# Tensor network study of CP(1) model with theta-term



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German Japanese Workshop 2024 25-27.09.2024 @Mainz

#### 2005-2008 Postdoc time







#### Contents

- Introduction of tensor networks
- Application to  $CP(1) + \theta$  model

### Tensor networks (TN)

- Why TN? Answer: No sign-problem Real-time, Finite density, θ-term, ...
- What's TN? we address later



 TN is used as a representation of wave function or partition function (path integral)

local interaction

$$Z \equiv \int [d\phi] e^{-S[\phi]}$$

local interaction

$$Z \equiv \int [d\phi] e^{-S[\phi]} = \sum_{i=1}^{n} \sum_{i=1}^{n} \sum_{j=1}^{n} \sum_{j=1}^{n} \sum_{j=1}^{n} \sum_{i=1}^{n} \sum_{j=1}^{n} \sum_{j=1}^{n} \sum_{j=1}^{n} \sum_{j=1}^{n} \sum_{j=1}^{n} \sum_{i=1}^{n} \sum_{j=1}^{n} \sum_{j=$$

Periodic BC



- $\sum_{\ldots,i,j,k,l,\ldots} \cdots T_{ijkl} T_{mnio} \cdots$
- Non-trivial step, but
- OK for scalar, gauge, and fermion fields as long as the interaction is local
  - Scalar : Shimizu 2012, Sakai+ 2018
  - Gauge: Meurice+ 2013, Nishimura+ 2021, Fukuma+ 2021
  - Fermions: Shimizu+Kuramashi 2014

 $Z \equiv \int [d\phi] e^{-S[\phi]} = \left(\sum_{i,j,k,l,l} \cdots T_{ijkl} T_{mnio} \cdots \right)$  $\operatorname{cost} \propto \chi^{2V}$  for  $1 \leq i, j, \ldots \leq \chi$ 

Periodic BC

Bond dimension





Information compression using singular value decomposition (SVD)

Tensor renormalization group (TRG) Levin+Nave 2007



### Application to $CP(1) + \theta$

# 2d CP(1) model

#### toy model of QCD

 $z \in \mathbb{C}^2$ 

- asymptotic free confinement instanton

 $|z|^2 = \sum_{a=1}^{2} z_a^* z_a = 1$ 

 $S = \int d^2x \left( \beta |D_{\mu}z|^2 + i \frac{\theta}{2\pi} \epsilon_{\mu\nu} \partial_{\mu}A_{\nu} \right)$ 

U(1) auxiliary field

on the lattice

continuum

on the lattice 
$$\begin{array}{c} & & \text{mod } 2\pi \\ S_{\text{lat}} = -2\beta \sum_{x,\mu} \left[ z^{\dagger}(x)z(x+\hat{\mu})U_{\mu}(x) + (c.c.) \right] - i\frac{\theta}{2\pi} \sum_{x} q(x) \quad \in \mathbb{C} \end{array}$$

$$U_{\mu} = e^{iA_{\mu}}$$

sign problem!









Haldane's conjecture : mass gap of O(3) vanishes at  $\theta = \pi$ Haldane 1983  $( \simeq CP(1) )$ 



Affleck+Haldane 1987 expected universality class : k = 1 SU(2) WZNW Wess+Zumino 1971, Novikov 1981, Witten 1984

Azcoiti+ 2007 Monte Carlo, imaginary  $\theta$ 



Critical region is observed but the universality class is not identified

⇒ This should be investigated by sign problem free method : tensor network!

