



# Effective Field Theory I

Hans-Werner Hammer

Technische Universität Darmstadt



Deutsche  
Forschungsgemeinschaft



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1. **Effective Field Theory I:** Basic ideas
2. **Effective Field Theory II:** Nonrelativistic short-range EFT
3. **Effective Field Theory III:** Applications in hadrons and nuclei

## Literature

G.P. Lepage, "What is renormalization", TASI Lectures 1989, arXiv:hep-ph/0506330

D.B. Kaplan, "Effective Field Theories", NNPS Lectures 1995, arXiv:nucl-th/9506035

G.P. Lepage, "How to renormalize the Schrödinger equation", Swieca Lectures 1997, arXiv:hep-ph/0506330

E. Braaten, HWH, "Universality in few-body systems with large scattering length", Phys. Rep. **428**, 259 (2006), arXiv:cond-mat/0410417

# Resolution of a Microscope



TECHNISCHE  
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DARMSTADT

- Resolution of microscope is limited

- Diffraction effects:

Light is EM wave ( $\lambda \approx 550 \text{ nm}$ )

- Resolution:  $d = \frac{\lambda}{A_N}$

- Aperture:  $A_N \approx 0.95 \dots 1.5$

$$\longrightarrow d_{\max} \approx 0.4 \mu\text{m}$$

- Quantum mechanics:

particles have wave character  $\Rightarrow$  matter waves

$$\lambda = \frac{h}{p} \quad (\text{de Broglie, 1924}) \quad \Rightarrow \quad \text{QM microscope}$$

$\Rightarrow$  intrinsic resolution scale determined by typical  $p$





## Difference between typical biological and physical Systems?

- DNA molecule



many energy scales of comparable size

- Hydrogen atom: few separated scales

- Hierarchy:  $m_e \ll m_p \ll m_{Z_0} \ll \dots$

- Physical quantities have units  $\rightarrow$  dimensional analysis

(vgl. M. Planck, Sitzungsber. Preuß. Akad. Wiss. **5**, 479 (1899)), “Planck units”)

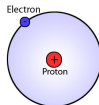
$$E_0 = m_e \times f\left(\alpha, \frac{m_e}{m_p}, \frac{m_e}{m_{Z_0}}, \dots\right)$$

- $f$  contains dynamics

- How to exploit scale separation?

$\Rightarrow$  Effective (Field) Theories  $(\iff)$  Calculation of  $f$

- Quantum bound state of  $e^-$  and  $p$



- Energy levels (lowest order): Schrödinger-Eq. for  $V(r) = -e^2/r$

$$E = E_0 = -\frac{m_e \alpha^2}{2n^2}, \quad \alpha = \frac{e^2}{4\pi}$$

- Corrections beyond leading order:  $E = E_0 \left[ 1 + \mathcal{O}(\alpha, \frac{m_e}{m_p}, \dots) \right]$

- corrections from EM interaction: fine structure  $\mathbf{L} \cdot \mathbf{S} \sim \alpha^4, \dots$
- corrections from proton structure: finite proton mass  $m_p$ 
  - hyperfine structure:  $\mu_p$
  - finite proton size:  $r_p$



- Estimate effect of  $r_p$  on  $E_0$
- Natural scales:  $a_0 = 1/(m_e \alpha) \sim 0.5 \text{ \AA}$  (Bohr radius)
- Proton charge radius:  $G_E(q^2) = 1 + q^2 \overbrace{G'_E(0)}^{r_p^2/6} + \dots$   
guess:  $G'_E(0) \sim 1/m_p^2, \quad q \sim 1/a_0 \implies \left(\frac{m_e \alpha}{m_p}\right)^2 \sim 10^{-11}$
- Size of correction to  $E_0$ :  $10^{-11} E_0 \rightarrow 10^{-11} \times 3 \cdot 10^{15} \text{ Hz} \sim 30 \text{ kHz}$
- Estimated enhancement in muonic hydrogen by factor  $(m_\mu/m_e)^2 \approx 200^2$   
 $\Rightarrow$  higher precision in muonic atoms possible  
(R. Pohl et al., Nature **466**, 213 (2010))
- How to construct an effective theory to calculate the corrections from  $m_p, \mu_p, r_p$  ?



- How to construct an effective theory to calculate these corrections?

