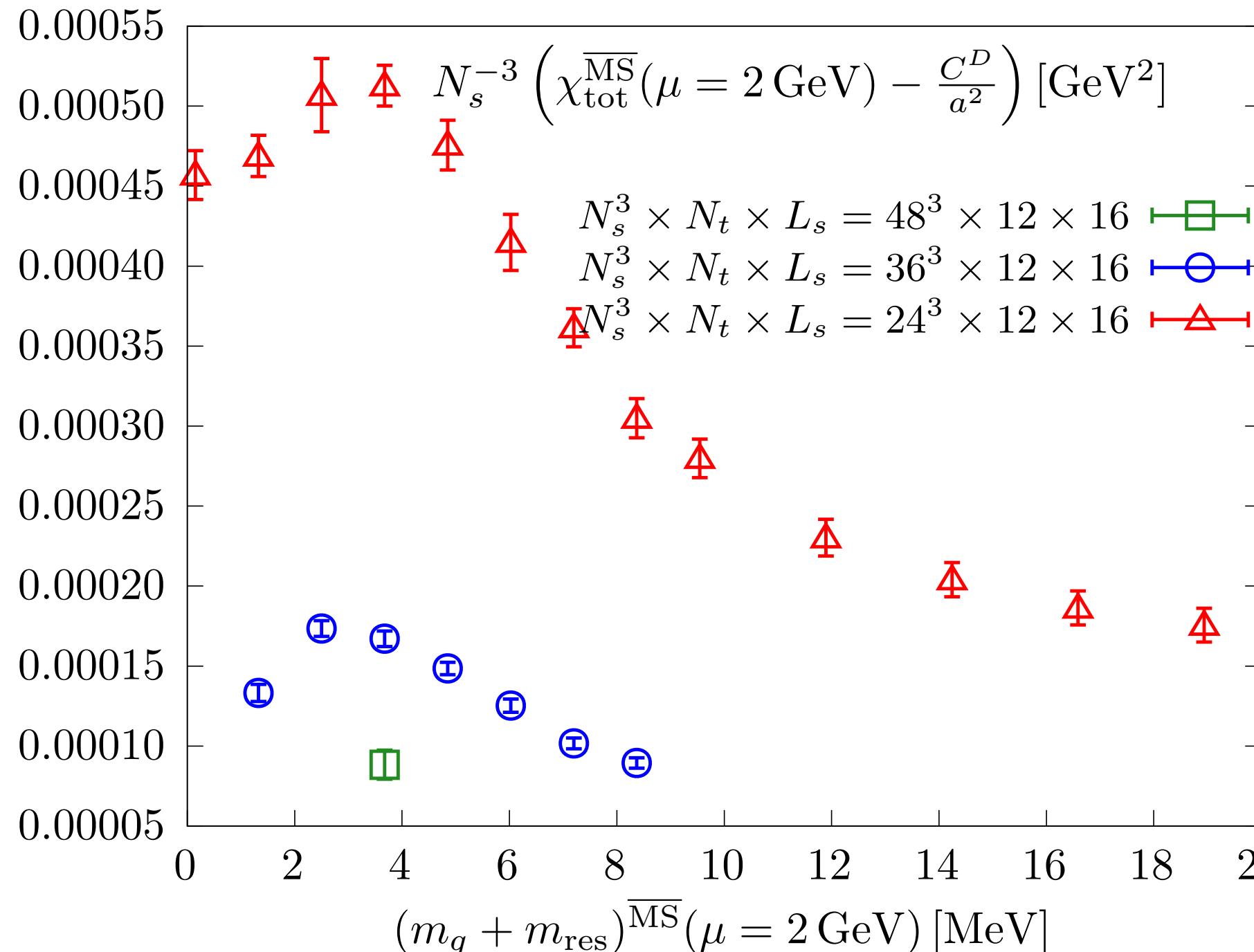
(FLAG Review '21: $m_{ud}^{\overline{MS}}(2 \text{ GeV}) = 3.381(40) \text{ MeV}$)

- χ_{disc} seems to be function of total quark mass

Finite size scaling: susceptibility

Crossover: $\chi_{tot}^{\max}(N_s, N_t)$ independent of V
 1st order PT: $\chi_{tot}^{\max}(N_s, N_t) \propto V$
 Z(2) 2nd order PT: a singular behavior should be observed in $\chi_{tot}^{\max}(N_s, N_t)$ with V
 $((N_s^3 \times N_t)^{\alpha}, \alpha = 1.966$ is the critical exponent)

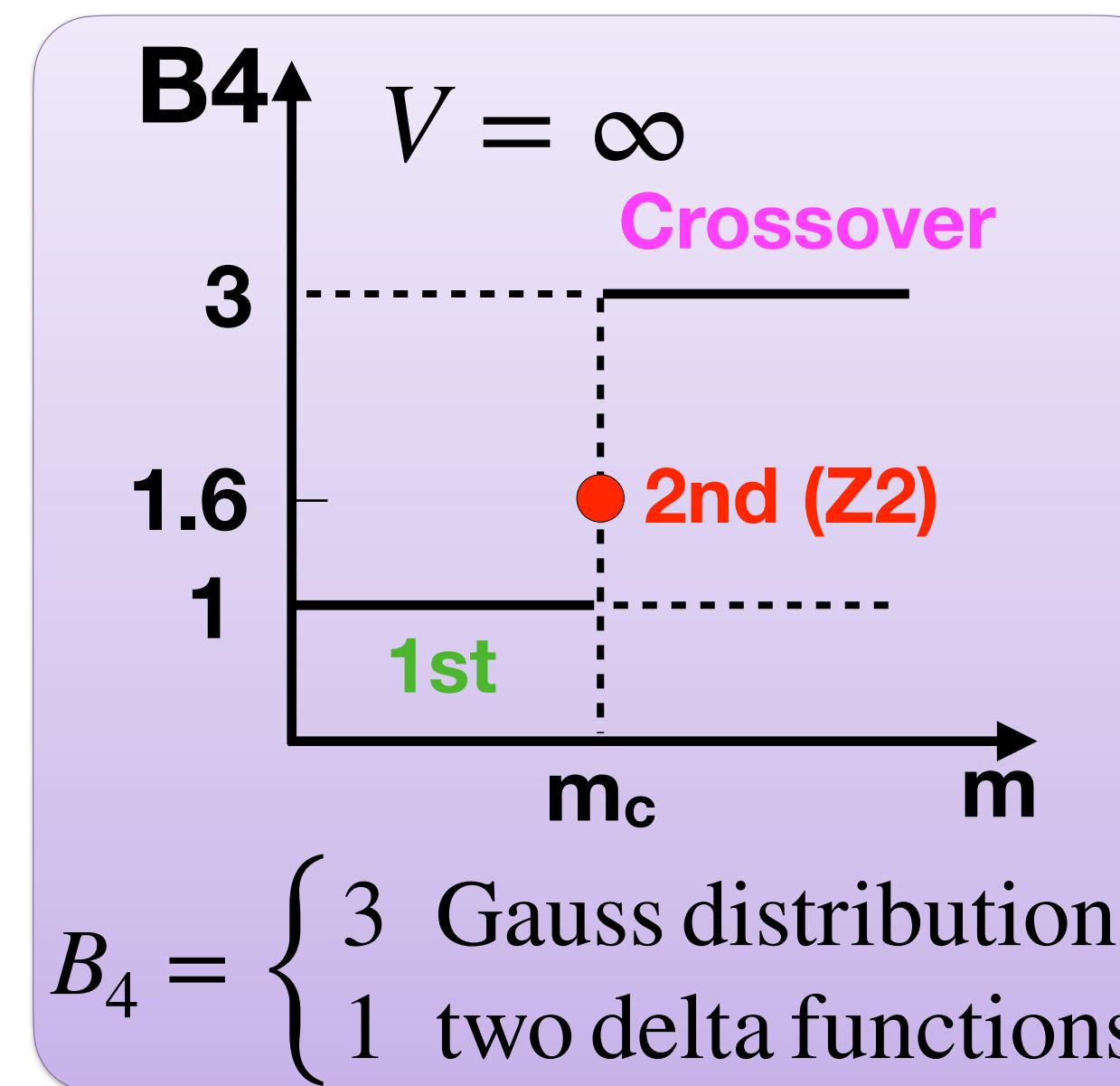
The peak height of χ_{tot} does not scale like a 1st or Z(2) second order PT
finite-size scaled renormalized chiral susceptibility



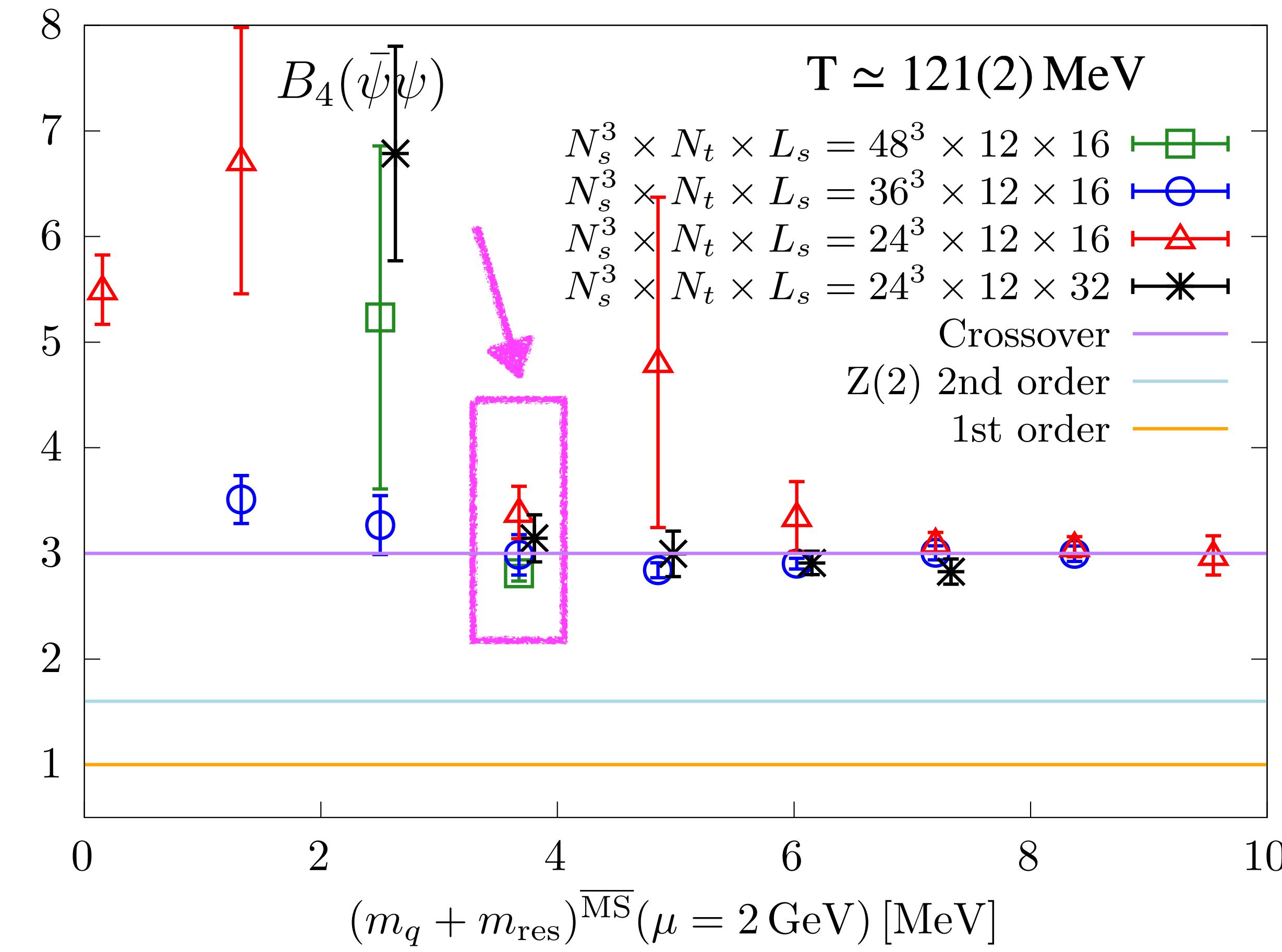
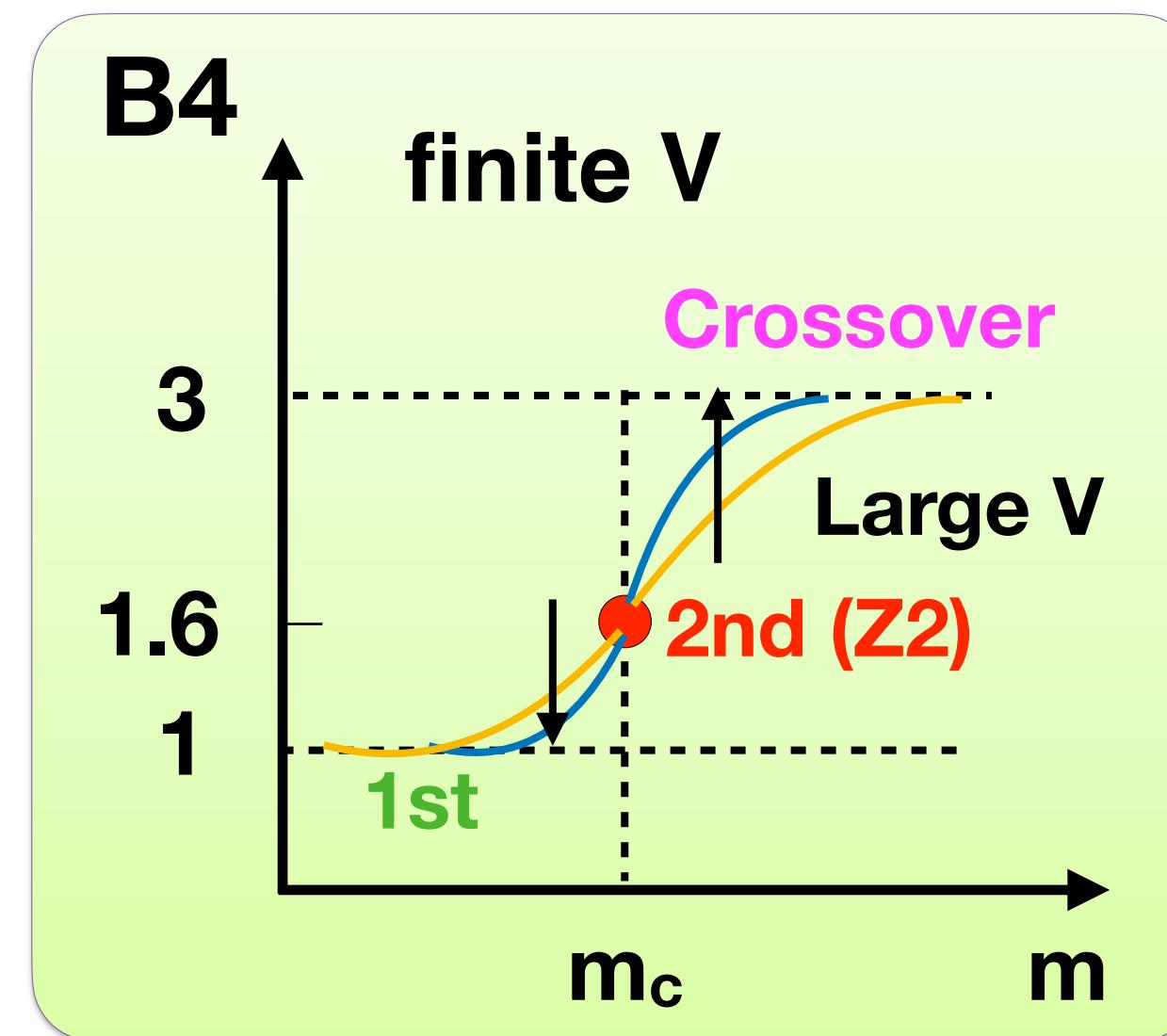
Binder Cumulant