Kawauchi+ST 2018 Tensor network

$$\chi = \frac{1}{L^2} \frac{\partial^2}{\partial \theta^2} \log \mathcal{Z}(\theta)|_{\theta = \pi} \propto L^b$$



1st 
$$: b = 2$$
  
2nd  $: b = \gamma/\nu < 2$ 

 $\beta = 0.5$  is roughly consistent with MC Azcoiti+ 2007

But universality class is not determined

 $\mathsf{D}_\beta$  ,  $\mathsf{D}_\theta$  : truncation order



no critical point is found up to  $\beta$ =1.1

### What's new in our study

We make two improvements

• Using better initial tensor

• New analysis based on CFT

#### Improvement of initial tensor

Kawauchi+ST 2017, Nakayama+ 2021

• Previous studies use character expansion for θ-term

 $N_{A}$ 

Hassan+ 1995

$$e^{i\frac{\theta}{2\pi}q_p} = \sum_{n \in \mathbb{Z}} e^{in(A_1 + A_2 - A_3 - A_4)} C_n(\theta) \xrightarrow{\qquad } \propto \frac{1}{n} \quad : \text{slow convergence!}$$

 $\Rightarrow$  large truncation error

• New method : Quadrature

 $hat{}\pi$ 

Gauss-Legendre quadrature

Gauge field : 
$$\int_{-\pi}^{\pi} dAf(A) \approx \sum_{i=1}^{n} w_i f(A_i) \qquad i \Rightarrow \text{tensor's index}$$
weight
Complex scalar field : 
$$\int_{|z|^2 = 1} dz g(z) \approx \sum_{p=1}^{N_z} w_p^{(z)} g(z_p) \qquad p \Rightarrow \text{tensor's index}$$
Genz+Keister 1996

#### Improvement of initial tensor



New method shows better precision

#### New analysis method

• Previous studies use susceptibility

Kawauchi+ST 2017, Nakayama+ 2021

$$\chi = \frac{1}{L^2} \frac{\partial^2}{\partial \theta^2} \log \mathcal{Z}(\theta)|_{\theta=\pi} \propto L^b \qquad \qquad \text{fitting around } \theta = \pi \Rightarrow \text{ fitting range ?}$$

• CFT : central charge and scaling dimensions



#### Central charge



#### Scaling dimensions

$$x_i = h_i + \bar{h}_i = \frac{1}{2\pi} \log\left(\frac{\lambda_0}{\lambda_i}\right) \qquad x_1, x_2, x_3, x_4$$



### Summary

- We have analyzed phase structure of 2d CP(1)+θ using tensor network
- Two improvements
  - Better (more precise) initial tensor
  - New analysis using CFT : central charge and scaling dimensions
- Critical point region starts from  $\beta \approx 0.55$  and k=1 SU(2) symmetry is observed at  $\beta \approx 0.595$



#### German-Japanese Workshop 202? in Kanazawa If our proposal is accepted!!!



#### Future

- Mass gap scaling
- Continuum limit?

#### Back up

Kawauchi+ST 2018 Tensor network



Critical region is observed but the universality class is not identified

#### c=1, k=1 WZNW

scaling dimension for c=1 free boson CFT

$$x_i = x_{M,N} = \frac{M^2}{4R^2} + R^2 N^2, \quad M, N \in \mathbb{Z}$$
  $R:$  compact radius

for  $R = \sqrt{\frac{1}{2}} \Rightarrow$  SU(2) symmetry appears  $\Rightarrow$  k=1 WZNW universality class Ginsparg 1988

$$x_{\pm 1,0} = x_{0,\pm} = \frac{1}{2} \Rightarrow \text{quartet}$$

X



#### Nature of transition at $\beta = 0.1$

Kawauchi+ST 2017

$$\chi = \frac{1}{V} \frac{\partial^2}{\partial \theta^2} \log \mathcal{Z}(\theta)|_{\theta = \pi} \propto V^a$$

00

$$\begin{cases} 1 \text{st} &: a = 1 \\ 2 \text{nd} &: a = \gamma/\nu < 1 \end{cases}$$



 $D_{cut} = 64$ , TRG for coarse-graining

Haldane's conjecture : mass gap of O(3) vanishes at  $\theta = \pi$ Haldane 1983  $( \simeq CP(1) )$ 