- Need to identify:

- ▣ Low and high scales in the problem:  $M_{lo}, M_{hi}, \dots$
- ▣ Active degrees of freedom: coordinates, particles, ...
- ▣ Symmetries to constrain interactions
- ▣ **Power counting**: organize the theory through expansion in  $M_{lo}/M_{hi}$

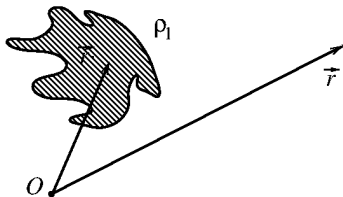
- Guiding principle:

Include long-range physics explicitly, parametrize short-range physics in low-energy constants

- Consider an example from classical electrodynamics

## Example: Multipole Expansion

- Separation of scales:  $\langle r' \rangle \ll r$
- Short-range physics: expand in  $\frac{\langle r' \rangle}{r}$   
→ “power counting”
- Breakdown scale:  $\langle r' \rangle \sim r$

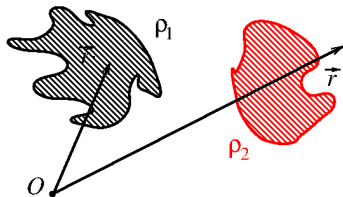


$$\phi(\mathbf{r}) = \sum_{l,m} \frac{C_l Y_{lm}(\Omega)}{r^{l+1}} \overbrace{\int d^3 r' \rho_1(\mathbf{r}') Y_{lm}^*(\Omega') (r')^l}^{q_{lm}}$$
$$\frac{1}{|\mathbf{r} - \mathbf{r}'|} = \sum_{l,m} C_l Y_{lm}(\Omega) Y_{lm}^*(\Omega') \frac{r'^l_{<}}{r^{l+1}_{>}}$$

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- Long-range physics: include explicitly
- Use these ideas in the framework of Quantum Field Theory



Quantum Field Theory is “only” way to satisfy quantum mechanics, Lorentz invariance, and cluster decomposition

⇒ to calculate the most general S-matrix consistent with given symmetries for any theory below some scale simply use the most general effective Lagrangian  $\mathcal{L}_{\text{eff}}$  consistent with these principles in terms of the appropriate asymptotic states

(Weinberg, 1979)

- **Nonrelativistic systems:** Lorentz invariance  $\longrightarrow$  Galilei invariance  
⇒ QM framework possible
- **Need “power counting scheme”** to organize the infinitely many terms in  $\mathcal{L}_{\text{eff}}$   
⇒ **predictive power**

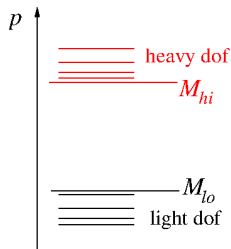


- Construct most general  $\mathcal{L}_{\text{eff}}$  respecting underlying symmetries

- Exploit separation of scales:  $p, M_{lo} \ll M_{hi}$
- Non-renormalizable theory, but only finite number of operators in  $\mathcal{L}$  contribute at each order

⇒ Low-energy constants (LEC)

- Symmetries and light dof's must be known



- Work at low energies:  $p \ll M_{hi}$
- Fix LEC's from matching (to experiment, other theory, ...)
- Calculate observables in expansion in  $p/M_{hi}, M_{lo}/M_{hi}$

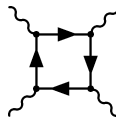
⇒ limited range of applicability

- **Classic Example** (Euler, Heisenberg, 1936)
- **Photon energy  $\omega$ , electron mass  $m_e$**
- **Separation of scales:  $\omega \ll m_e$**
- **Only photons are active dof**  $\Rightarrow$  **theory simplifies**

$$\Rightarrow \mathcal{L}_{QED}[\psi, \bar{\psi}, A_\mu] \rightarrow \mathcal{L}_{eff}[A_\mu]$$

- **Invariants:**  $F_{\mu\nu}F^{\mu\nu}, (F_{\mu\nu}F^{\mu\nu})^2, (F_{\mu\nu}\tilde{F}^{\mu\nu})^2, \dots$

with field strength tensor  $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$ , dual tensor  $\tilde{F}^{\mu\nu} = \frac{1}{2}\epsilon^{\mu\nu\alpha\beta}F_{\alpha\beta}$

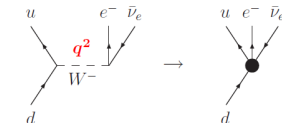


$$\mathcal{L}_{eff} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \underbrace{c_1}_{LEC}(F_{\mu\nu}F^{\mu\nu})^2 + \underbrace{c_2}_{LEC}(F_{\mu\nu}\tilde{F}^{\mu\nu})^2 + \dots$$