$$B_4(\bar{\psi}\psi) = \frac{\langle (\delta\bar{\psi}\psi)^4 \rangle}{\langle (\delta\bar{\psi}\psi)^2 \rangle^2}, \quad \delta\bar{\psi}\psi = \bar{\psi}\psi - \langle \bar{\psi}\psi \rangle$$

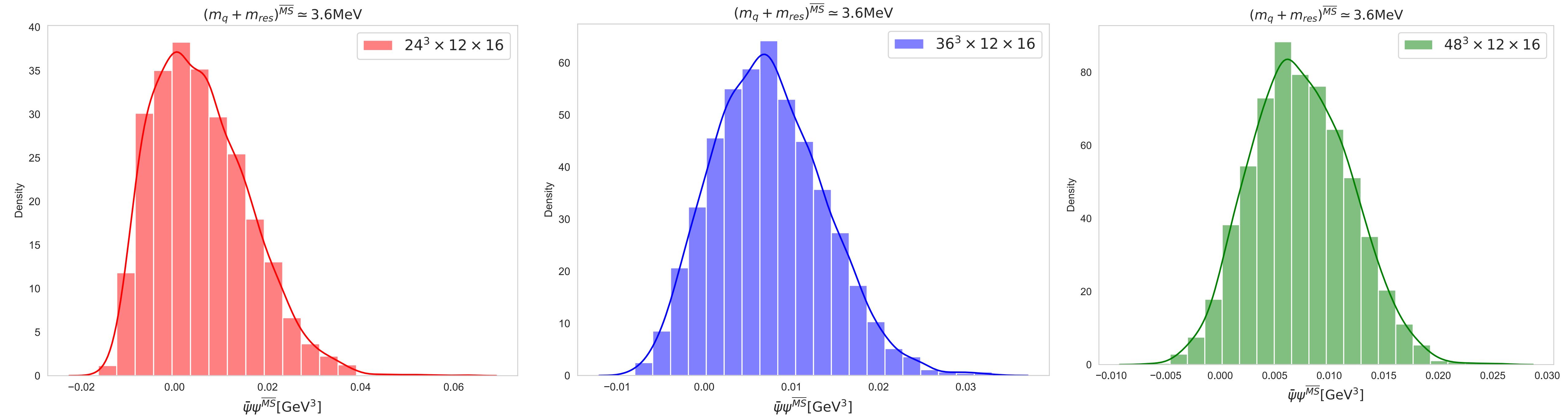
Suggests a crossover transition at $(m_q + m_{\text{res}})^{\overline{\text{MS}}} (2 \text{ GeV}) \simeq 3.6 \text{ MeV}$



A.Kiyohara et al., PRD 104, 114509 (2021)



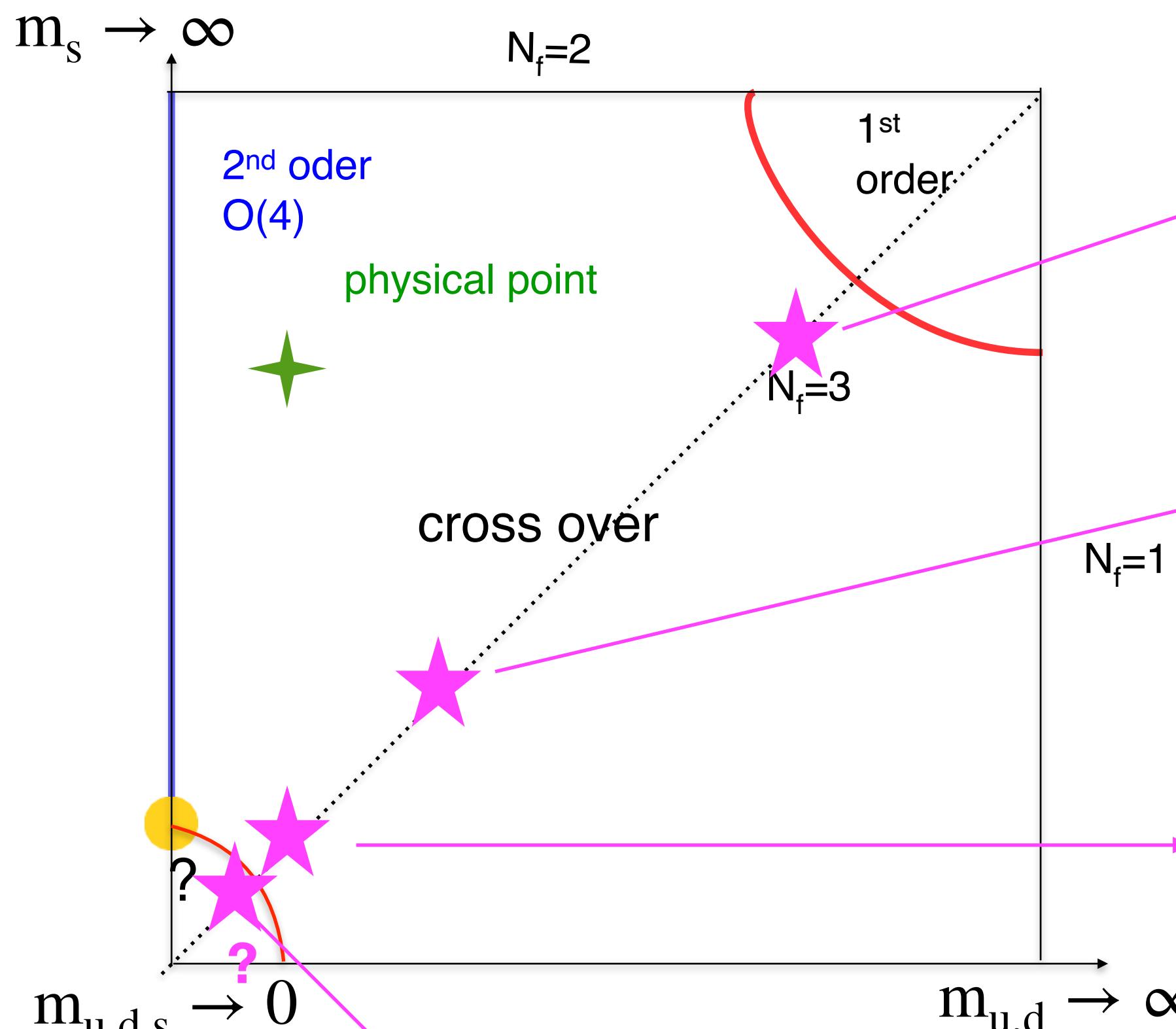
Histogram of chiral condensate near transition point



Behaves like a Gaussian distribution, no evidence of a double peak structure would appear as $V \uparrow$

Summary and outlook

No evidence of a 1st order PT in our explored quark mass range
If a 1st order region exist, the critical mass should be less than 3.6 MeV



- $T \sim 242 \text{ MeV}(N_t = 6), (m_q + m_{\text{res}})_{pc}^{\overline{\text{MS}}} \sim 183 \text{ MeV}$
 $\Leftrightarrow m_\pi^{pc} \sim 997 \text{ MeV}$, crossover transition

Y. Nakamura, Y. Zhang et al., PoS LATTICE2021

- $T \sim 181 \text{ MeV}(N_t = 8), (m_q + m_{\text{res}})_{pc}^{\overline{\text{MS}}} \sim 42 \text{ MeV}$
 $\Leftrightarrow m_\pi^{pc} \sim 476 \text{ MeV}$, crossover transition

Y. Zhang et al., PoS LATTICE2022

- $T \sim 121 \text{ MeV}(N_t = 12), (m_q + m_{\text{res}})_{pc}^{\overline{\text{MS}}} \sim 3.6 \text{ MeV}$
 $\Leftrightarrow m_\pi^{pc} \sim 141 \text{ MeV}$, crossover transition