Bietenholtz et al. 1995, Wiese et al. 2012, de Forcrand et al. 2012, Azcoiti et al. 2012, Alles et al. 2014

 $O(3) + \theta$  was intensively studied by MC and Haldane's conjecture is confirmed And universality class is consistent with k=1 WZNW model

Wess & Zumino 1971, Novikov 1981, Witten 1984

Tensor renormalization group (TRG) PRL99,120601(2007)



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Bond dimension

 $T_{ijkl}$ 

 $1 \leq i,j,\ldots \leq \chi$ 

 $\Leftrightarrow$ 



Tensor renormalization group (TRG) PRL99,120601(2007)



$$1 \le i, j, \dots \le \chi$$

$$\chi^2 \times \chi^2 \text{ matrix}$$

$$\Leftrightarrow \quad T_{ijkl} = M_{(ij)(kl)}$$



Tensor renormalization group (TRG) PRL99,120601(2007)

 $M \in \mathbb{C}^{\chi^2 \times \chi^2}$ 

Singular Value Decomposition(SVD)  $M_{ab} = \sum_{m} U_{am} \sigma_m (V^{\dagger})_{mb}$ unitary matrix  $\sigma_1 \ge \sigma_2 \ge \dots \ge 0$ : singular value (non-negative)

$$\Rightarrow T_{ijkl} = M_{(ij)(kl)}$$



Tensor renormalization group (TRG) PRL99,120601(2007)

 $M \in \mathbb{C}^{\chi^2 \times \chi^2} \ \Rightarrow \mathrm{TN}$  is sign-problem-free

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$$\begin{split} T_{ijkl} &= M_{(ij)(kl)} \\ & \overset{\text{SVD}}{=} \sum_{m=1}^{\chi^2} U_{(ij)m} \sigma_m V_{m(kl)}^{\dagger} \end{split}$$

Tensor renormalization group (TRG) PRL99,120601(2007)

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Singular Value Decomposition(SVD)  $M_{ab} = \sum_{m} U_{am} \sigma_m (V^{\dagger})_{mb}$ unitary matrix  $\sigma_1 \ge \sigma_2 \ge \dots \ge 0$ : singular value (non-negative)

$$\approx \qquad \Longleftrightarrow \qquad T_{ijkl} = M_{(ij)(kl)}$$
  
truncation  
$$\approx \sum_{m=1}^{X} U_{(ij)m} \sigma_m V_{m(kl)}^{\dagger}$$
  
truncation of SVD = information compression

## Tensor network rep. of ${\cal Z}$

depends on property of field and interaction

- Scalar field (non-compact)
  - Orthonormal basis expansion

Shimizu mod.phys.lett. A27,1250035(2012), Lay & Rundnick PRL88,057203(2002)

- Gauss Hermite quadrature Sakai et al., JHEP03(2018)141

• Gauge field (compact : SU(N), CP(N) etc.)

 Character expansion : maintain symmetry, better convergence Meurice et al., PRD88,056005(2013)

• Fermion field (Dirac/Majorana)

Shimizu & Kuramashi PRD90,014508(2014), ST & Yoshimura PTEP(2015)043B01

- Grassmann number  $\theta^2=0 \rightarrow finite$  sum

In principle, we can treat any fields

$$e^{\phi\theta} = 1 + \phi\theta = \sum_{n=0}^{1} (\phi\theta)^n$$



e.g. 2D Ising model

$$\mathcal{Z} = \sum_{\{s\}} \exp\left(\sum_{\langle x,y \rangle} \beta s_x s_y\right) = \sum_{\{s\}} \prod_{\langle x,y \rangle} \exp(\beta s_x s_y)$$
  

$$= (\cosh\beta)^{2V} \sum_{\{s\}} \prod_{\langle x,y \rangle} \sum_{i_{xy}=0}^{1} (s_x s_y \tanh\beta)^{i_{xy}} \qquad V = \text{# of lattice sites}$$
  

