

1 Dimensional Analysis I: Planck units

Construct Planck length, mass, and time out of G , $\hbar$, c

$$[G] = \frac{m^3}{kg \cdot s^2} \quad \text{Newton's constant}$$

$$\text{Aside!} \quad \frac{kg \cdot m}{s^2} = [G] \frac{(kg)^2}{m^2}$$

$$[\hbar] = J \cdot s = \frac{kg \cdot m^2}{s}$$

$$F = G \frac{m_1 m_2}{r^2}$$

$$[c] = \frac{m}{s}$$

$$\hookrightarrow m_p = \left(\frac{\hbar c}{G} \right)^{1/2} \approx 2.2 \cdot 10^{-8} \text{ kg}$$

$$\text{Note!} \quad m_p \approx 1.7 \cdot 10^{-27} \text{ kg}$$

$$m_e \approx 9.1 \cdot 10^{-31} \text{ kg}$$

$$\hookrightarrow l_p = \left(\frac{\hbar G}{c^3} \right)^{1/2} \approx 1.6 \cdot 10^{-35} \text{ m}$$

$$\text{Note!} \quad r_e < 10^{-18} \text{ m}$$

$$r_p \approx 10^{-15} \text{ m}$$

$$\hookrightarrow t_p = \left(\frac{\hbar G}{c^5} \right)^{1/2} \approx 5.4 \cdot 10^{-44} \text{ s}$$

$$\text{Note!} \quad \tau_\Delta \approx 5.6 \cdot 10^{-24} \text{ s} \quad \text{Strong decay}$$

$$\left. \begin{array}{l} \tau_\mu \approx 2 \cdot 10^{-6} \text{ s} \\ \tau_{\pi^\pm} \approx 2.6 \cdot 10^{-8} \text{ s} \end{array} \right\} \text{ weak decay}$$

$$\left. \begin{array}{l} \tau_{\pi^0} \approx 8 \cdot 10^{-17} \text{ s} \\ \tau_\eta \approx 5 \cdot 10^{-19} \text{ s} \end{array} \right\} \text{ EM decay}$$

2 Dimensional analysis II: photon-photon scattering

$$\mathcal{L}_{\text{eff}} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + c_1 (F_{\mu\nu} F^{\mu\nu})^2 + c_2 (F_{\mu\nu} \tilde{F}^{\mu\nu})^2 + \dots$$

Dimensional analysis: $[2] = 4$

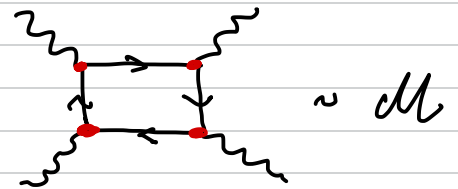
$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu \rightarrow [A_\mu] = 1, [\partial_\mu] = 1$$

$$\hookrightarrow [F_{\mu\nu}] = 2, [\tilde{F}_{\mu\nu}] = [\frac{1}{2} \epsilon_{\mu\nu\alpha\beta} F^{\alpha\beta}] = 2$$

Therefore we have $[c_1] = [c_2] = -4$

Leading QED contribution:

- m_e is only mass scale
- 4 couplings $\rightarrow$ 4 powers of e



$$\hookrightarrow c_1 \sim c_2 \sim \frac{e^4}{m_e^4}$$

$$\hookrightarrow c_1 (F_{\mu\nu} F^{\mu\nu})^2 + c_2 (F_{\mu\nu} \tilde{F}^{\mu\nu})^2$$

Amplitude for $E_\gamma \ll m_e$: $\partial_\mu \rightarrow (P_\mp)_\mu \sim E_\mp$

$$\hookrightarrow \mathcal{M} \sim \frac{e^4}{m_e^4} E_\gamma^4$$

$\gamma\gamma$ cross section for $E_\gamma \ll m_e$

$$\rightarrow \int \left(\frac{d^3 k}{2E} \right)^2 \delta^4(\dots)$$

$$\hookrightarrow d\sigma = \frac{1}{|i|} |\mathcal{M}|^2 \int d\text{LIPS}$$

$$\sim \frac{1}{E_\gamma^2} \left(\frac{e^4}{m_e^4} E_\gamma^4 \right)^2 (E_\gamma)^0 \sim \frac{e^8}{m_e^8} E_\gamma^6$$

alternatively from $c_{b2} \sim \frac{e^4}{m_e^4}$ and insight that correct dimension can only come from E_γ

3 Naturalness for square well potential

$$r_0 = \text{const.} \rightarrow \text{take } x = \kappa_0 r_0$$

$$\hookrightarrow a/r_0 = 1 - \frac{\tan(x)}{x}$$

$$P(a) da = \underbrace{P(x)}_{\rightarrow 1} dx = \left| \frac{dx}{da} \right| da = \frac{1}{|da/dx|} da$$

normalize later

$$(\tan x)' = 1 + \tan^2 x$$

$$\hookrightarrow \frac{da}{dx} = r_0 \frac{d}{dx} \left(1 - \frac{\tan x}{x} \right) = r_0 \left(\cancel{\frac{\tan x}{x^2}} - \frac{1 + \tan^2 x}{x} \right)$$

neglect

$$= r_0 \left(-\frac{1}{x} - x \left(1 - \frac{a}{r_0} \right)^2 \right) = -r_0 x \left(\frac{1}{x^2} + \left(1 - \frac{a}{r_0} \right)^2 \right)$$

$$\hookrightarrow P(a) da = \frac{da}{r_0^2 \kappa_0} \mathcal{N} \left(\frac{1}{\kappa_0^2 r_0^2} + \left(1 - \frac{a}{r_0} \right)^2 \right)^{-2} = \frac{1}{\kappa_0} \frac{\mathcal{N} da}{\frac{1}{\kappa_0^2} + (a - r_0)^2}$$

$x = \kappa_0 r_0$ normalization

$$\text{Normalization: } \int_{-\infty}^{\infty} P(a) da = 1 \rightarrow \mathcal{N} = \frac{1}{\pi}$$

$$\hookrightarrow P(a) da = \frac{1}{\frac{1}{\kappa_0^2} + (a - r_0)^2} \frac{da}{\pi \kappa_0}$$

Example: $\kappa_0 r_0 = 10$

$\hookrightarrow a \approx r_0$ is most likely if κ_0 uniformly distributed