Y. Zhang et al., PoS LATTICE2023

- $T \sim 104 \text{ MeV}(N_t = 14)$, lighter quark mass simulation is underway

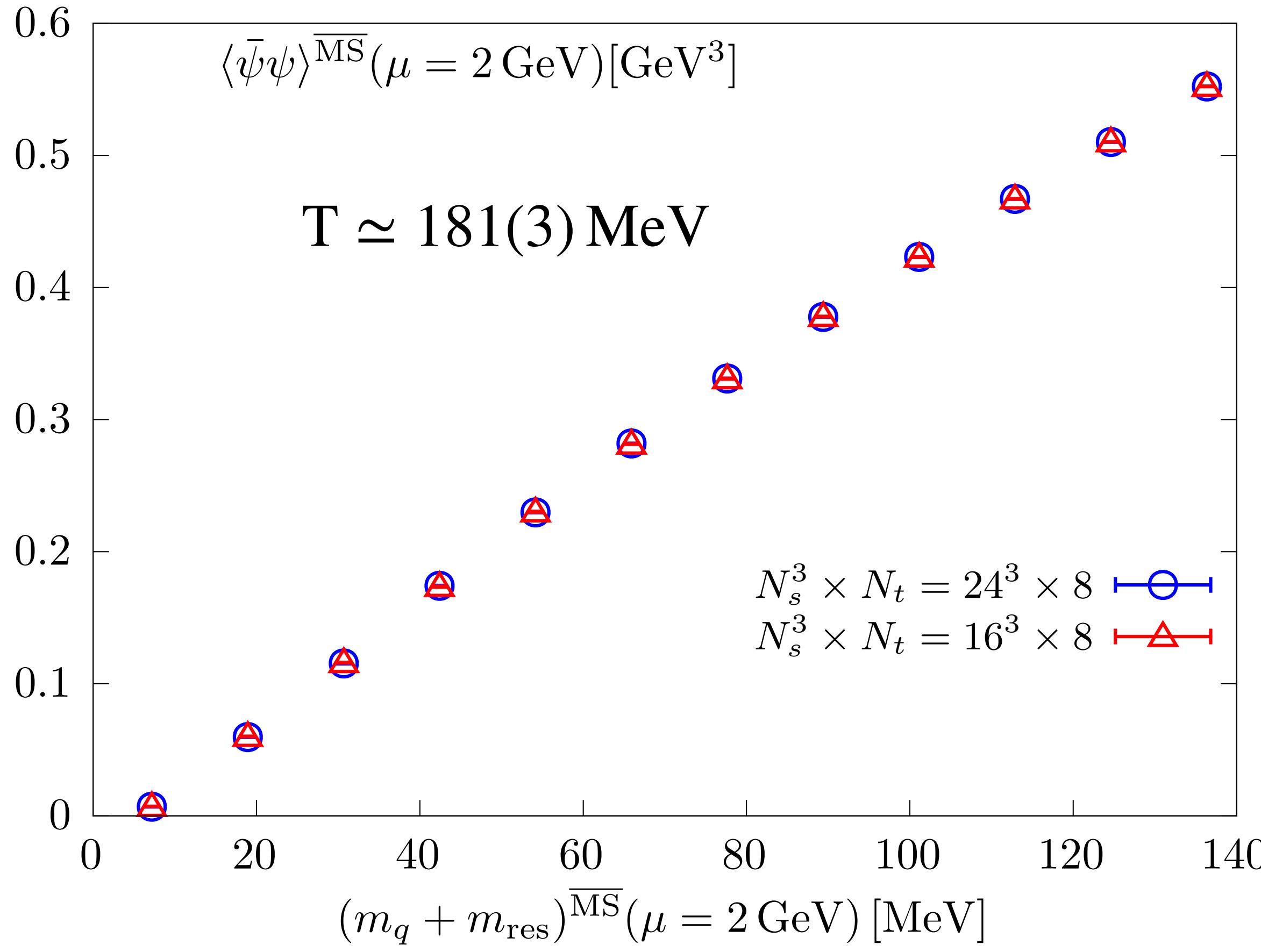
Acknowledgements

- Codes
 - HMC
 - Grid (implementation for A64FX: thanks to the Regensburg group)
 - Measurements
 - Bridge++
 - Hadrons / Grid
- Computers
 - Supercomputer Fugaku provided by the RIKEN Center for Computational Science through HPCI project #hp210032 and Usability Research ra000001.
 - Wisteria/BDEC-01 Oddysey at Univ. Tokyo/JCAHPC through HPCI project #hp220108 Ito supercomputer at Kyushu University through HPCI project #hp190124 and hp200050
 - Hokusai BigWaterfall at RIKEN
- Grants
 - JSPS Kakenhi (20H01907)

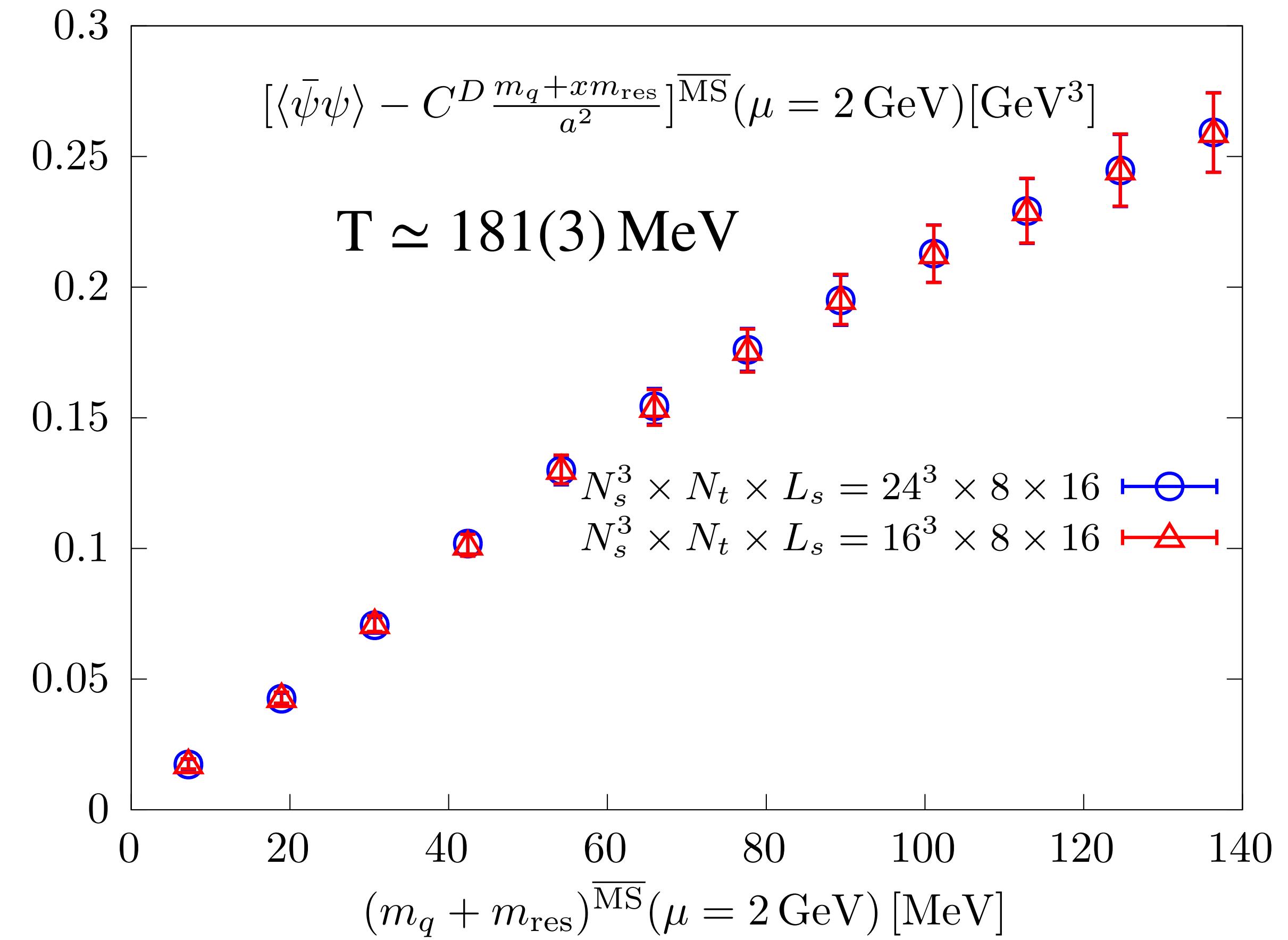
Backup slide

Renormalized chiral condensate

Multiplicatively renormalized chiral condensate

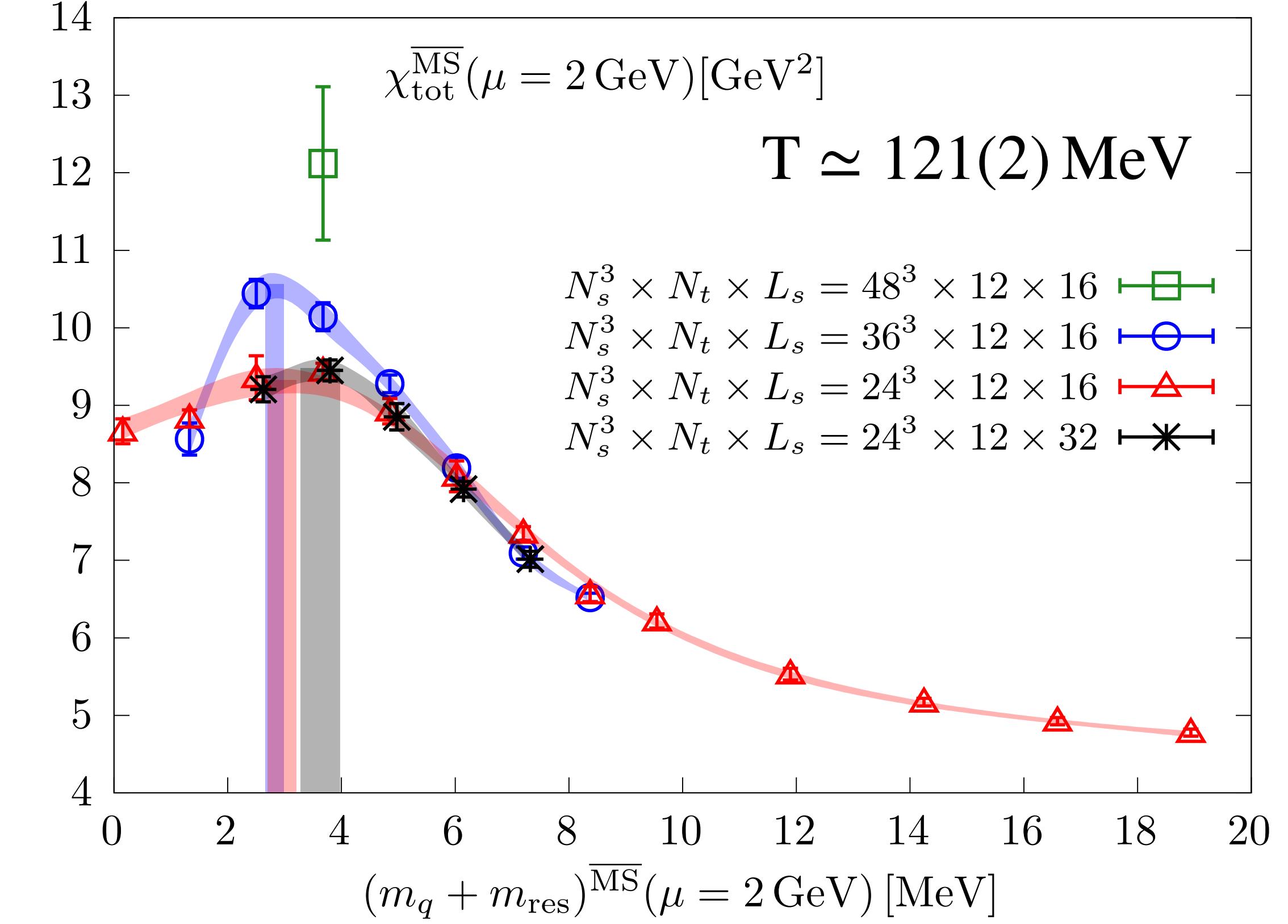
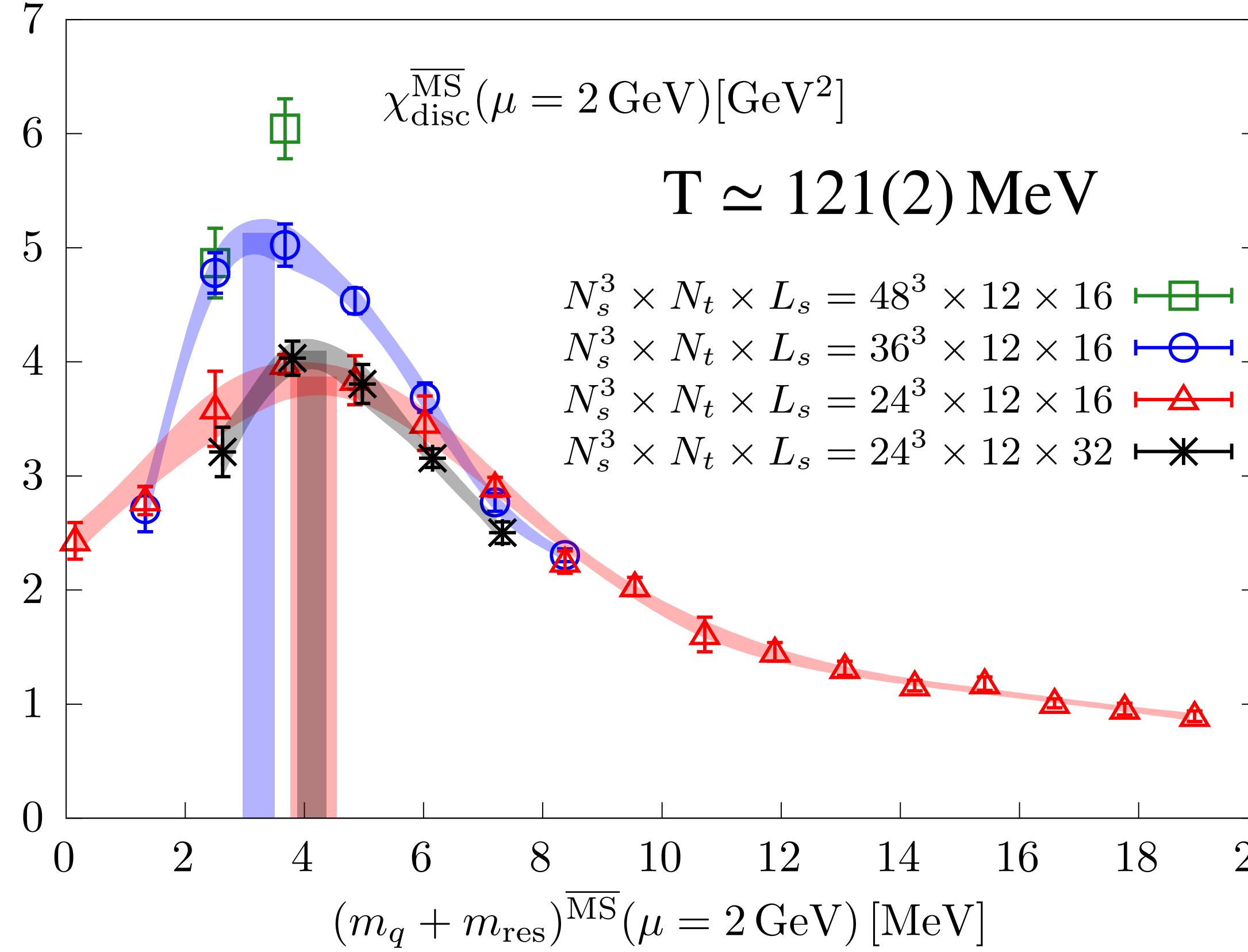