$$\exp(\beta s_x s_y) = \cosh(\beta s_x s_y) + \sinh(\beta s_x s_y)$$
  

$$= \cosh\beta + s_x s_y \sinh\beta$$
  

$$= \cosh\beta (1 + s_x s_y \tanh\beta)$$
  

$$= \cosh\beta \sum_{i_{xy}=0}^{1} (s_x s_y \tanh\beta)^{i_{xy}} \qquad x = \pm 1$$
  

$$\frac{1}{i_{xy}} = 0$$
  

$$\exp(\partial s_x s_y) = \frac{1}{i_{xy}} \left( \sum_{i_{xy}=0}^{1} (s_x s_y \tanh\beta)^{i_{xy}} + \sum_{i_{xy}=0}^{1} (s_x$$

$$\mathcal{Z} = \sum_{\{s\}} \exp\left(\sum_{\langle x,y \rangle} \beta s_x s_y\right) = \sum_{\{s\}} \prod_{\langle x,y \rangle} \exp(\beta s_x s_y)$$
$$= (\cosh\beta)^{2V} \sum_{\{s\}} \prod_{\langle x,y \rangle} \sum_{i_xy=0}^{1} (s_x s_y \tanh\beta)^{i_xy}$$
$$= (\cosh\beta)^{2V} \sum_{\{i\}} \sum_{\{s\}} \prod_{\langle x,y \rangle} (s_x \sqrt{\tanh\beta} \cdot s_y \sqrt{\tanh\beta})^{i_xy}$$

e.g. 2D Ising model



e.g. 2D Ising model  $\mathcal{Z} = \sum_{\{s\}} \exp\left(\sum_{\langle x,y \rangle} \beta s_x s_y\right) = \sum_{\{s\}} \prod_{\langle x,y \rangle} \exp(\beta s_x s_y)$  $= (\cosh \beta)^{2V} \sum_{\{s\}} \prod_{\{x,y\} > i_{xy} = 0}^{1} (s_x s_y \tanh \beta)^{i_{xy}}$  $= (\cosh\beta)^{2V} \sum \sum \prod (s_x \sqrt{\tanh\beta} \cdot s_y \sqrt{\tanh\beta})^{i_{xy}}$  $\{i\} \ \{s\} \ \langle x, y \rangle$  $= (\cosh\beta)^{2V} \sum \sum \prod (s_x \sqrt{\tanh\beta})^{i_{xy}} (s_x \sqrt{\tanh\beta})^{i_{xz}} (s_x \sqrt{\tanh\beta})^{i_{xw}} (s_x \sqrt{\tanh\beta}$  $\{i\} \ \{s\} \ x$  $= (\cosh\beta)^{2V} \sum \sum \prod (\sqrt{\tanh\beta})^{i_{xy}+i_{xz}+i_{xw}+i_{xv}} s_x^{i_{xy}+i_{xz}+i_{xw}+i_{xv}}$  $\{i\}\ \{s\}\ x$  $= (\cosh\beta)^{2V} \sum_{\{i\}} \prod_{x} (\sqrt{\tanh\beta})^{i_{xy}+i_{xz}+i_{xw}+i_{xv}} \sum_{s_x=\pm 1} s_x^{i_{xy}+i_{xz}+i_{xw}+i_{xv}} \text{ summation is done}$  $= (\cosh\beta)^{2V} \sum_{\{i\}} \prod_{x} (\sqrt{\tanh\beta})^{i_{xy}+i_{xz}+i_{xw}+i_{xv}} 2\delta(\operatorname{mod}(i_{xy}+i_{xz}+i_{xw}+i_{xv},2))$  $\{i\}$  x  $=T_{i_{xy}i_{xz}i_{xw}i_{xv}}$ new d.o.f. : index of tensor

e.g. 2D Ising model

 $\mathcal{Z} = 2^{V} (\cosh \beta)^{2V} \sum_{\dots, i, j, k, l, m, n, o, \dots} \cdots T_{ijkl} T_{mnio} \cdots$ 