- **Estimate size & energy dependence of cross section**  $\longrightarrow$  **exercise**
- **Observation in Pb-Pb Collisions at LHC**  
(ATLAS Kollaboration, Nature Physics **13**, 852 (2017))

## ■ Weak decays

- mediated by  $W^\pm$  bosons ( $M_W \approx 80 \text{ GeV}$ )
- energy release in neutron  $\beta$ -decay  $\sim 1 \text{ MeV}$  [ $n \rightarrow pe^- \bar{\nu}_e$ ]
- energy release in kaon decays  $\sim 300 \text{ MeV}$  [ $K \rightarrow \pi e \nu$ ]



$$\frac{e^2}{8 \sin^2 \theta_W} \times \frac{1}{M_W^2 - q^2} \xrightarrow{q^2 \ll M_W^2} \frac{e^2}{8 M_W^2 \sin^2 \theta_W} \left\{ 1 + \frac{q^2}{M_W^2} + \dots \right\} = \frac{G_F}{\sqrt{2}} + \dots$$

$\Rightarrow$  Fermi Theory of Weak Interaction



- **Energy/momentum expansion** (derivatives acting on fields)
  - **Dimensional analysis** (breakdown scale  $M_{hi}$ )
    - ▣ derivatives  $\rightarrow$  powers of typical momentum  $q$ 
      - one derivative:  $\partial \sim q/M_{hi}$
      - vertex w/  $N$  derivatives:  $\partial^N \sim (q/M_{hi})^N$
      - terms with more derivatives are suppressed if  $q \ll M_{hi}$
    - ▣ **Loops**: additional powers of momenta from vertices and loop momenta
- $\Rightarrow$  **loops are generally suppressed compared to trees (exceptions!)**

$$\begin{array}{ccc} \text{X} & \sim (q/M_{hi})^N & \text{Loop diagram} \end{array} \quad \sim (q/M_{hi})^{2N} \underbrace{\int d^4 p f(p, q)}_{(q/M_{hi})^M}$$



## 1. Construct $\mathcal{L}$ respecting symmetries (e.g. gauge invariance of QED)

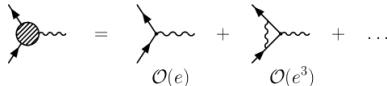
$$\psi \rightarrow \psi' = e^{i\alpha(x)}\psi, \quad A_\mu \rightarrow A'_\mu = A_\mu - \frac{1}{e}\partial_\mu\alpha(x)$$

## 2. Retain renormalizable interactions ( $D \leq 4$ ): $[\psi] = 3/2$ , $[A] = 1$

$$\text{keep } \underbrace{\bar{\psi}\gamma_\mu\psi A^\mu}_{D=4}, \quad \text{but drop } \underbrace{\bar{\psi}\sigma_{\mu\nu}\psi F^{\mu\nu}}_{D=5}, \quad \underbrace{(F_{\mu\nu}F^{\mu\nu})^2}_{D=8}$$

$$\Rightarrow \quad \mathcal{L} = \bar{\psi}(i\not{\partial} - \cancel{e}\not{A} - m)\psi - \frac{1}{4}F^{\mu\nu}F_{\mu\nu}$$

## 3. Calculate Feynman diagrams:


$$\text{Feynman diagram} = \text{Feynman diagram} + \text{Feynman diagram} + \dots$$

$\mathcal{O}(e)$                        $\mathcal{O}(e^3)$

## 4. Fix parameters from data $\rightarrow$ predictions

$$\text{Example: } \mu_e = -\frac{eg_e\mathbf{s}_e}{2m_e}, \quad g_e = 2 \left[ 1 + \frac{e^2}{8\pi^2} + \mathcal{O}(e^4) \right]$$



- **Renormalization:** method to make sense of infinities/UV sensitivities in QFT
  - ▣ Some observables are sensitive to short-distances  $\Rightarrow$  match to experiment and predict other observables
  - ▣ **Classical renormalizability:** redefinition of a finite number of terms sufficient to all orders (QED:  $e, m_e, Z_3, Z_2$ )
  - ▣ **EFT:** new structures appear at every order
- **Renormalizable theories have been very successful**  $\implies$  Standard Model (SM)
- **Non-renormalizable interactions** ( $D > 4$ ) are suppressed at low energies  $\longleftarrow$   
factors  $1/M_{hi}^{D-4}$
- **Modern understanding:** SM is also low-energy EFT, physics beyond SM can be included through non-renormalizable interactions