Additive and multiplicatively renormalized chiral condensate

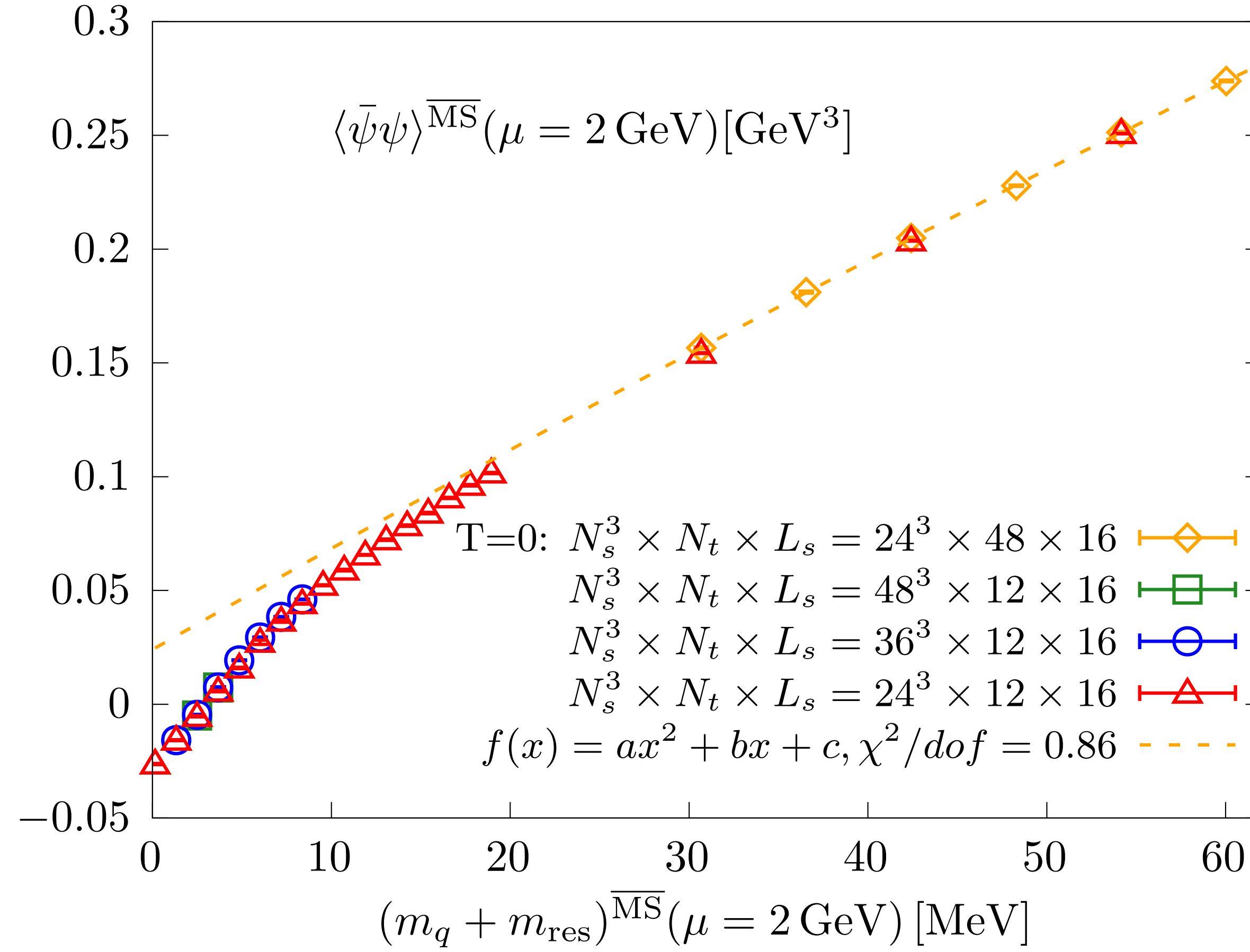


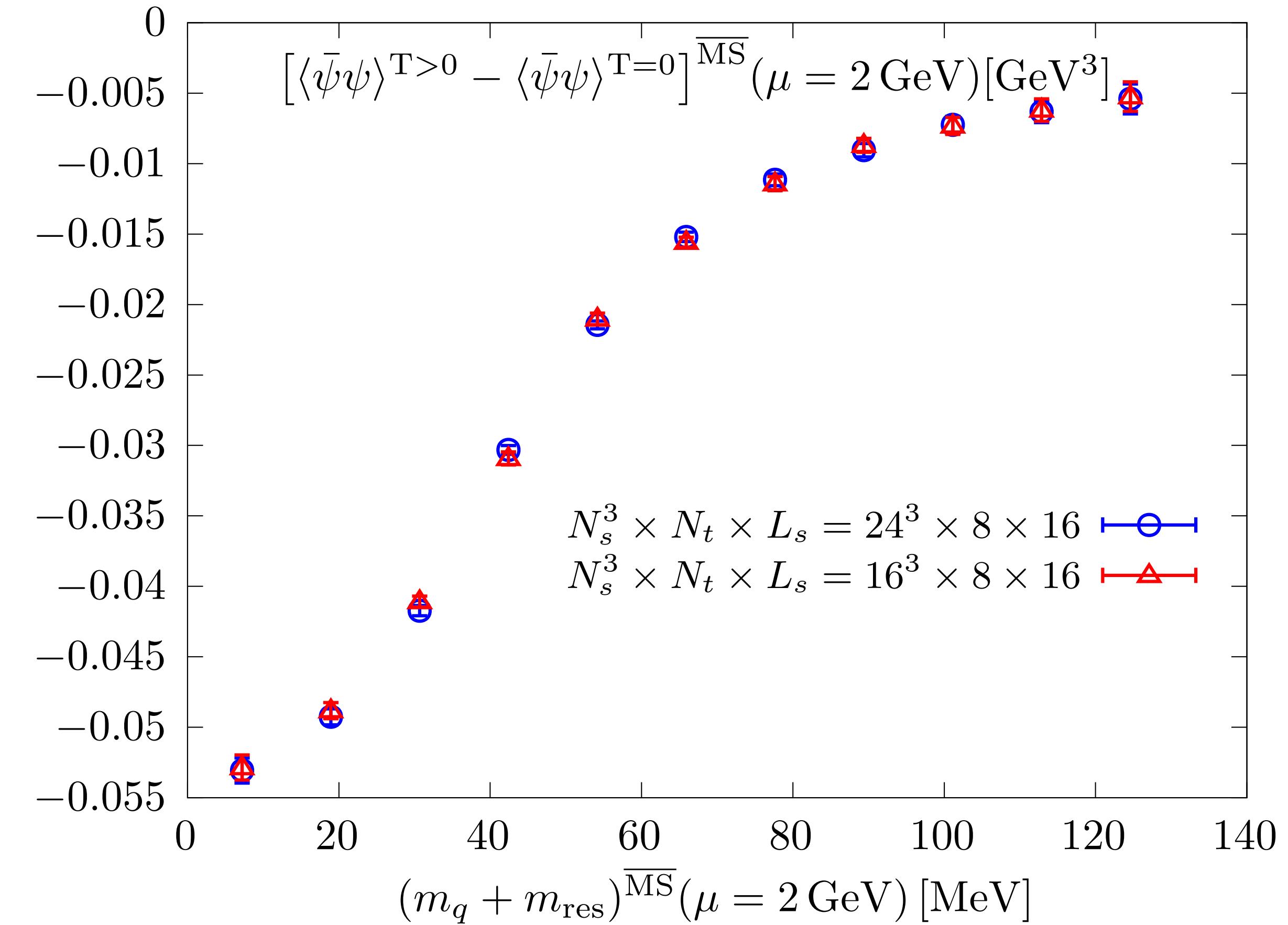
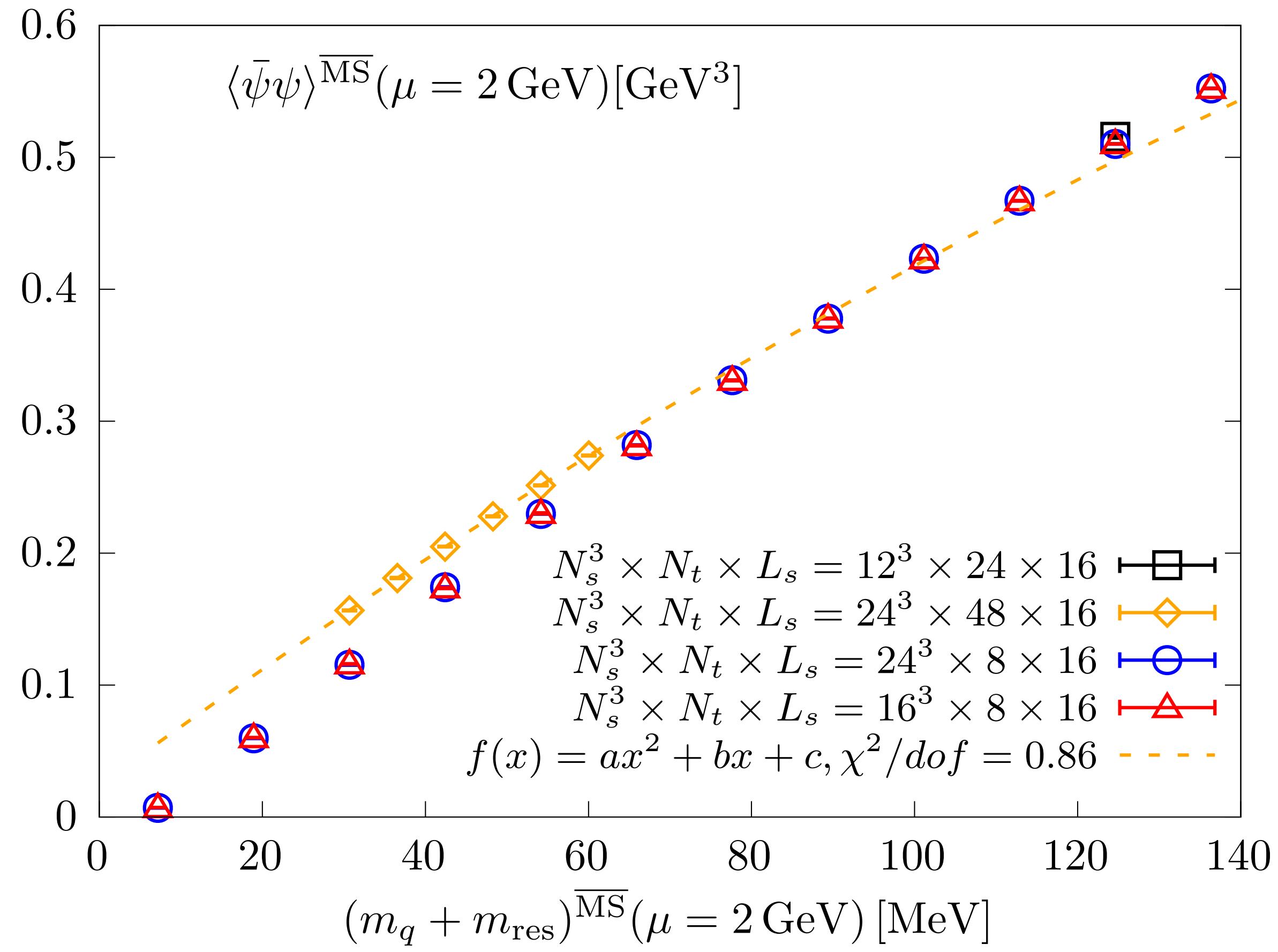
Subtracted chiral condensate vanishes in the chiral limit as expected since $T > T_c$