 $T_{ijkl} = (\sqrt{\tanh\beta})^{i+j+k+l} \delta(\operatorname{mod}(i+j+k+l), 2)$ 

 $\begin{bmatrix} T_{0000} & T_{0001} & T_{0010} & T_{0011} \\ T_{0100} & T_{0101} & T_{0110} & T_{0111} \\ T_{1000} & T_{1001} & T_{1010} & T_{1011} \\ T_{1100} & T_{1101} & T_{1110} & T_{1111} \end{bmatrix} = \begin{bmatrix} 1 & 0 & 0 & \tanh \beta \\ 0 & \tanh \beta & \tanh \beta & 0 \\ 0 & \tanh \beta & \tanh \beta & 0 \\ \tanh \beta & 0 & 0 & (\tanh \beta)^2 \end{bmatrix}$ 

size and elements of tensor depend on a model

#### MC

Boltzmann weight is interpreted as probability

Importance sampling

Statistical errors

Sign problem may appear

Critical slowing down

#### TRG

Tensor network rep. of partition function (no probability interpretation)

Compression of tensor by SVD, Optimization

Systematic errors (truncated SVD)

No sign problem

∵ no probability

Efficiency of compression gets worse around criticality

can be improved by TNR, Loop-TNR in 2D system Evenbly & Vidal 2014, Gu et al., 2015 From point of view of MC

From point of view of TN

Space of theory



Is there a region where any method cannot access ?

Is entanglement problem NP-hard ?

### Hierarchy of singular value



• Near criticality: hierarchy gets worse (large S)

like critical slowing down in MC

Tensor network renormalization (TNR) Evenbly&Vidal 2014 can help the situation

#### **Renormalization group**

#### TRG

$$H(K; \{s\}) = \sum_{i} K_i \mathcal{O}_i(\{s\})$$
  
s : spins

$$\mathcal{Z} = \sum_{\{s\}} e^{-\beta H}$$

Block spin transf., Migdal-Kadanoff RG

$$\begin{split} H &= \sum_{i} K_{i} \mathcal{O}_{i} \xrightarrow{\mathcal{R}} H' = \sum_{i} K'_{i} \mathcal{O}'_{i} \\ K' &= \mathcal{R}_{K}(K) \\ K^{*} &= \mathcal{R}_{K}(K^{*}) \text{ : fixed point} \end{split}$$

target: critical exponent etc

$$T_{ijkl}(K)$$

$$i, j, k, l : \text{indexes}$$

$$\mathcal{Z} = \sum_{i,j,k,l,\dots} \prod T_{ijkl} \cdots$$
SVD + contraction
$$\sum_{i,j,k,l,\dots} \prod T_{ijkl} \cdots \xrightarrow{\mathcal{R}} \sum_{i,j,k,l,\dots} \prod T'_{ijkl} \cdots$$

$$T' = \mathcal{R}_T(T)$$

$$T^* = \mathcal{R}_T(T^*) : \text{fixed point tensor}$$
target : partition function etc

Kadanoff et al, Rev.Mod.Phys.86,647(2014)

#### SVD

For simplicity rank- $k m \times n$  real matrix  $A (m \ge n \ge k)$  is given by

 $A = U \Lambda V^T$  (full SVD)

 $\begin{cases} U: m \times m \text{ orthonormal matrix}: U=(u_1, u_2, ..., u_m), \quad U^T U=U U^T = I_m \\ V: n \times n \text{ orthonormal matrix}: V=(v_1, v_2, ..., v_n), \quad V^T V=V V^T = I_n \\ \Lambda: m \times n \text{ diagonal matrix}: \Lambda = \text{diag}(\sigma_1, \sigma_2, ..., \sigma_n) \end{cases}$ where  $\sigma_1 \ge \sigma_2 \ge \dots \ge \sigma_k > \sigma_{k+1} = \dots = \sigma_n = 0$ )



#### Best approximation of A ?