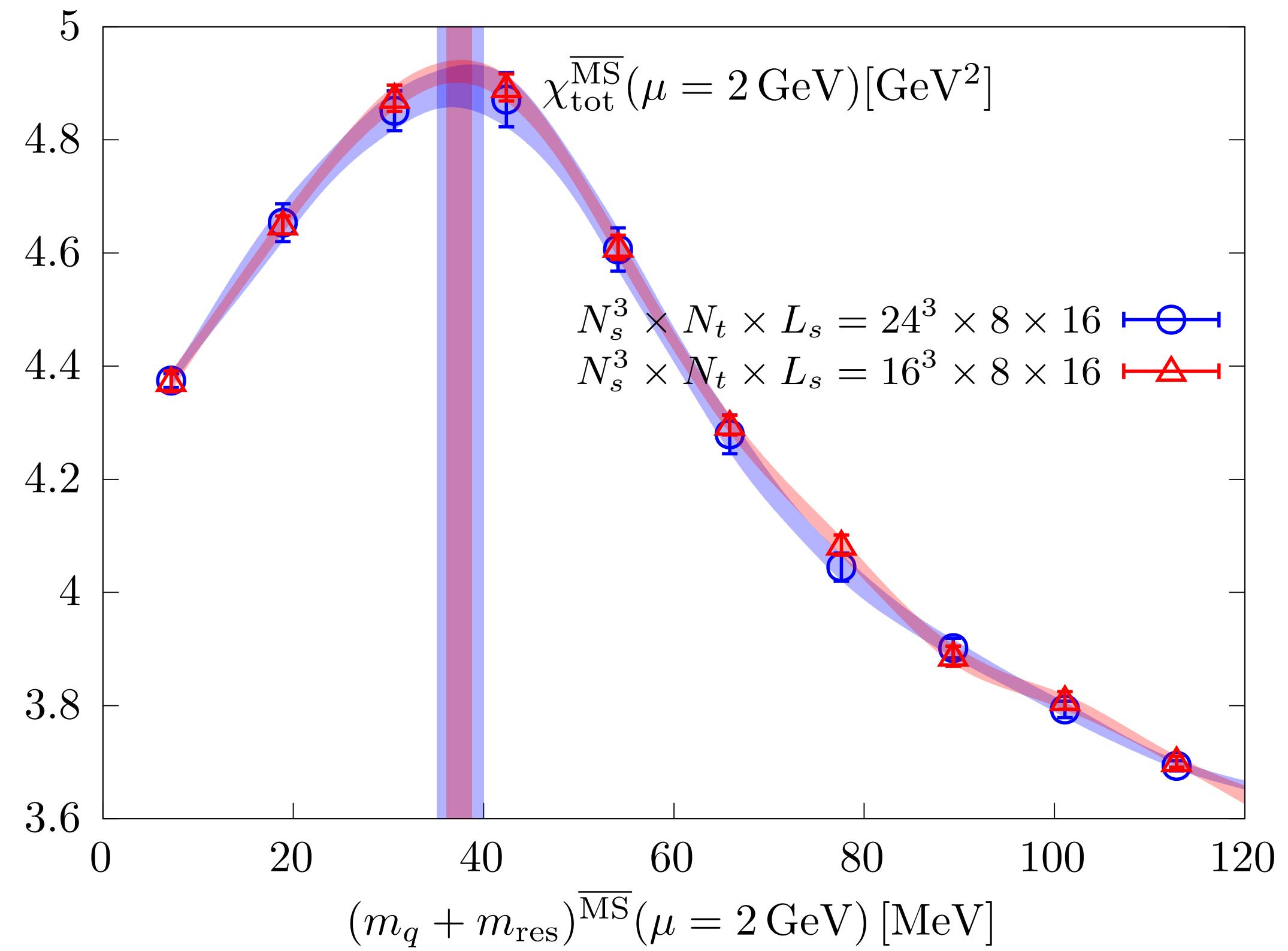
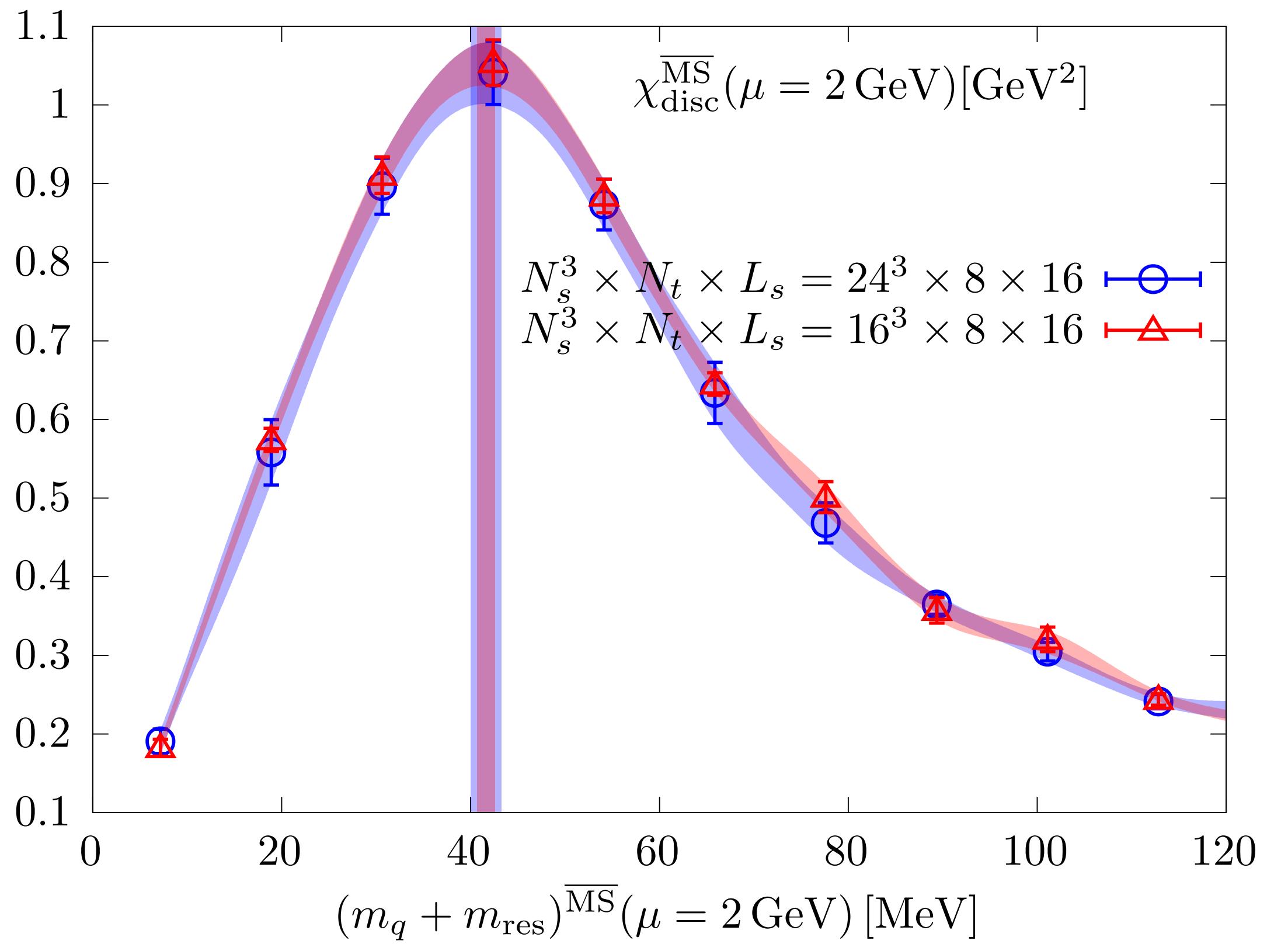
Chiral susceptibility

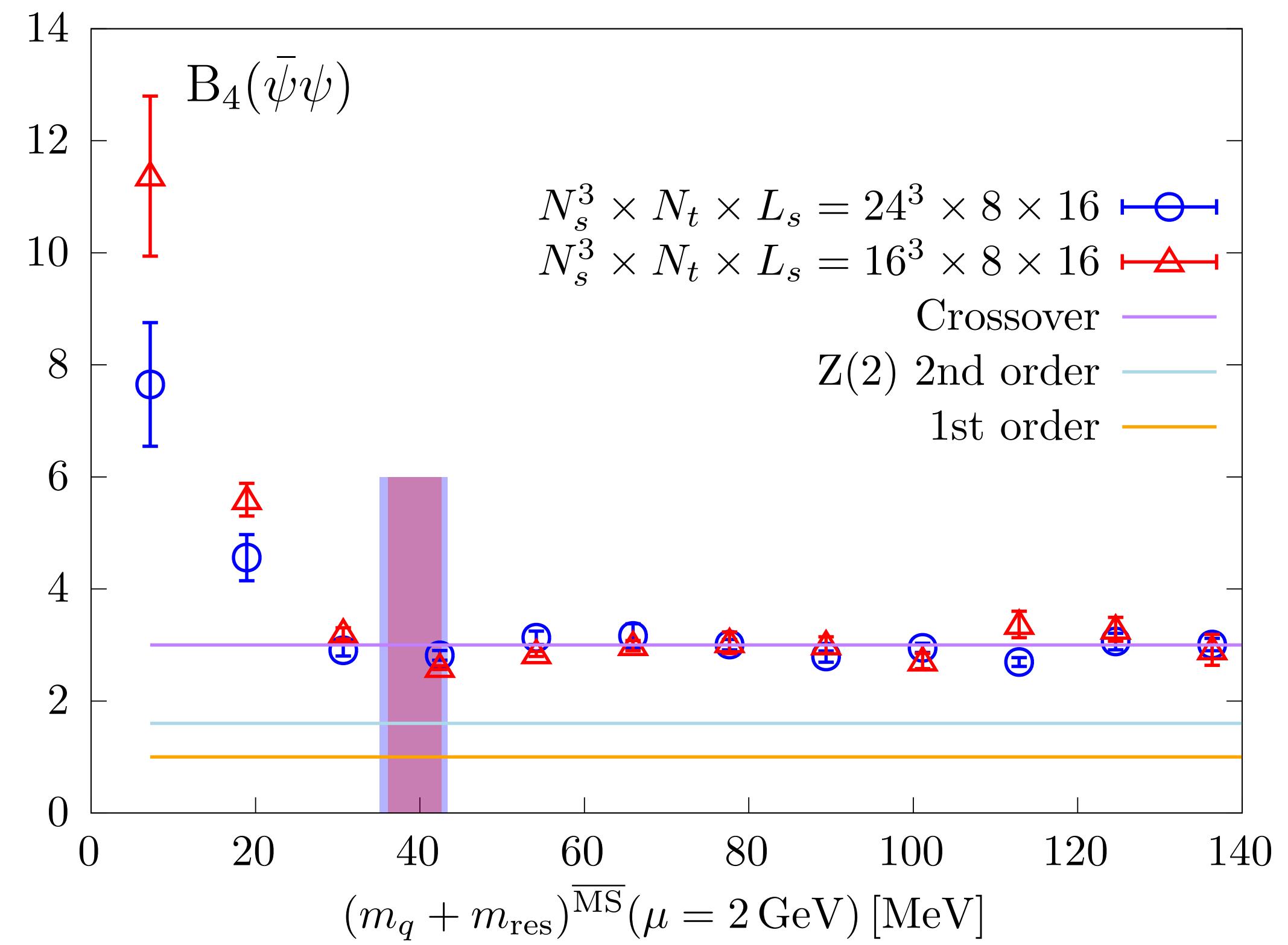
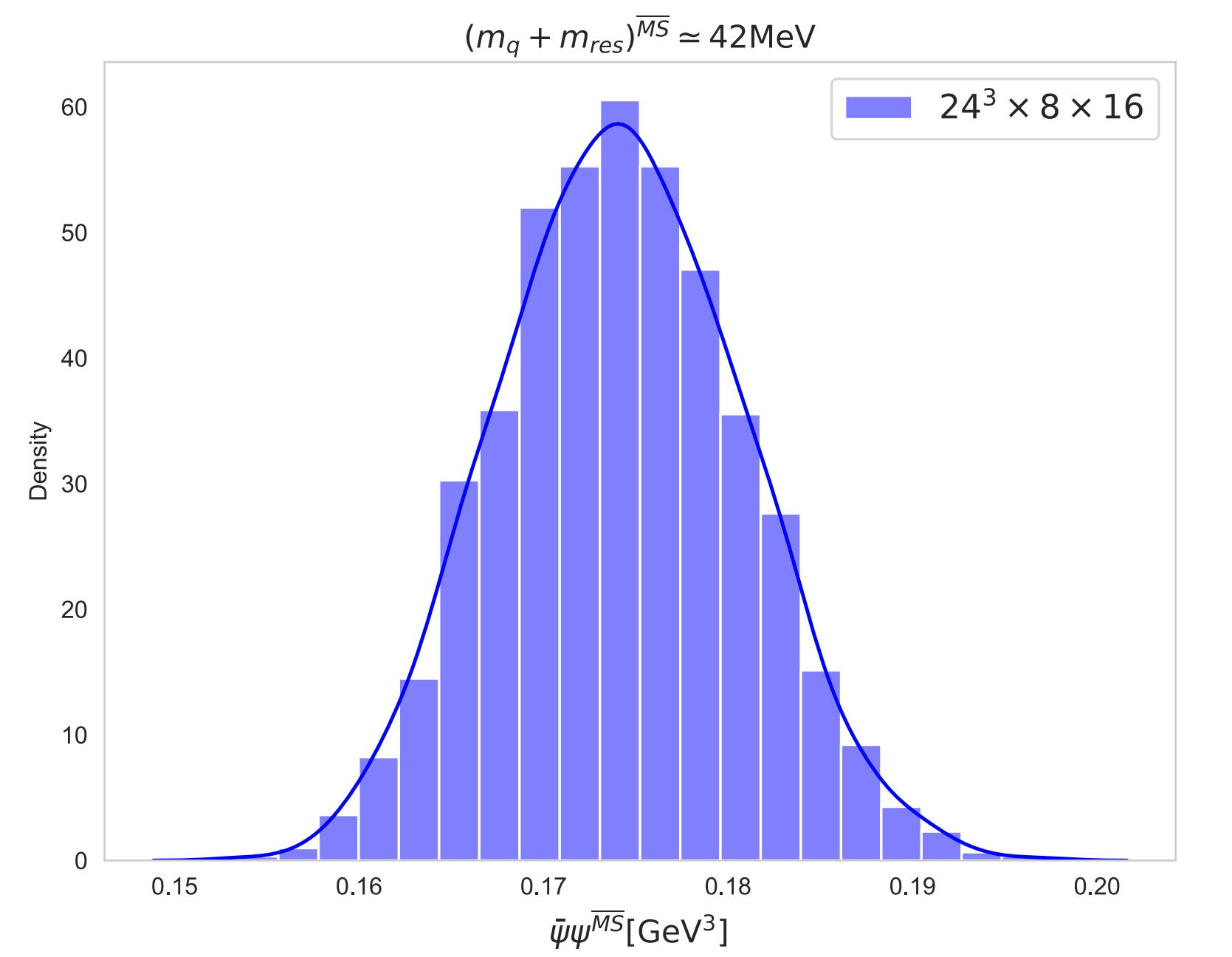
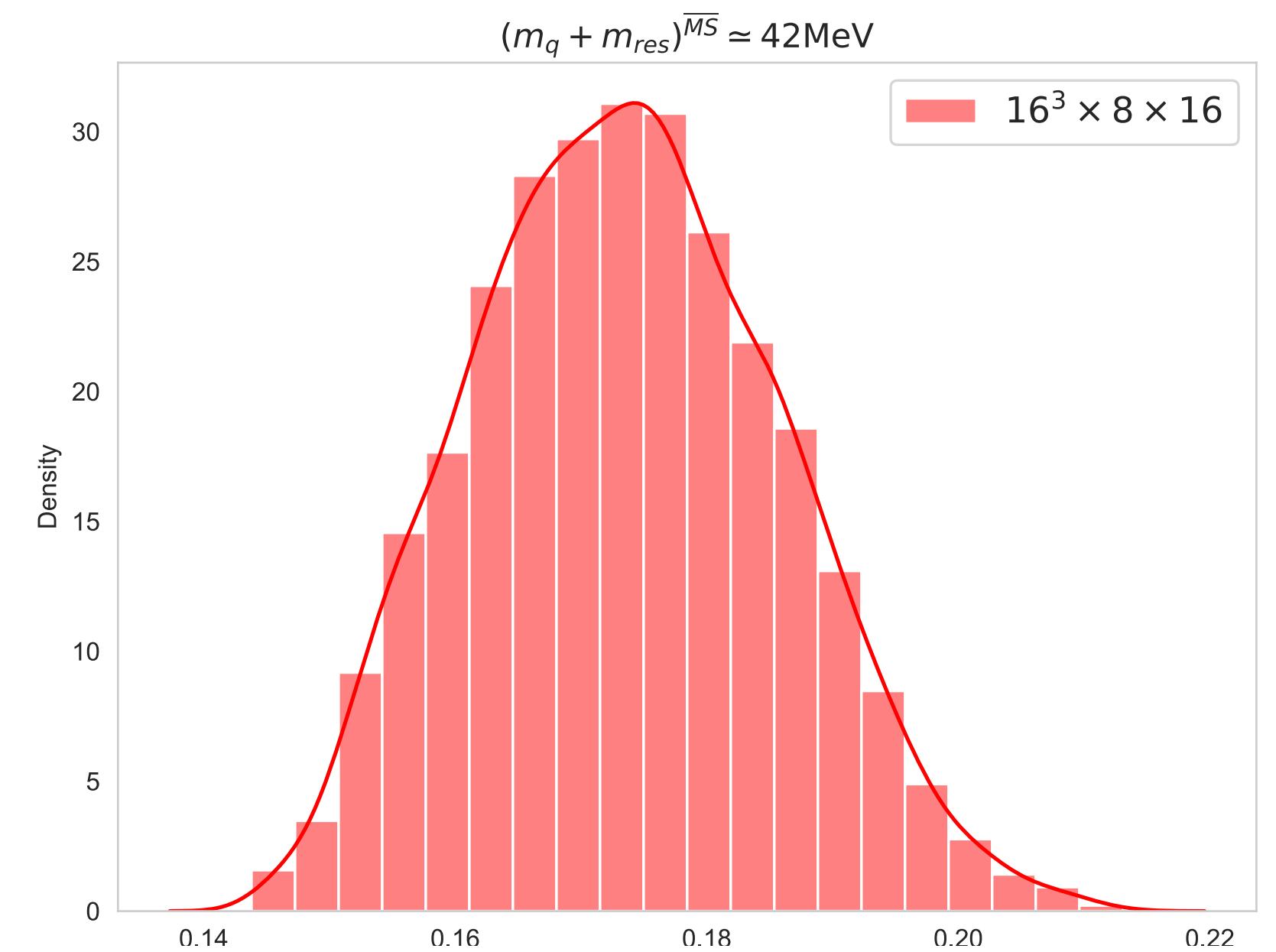


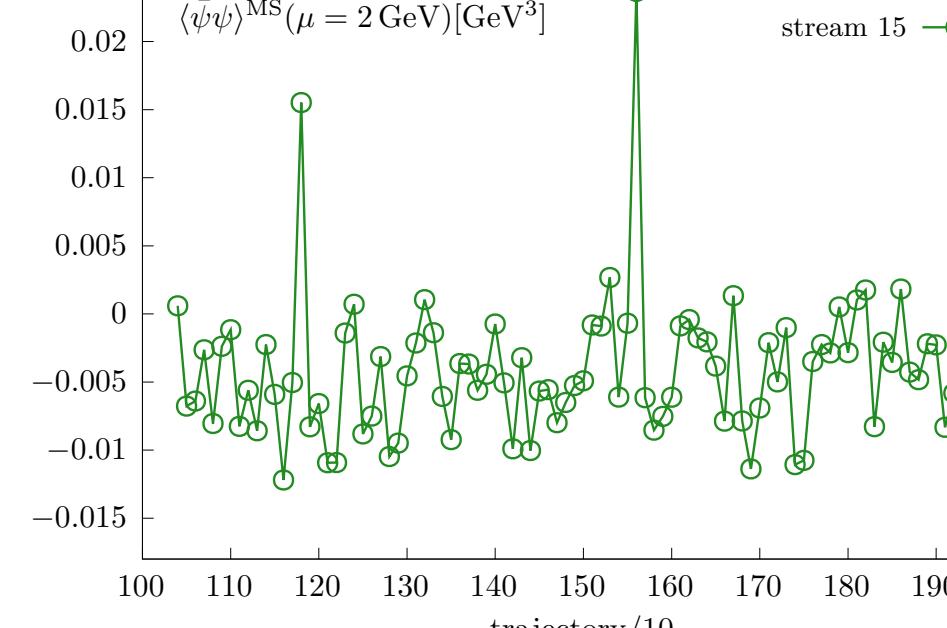
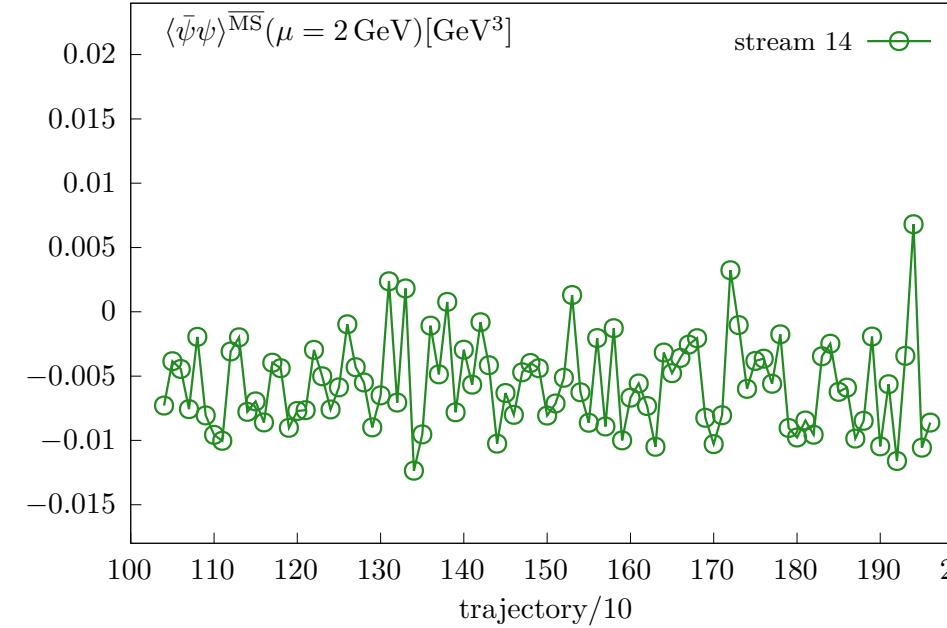
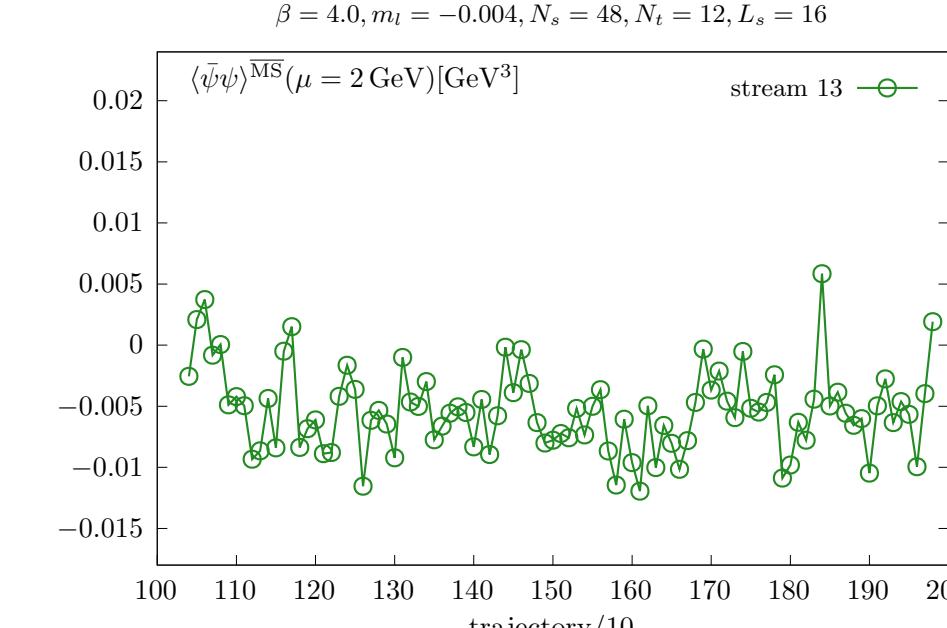
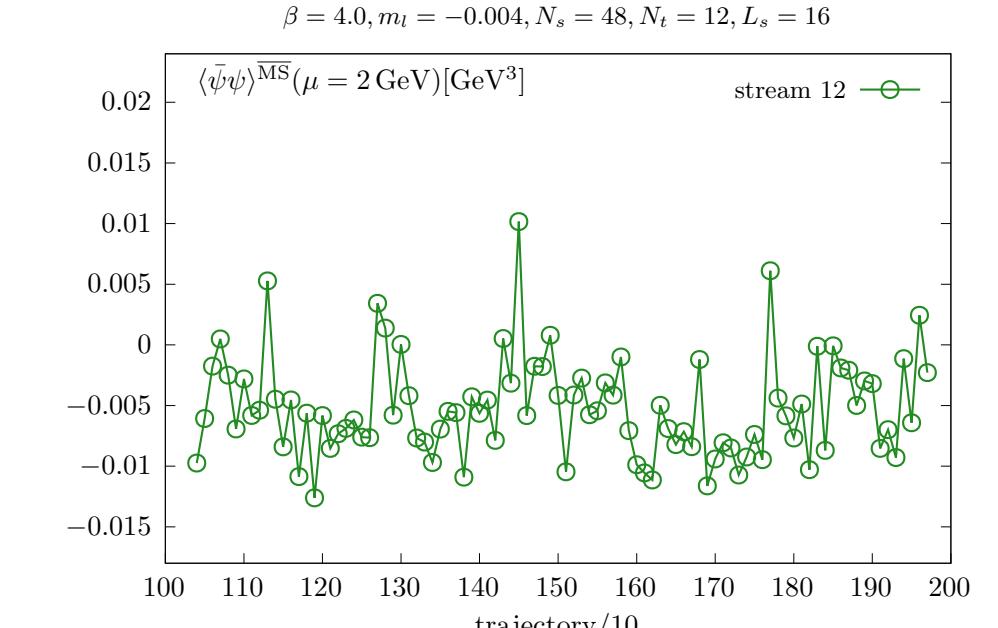
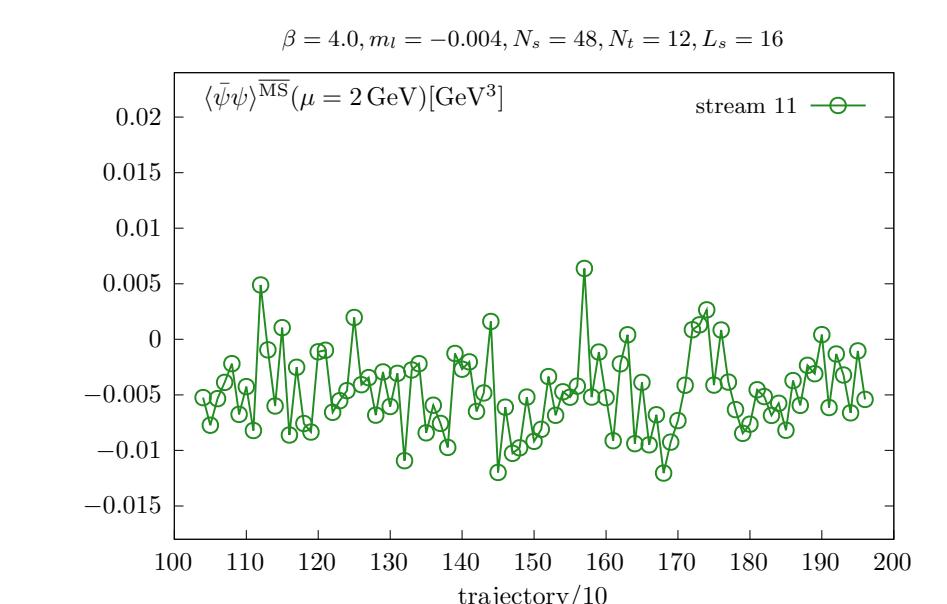
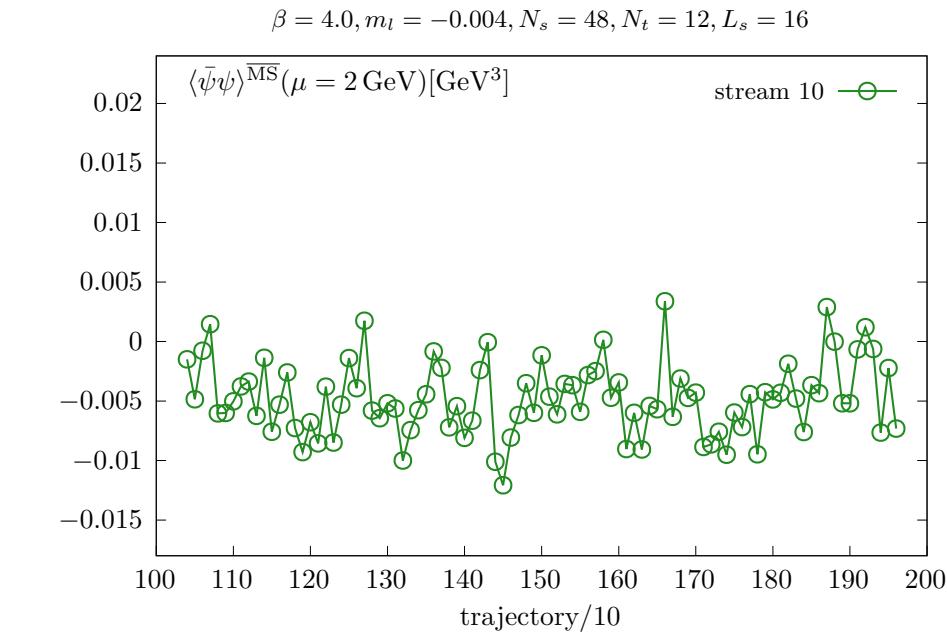
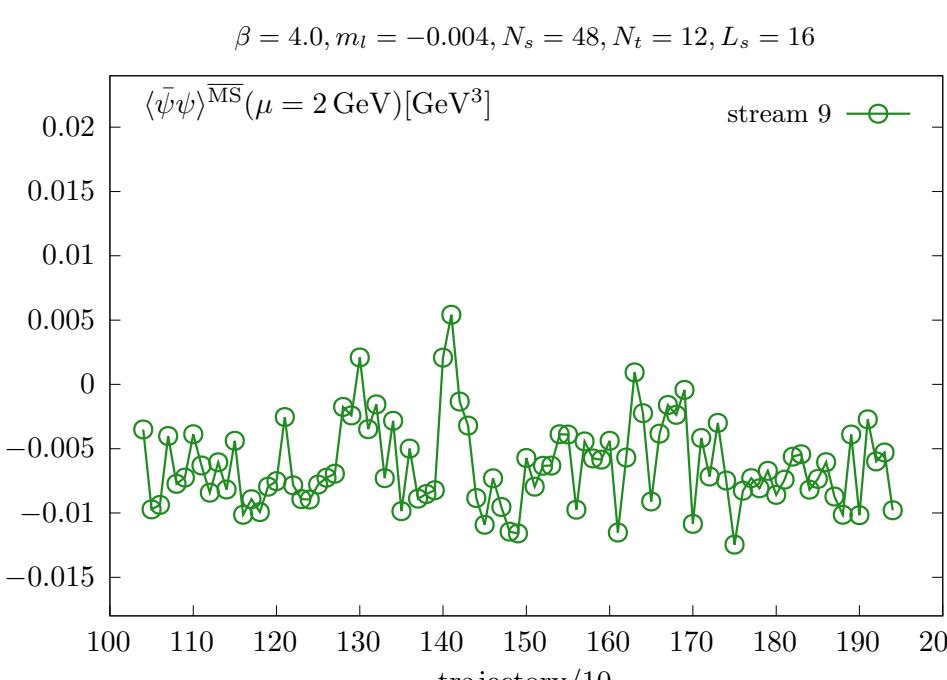
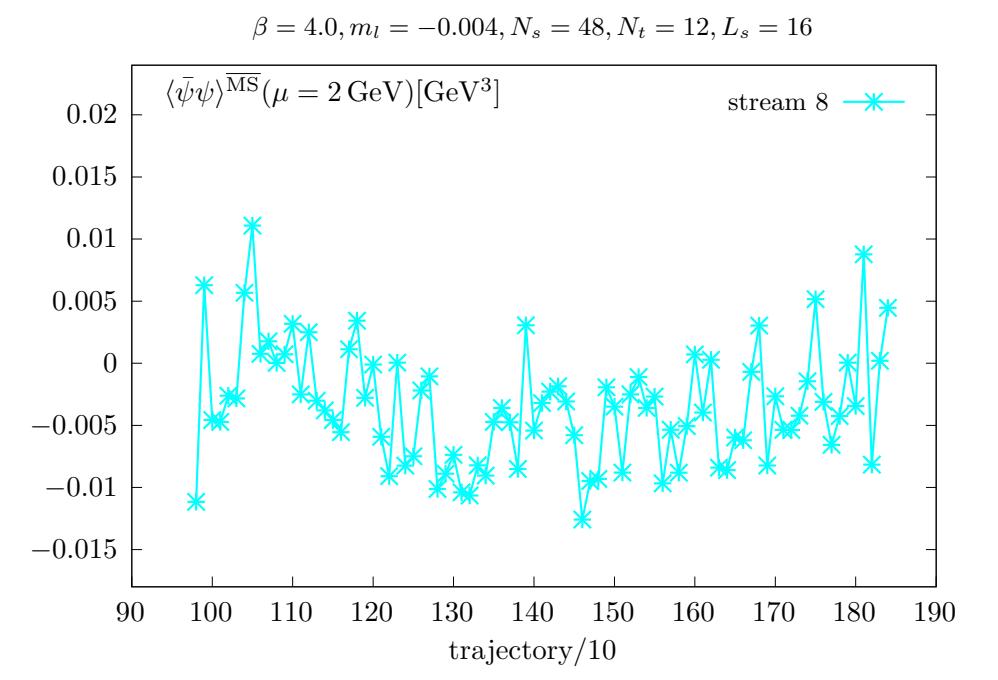
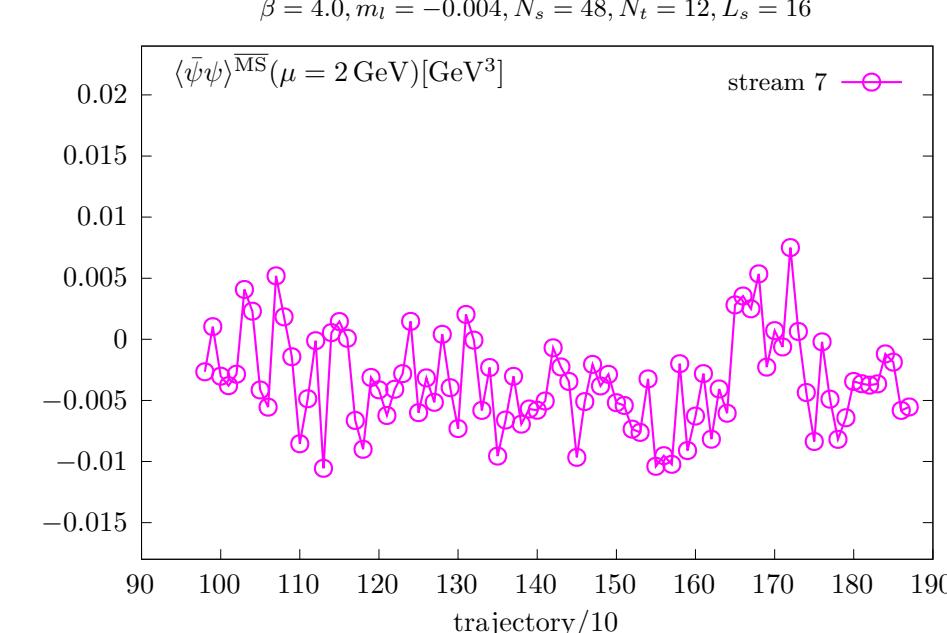
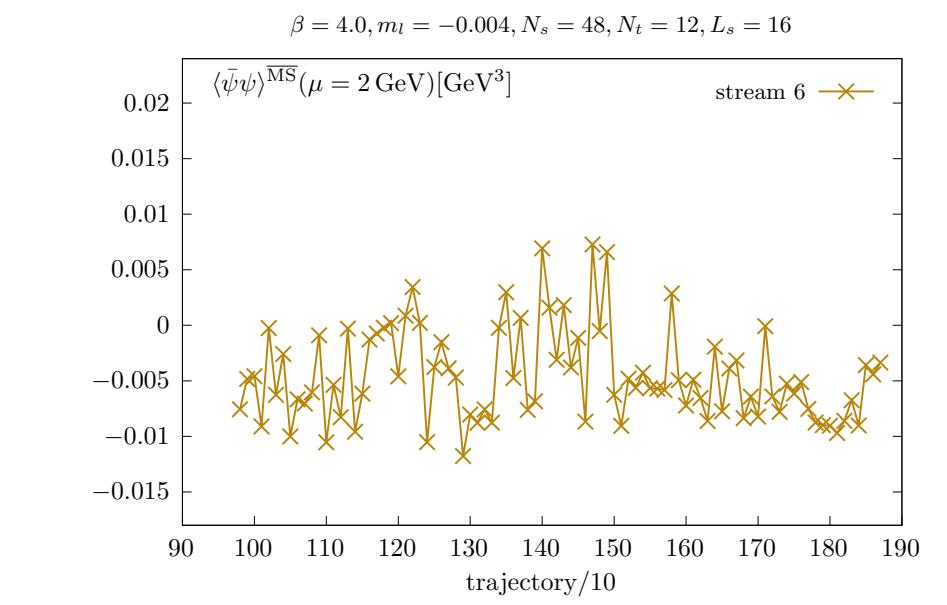
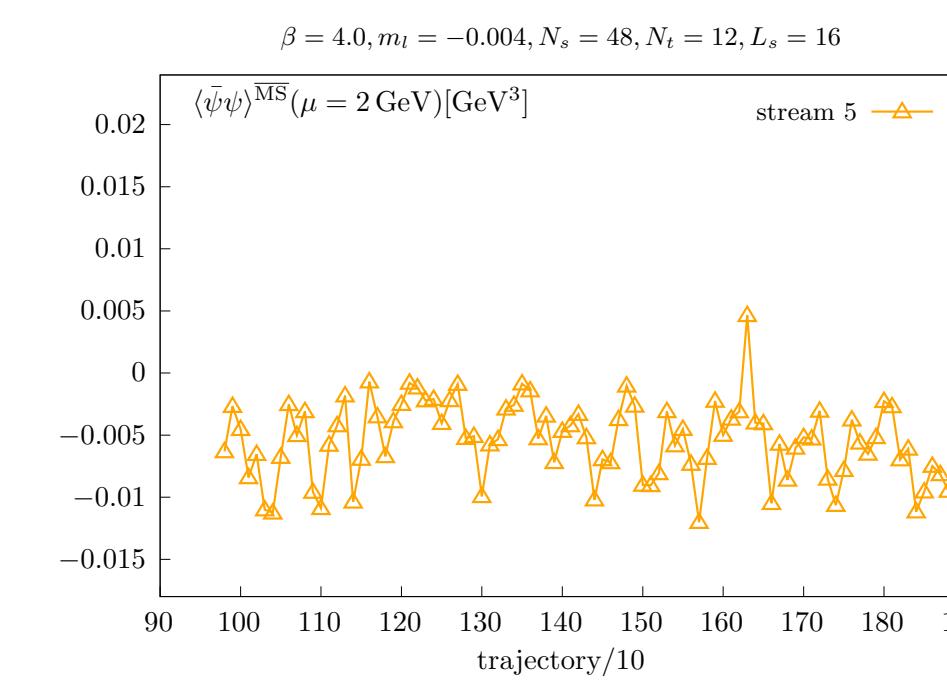
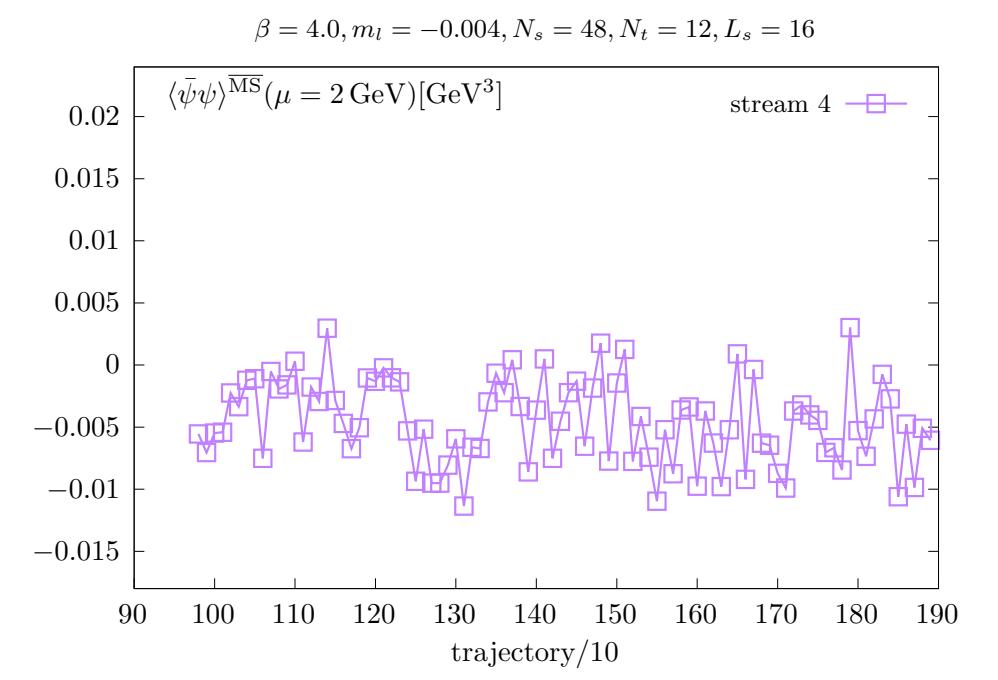
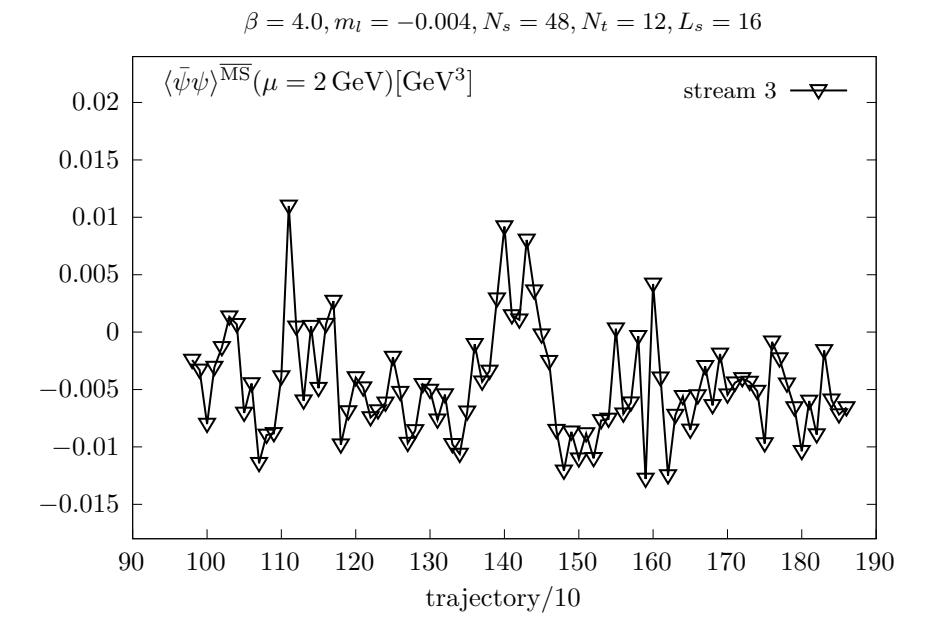
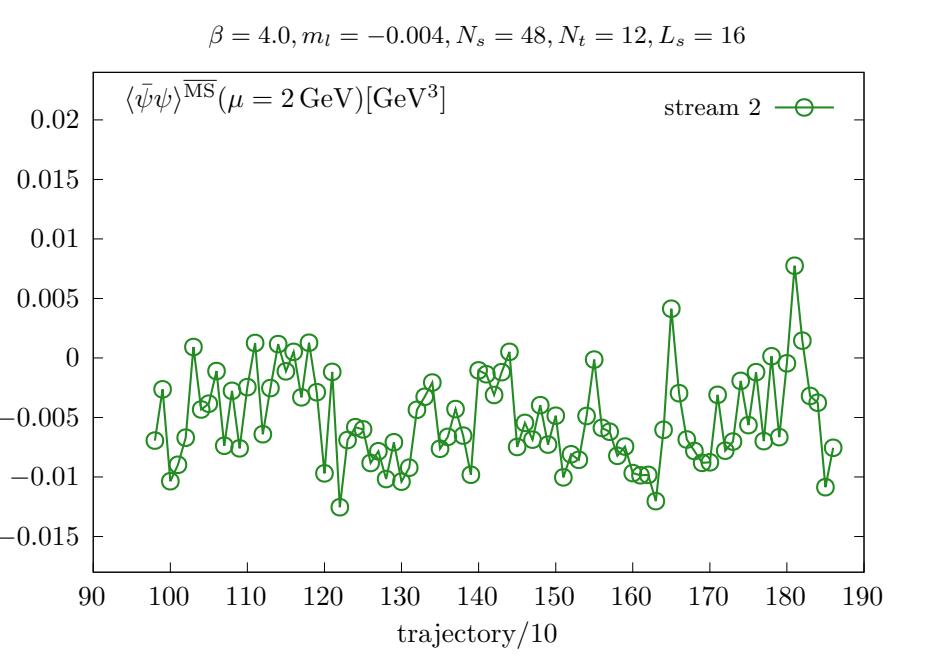
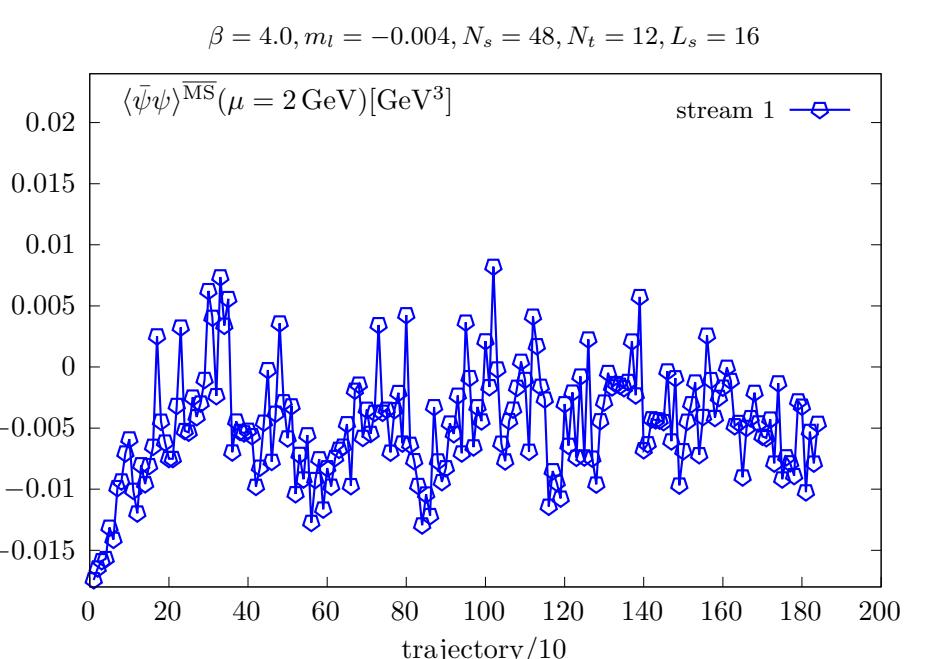
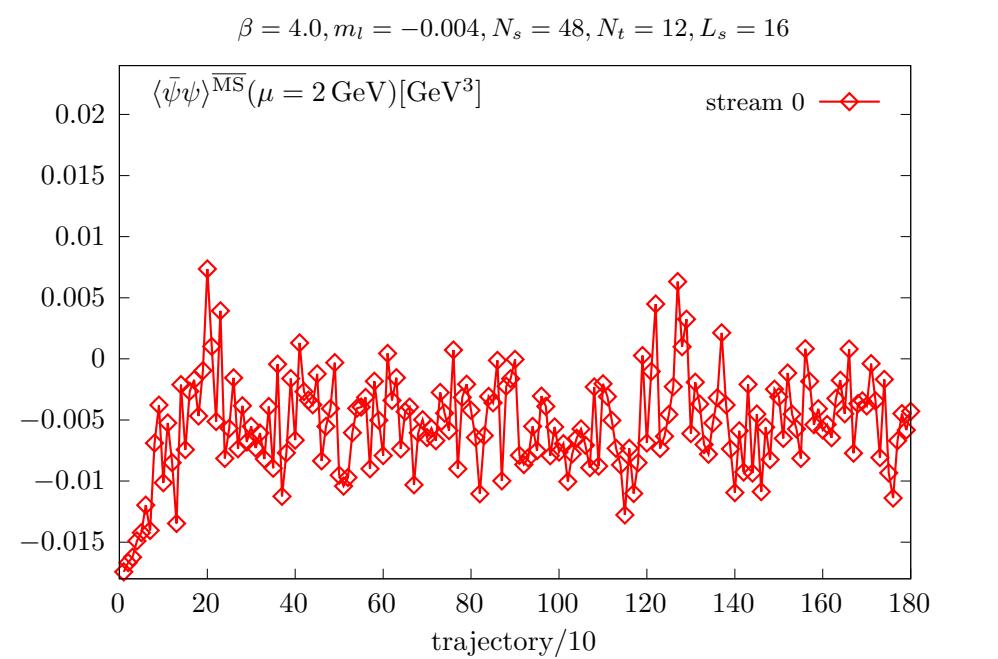
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- The transition mass point is around 2.7-4.5 MeV (FLAG Review '21: $m_{ud}^{\overline{\text{MS}}}(2 \text{ GeV}) = 3.381(40) \text{ MeV}$)
- Chiral susceptibilities seem to be function of total quark mass











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